

Formation and Retention of Planets in Disks

Wladimir (Wlad) Lyra

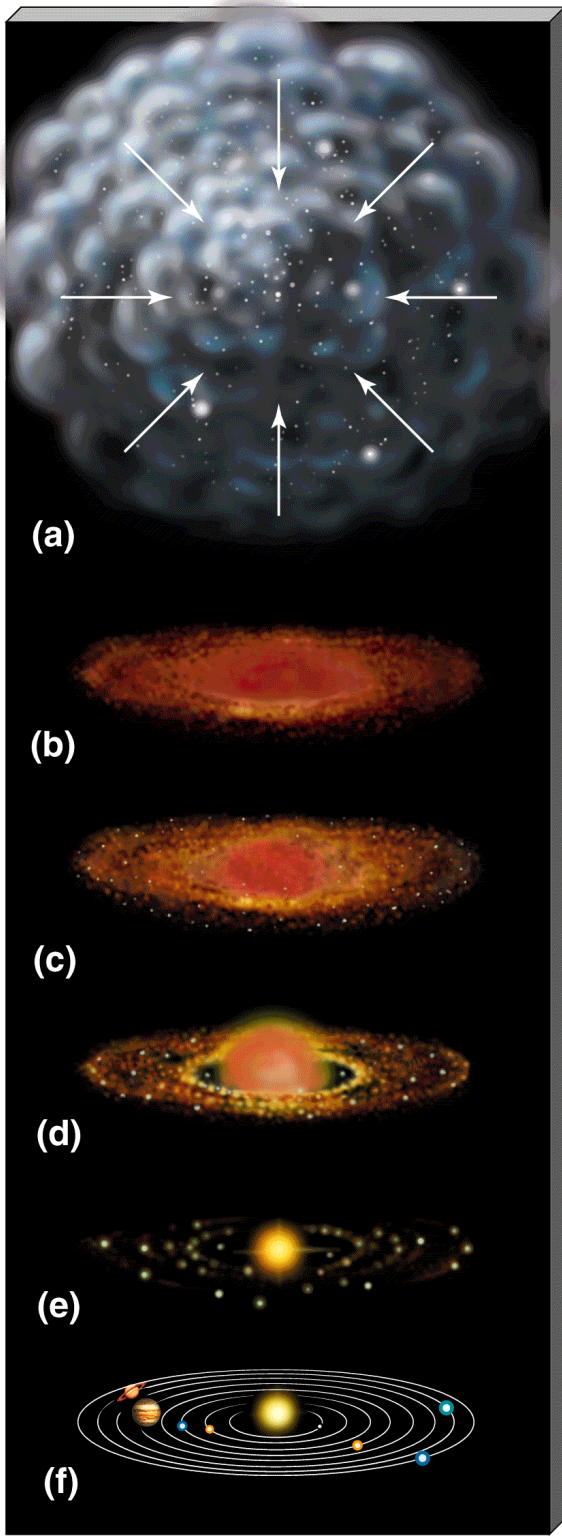
Sagan Fellow

NASA-JPL/Caltech



Collaborators:

Axel Brandenburg (Stockholm), Kees Dullemond (Heidelberg), Anders Johansen (Lund),
Brandon Horn (Columbia), Hubert Klahr (Heidelberg), Marc Kuchner (Goddard),
Mordecai-Mark Mac Low (AMNH), Sijme-Jan Paardekooper (Cambridge),
Nikolai Piskunov (Uppsala), Natalie Raettig (Heidelberg), Zsolt Sandor (Innsbruck),
Neal Turner (JPL), Andras Zsom (MIT).



A disk life story

Gas-rich phase (< 10 Myr)

T-Tauri Disks

Accretion and Planet Formation

Thinning phase (~ 10 Myr)

Transitional Disks

Planet retention

Gas-poor phase (> 10 Myr)

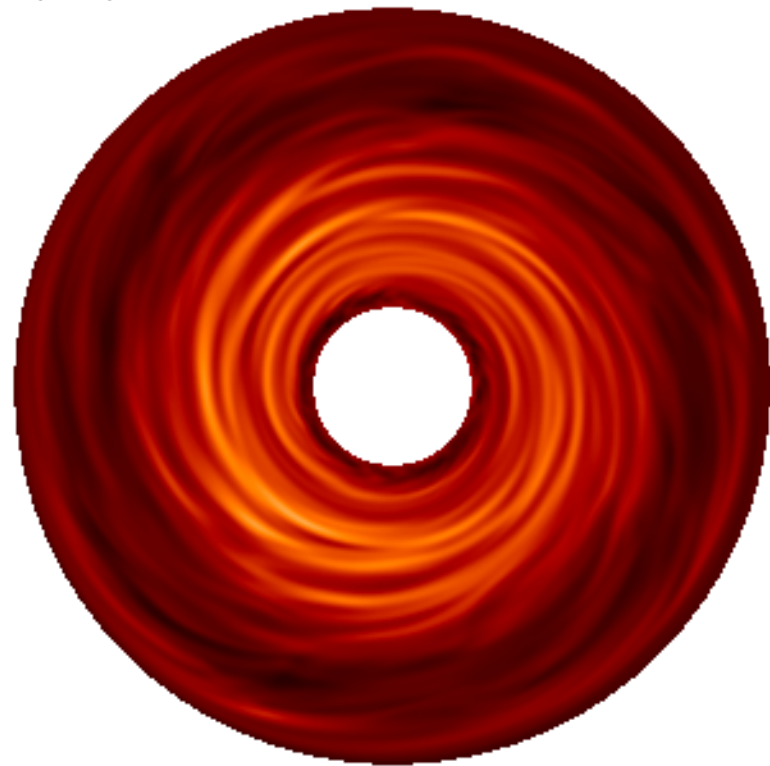
Debris Disks

Stabilization of architecture and Planet Detection

Accretion in disks occurs via turbulent viscosity

Turbulence in disks is enabled by
the Magneto-Rotational Instability

$t=46.3/88\text{yr}$



0.00 1.25 2.50

Magnetized disk

Lyra et al. (2008a)

MRI sketch

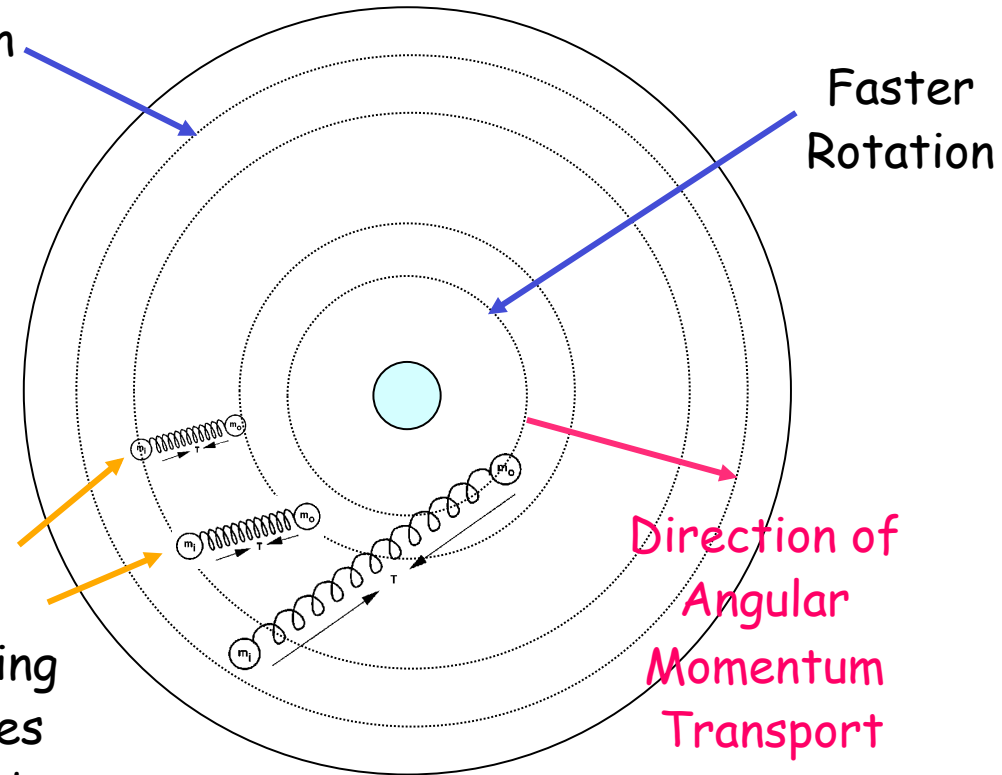
Slower
Rotation

Faster
Rotation

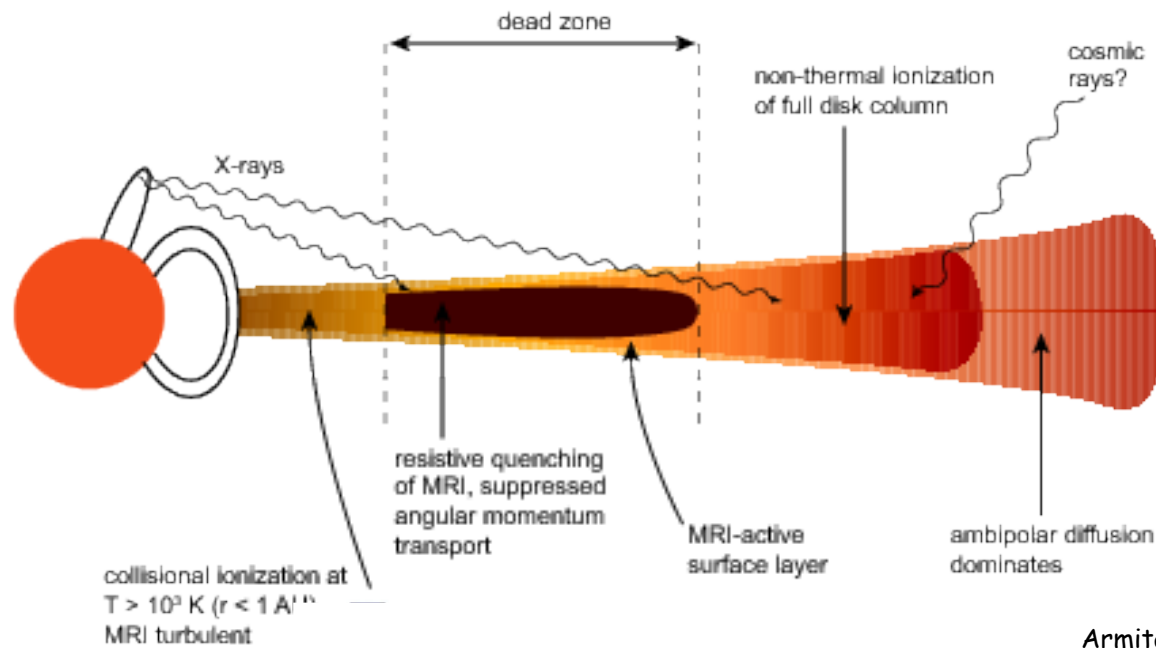
Stretching
amplifies
B-field

Direction of
Angular
Momentum
Transport

Unstable if angular
velocity decreases
outward



Alas... Dead zones are robust features of accretion disks

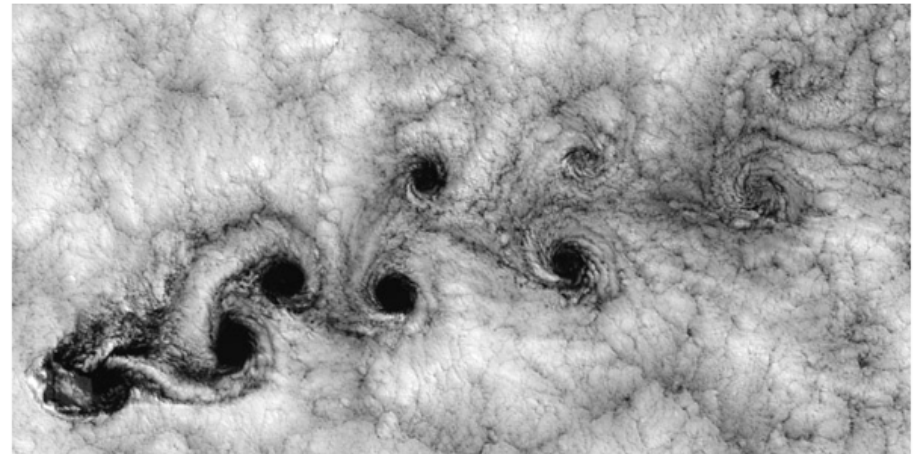
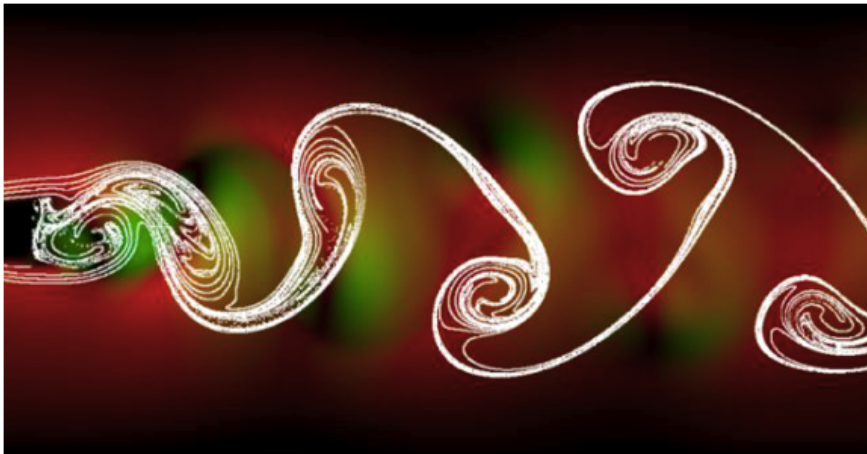


Disks are cold and thus poorly ionized
(Blaes & Balbus 1994)

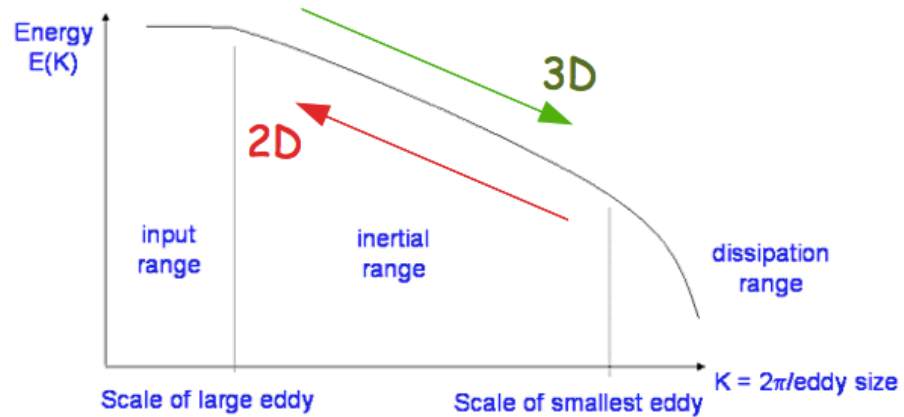
Therefore, accretion is **layered** (Gammie 1996)

There should be a non-magnetic,
hydrodynamical, source of turbulence in the **dead zone**.

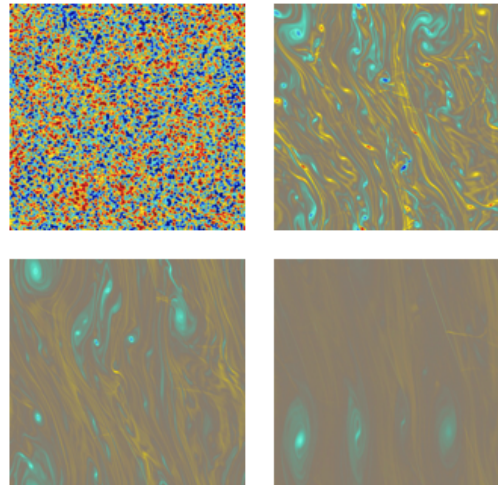
Vortices - An ubiquitous fluid mechanics phenomenon



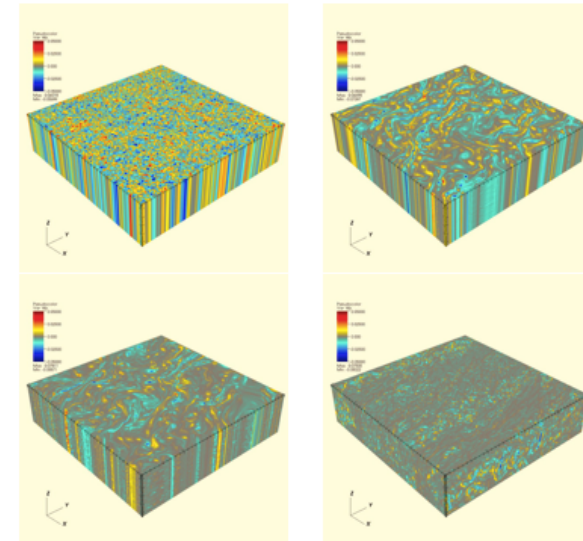
The energy cascade



Shen et al. (2006).
See also Batchelor (1967)



2D



3D

Inverse Cascade

No 3D instability
Eddies merge

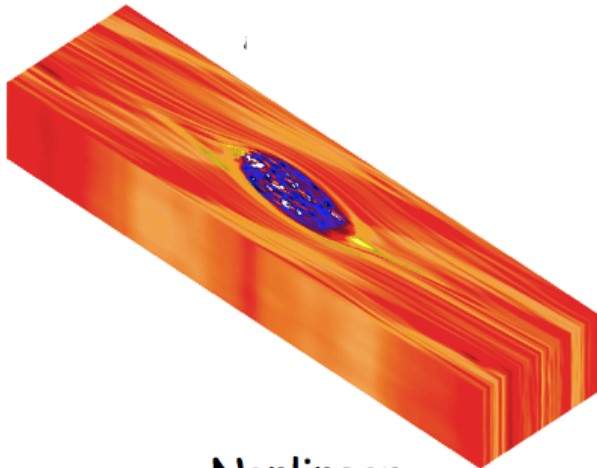
Direct Cascade

Destruction occurs
faster than merging

Sustaining vortices

Mechanisms to
inject vorticity
to counteract the vorticity lost in the direct cascade

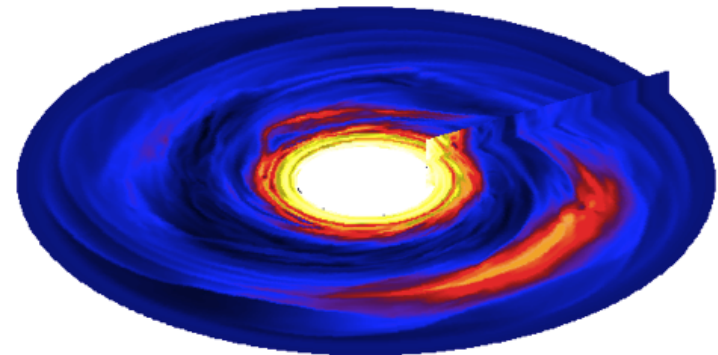
Baroclinic Instability



Nonlinear

Powered by:
buoyancy, thermal diffusion
(baroclinic source term)

Rossby Wave Instability

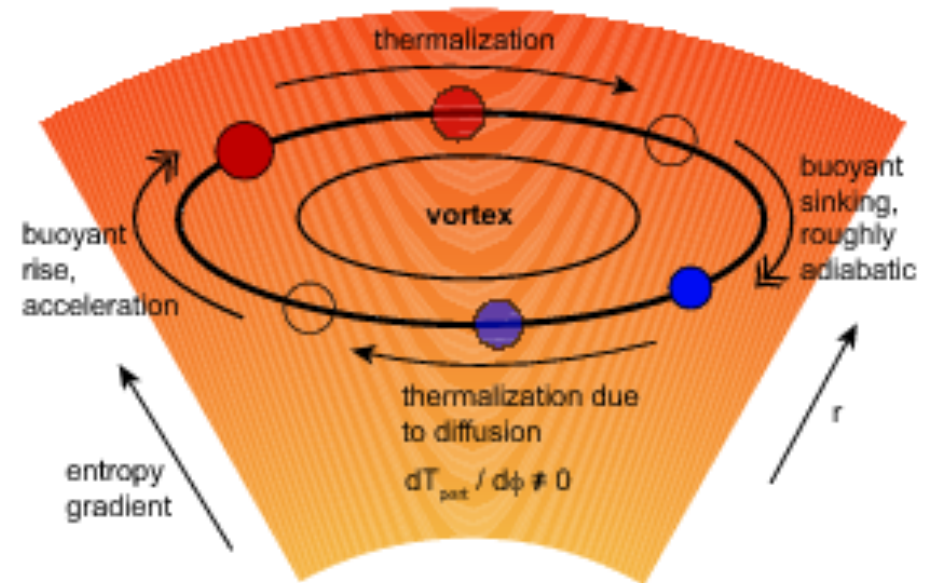


Linear

Powered by:
Modification of shear profile
(*external vortensity reservoir*)

Baroclinic Instability - Excitation and self-sustenance of vortices

Sketch of the
Baroclinic Instability



Armitage (2010)

$$\frac{\partial \omega}{\partial t} = \underbrace{-(\mathbf{u} \cdot \nabla) \omega}_{\text{advection}} - \underbrace{\omega (\nabla \cdot \mathbf{u})}_{\text{compression}} + \underbrace{(\omega \cdot \nabla) \mathbf{u}}_{\text{stretching}} + \frac{1}{\rho^2} \underbrace{\nabla \rho \times \nabla p}_{\text{baroclinicity}} + \underbrace{\nu \nabla^2 \omega}_{\text{dissipation}}$$

advection

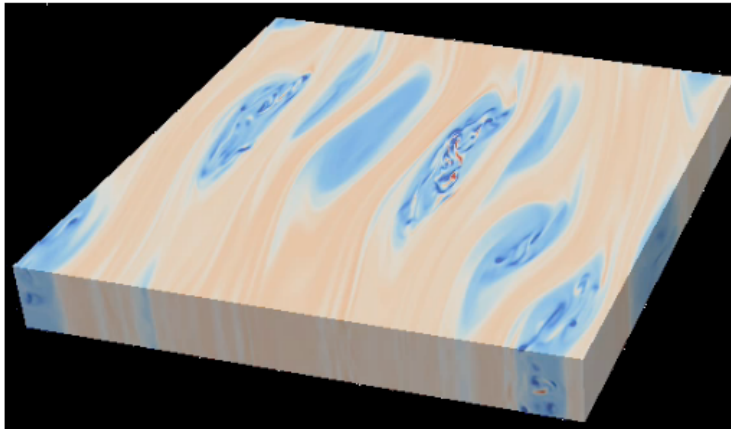
compression

stretching

baroclinicity

dissipation

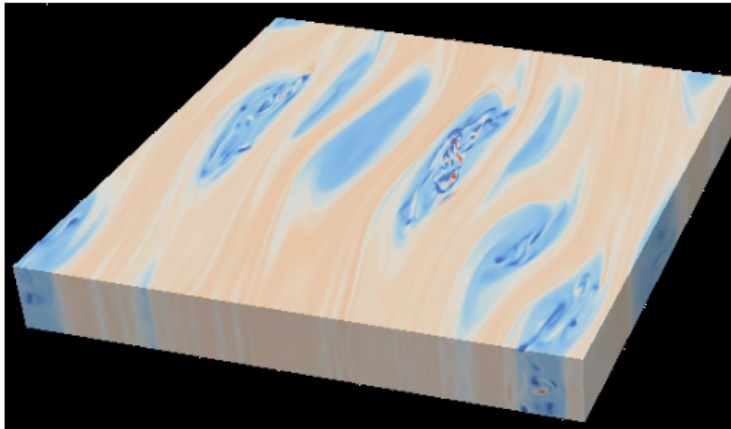
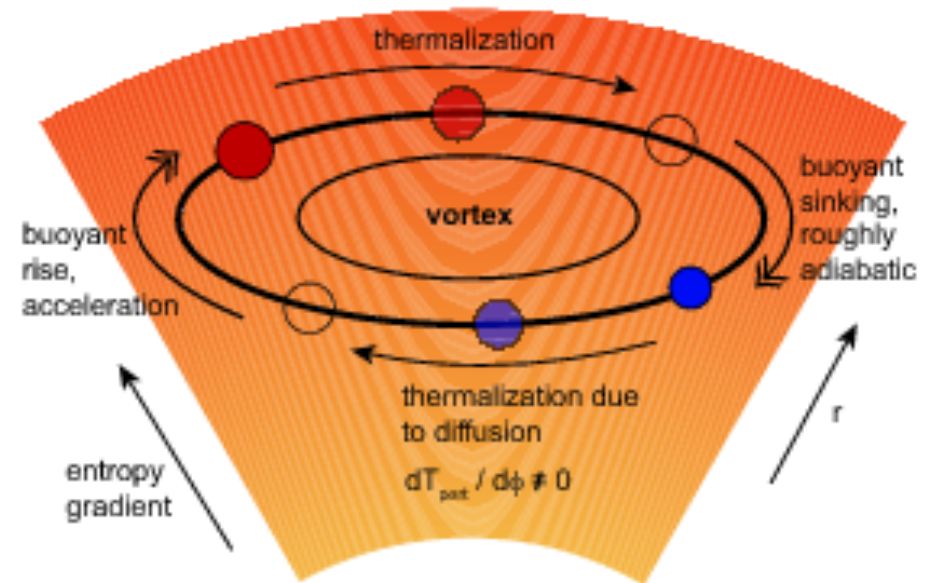
Lesur & Papaloizou (2010)



Baroclinic Instability - Excitation and self-sustenance of vortices

1. Radial entropy gradient
2. Thermal diffusion

Sketch of the Baroclinic Instability



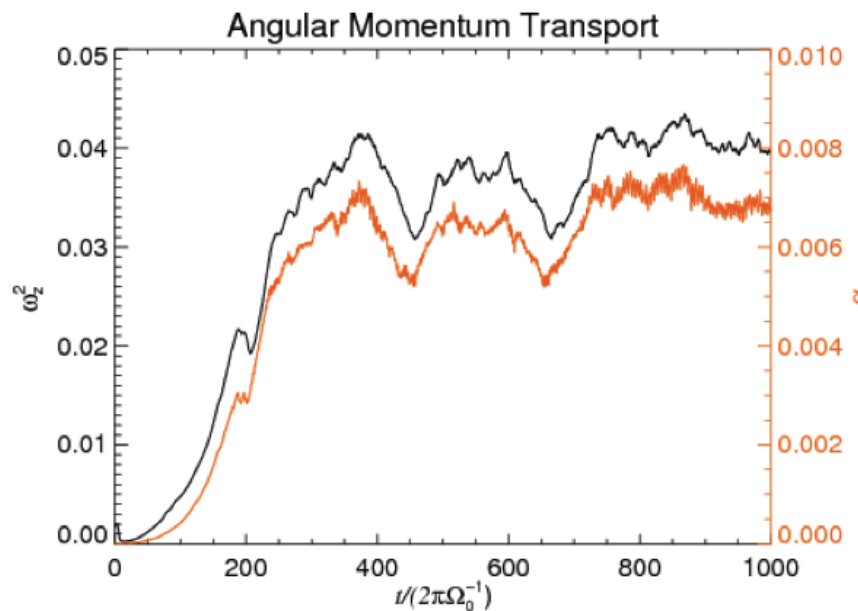
Lesur & Papaloizou (2010)

Armitage (2010)

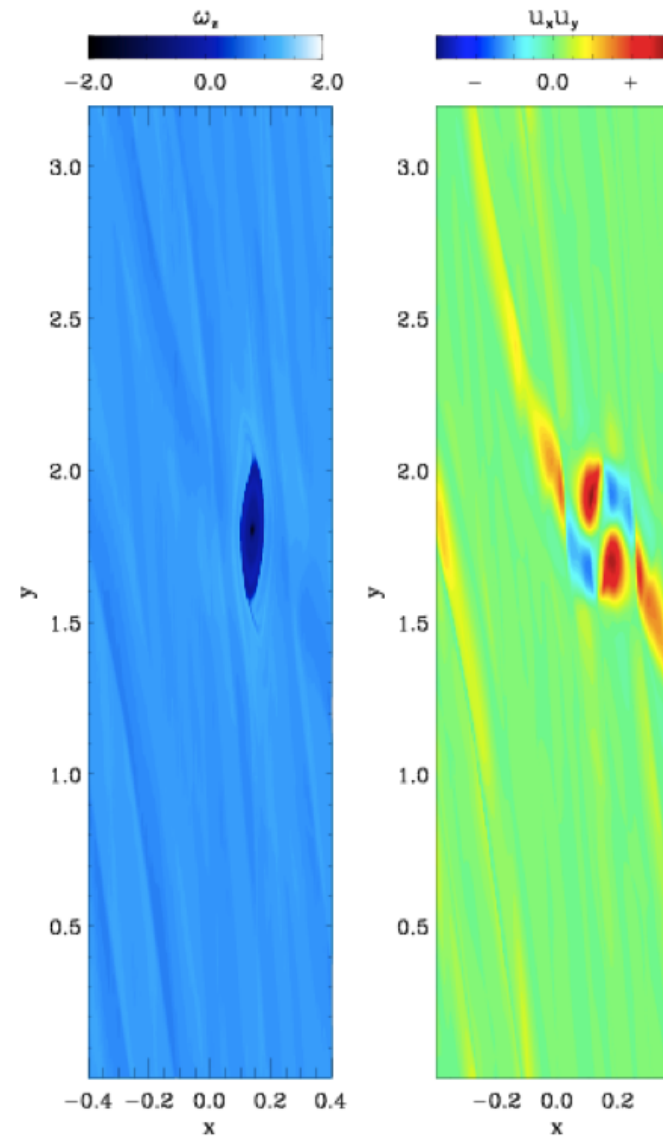
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Baroclinic Instability and Accretion

Large mass accretion rates,
comparable to the MRI!



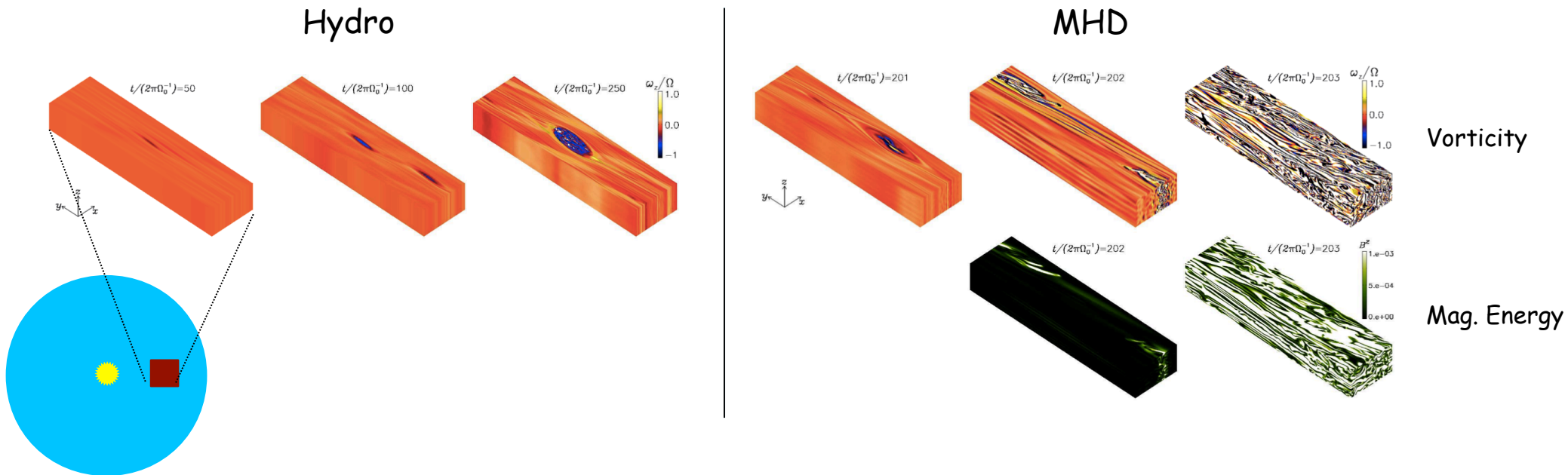
Raettig, Lyra, & Klahr (2012)



The angular momentum is carried by
waves excited by the vortex
(see also Heinemann & Papaloizou 2008,2009)

Baroclinic instability and layered accretion

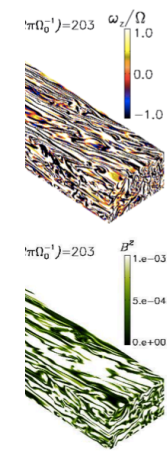
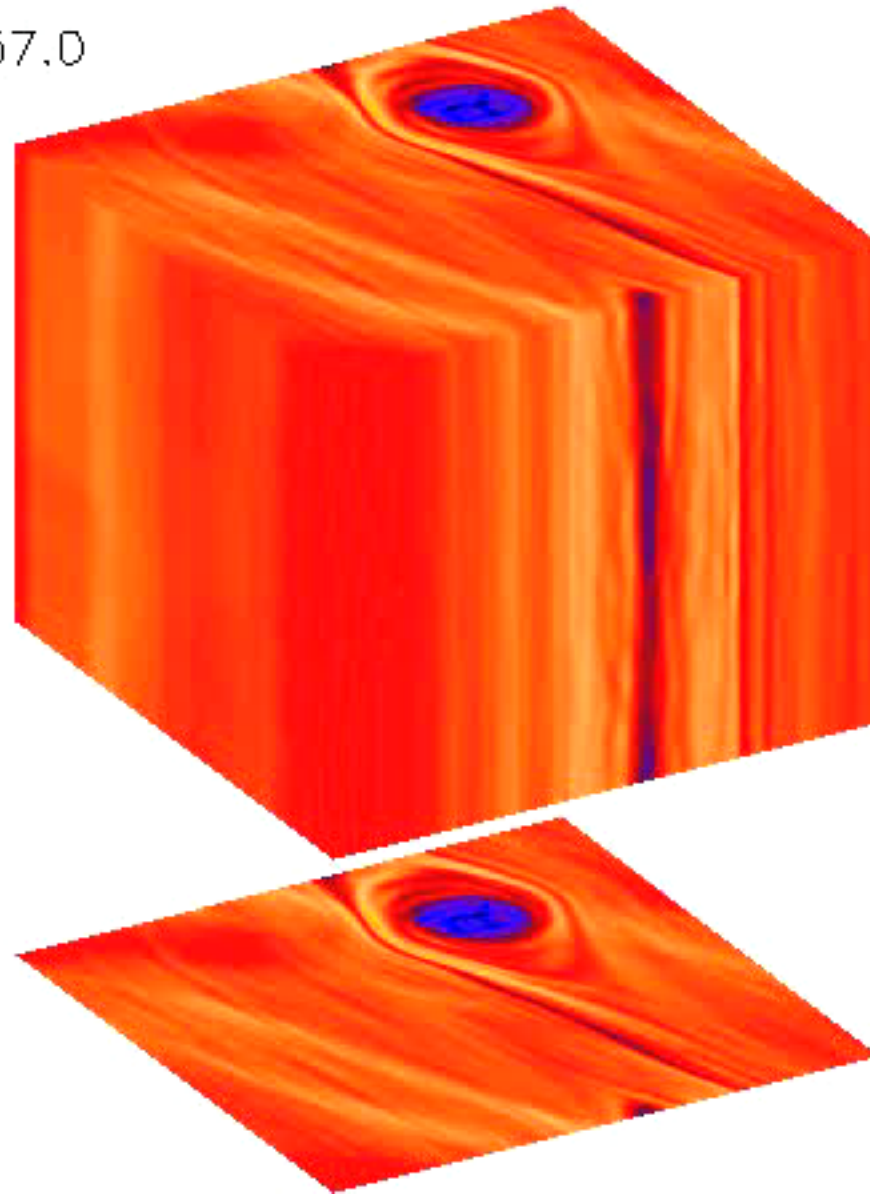
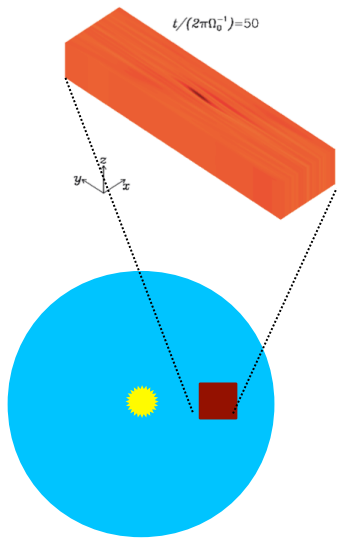
What happens when the vortex is magnetized?



Lyra & Klahr (2011)

Baroclinic instability and layered accretion

$t = 1257.0$

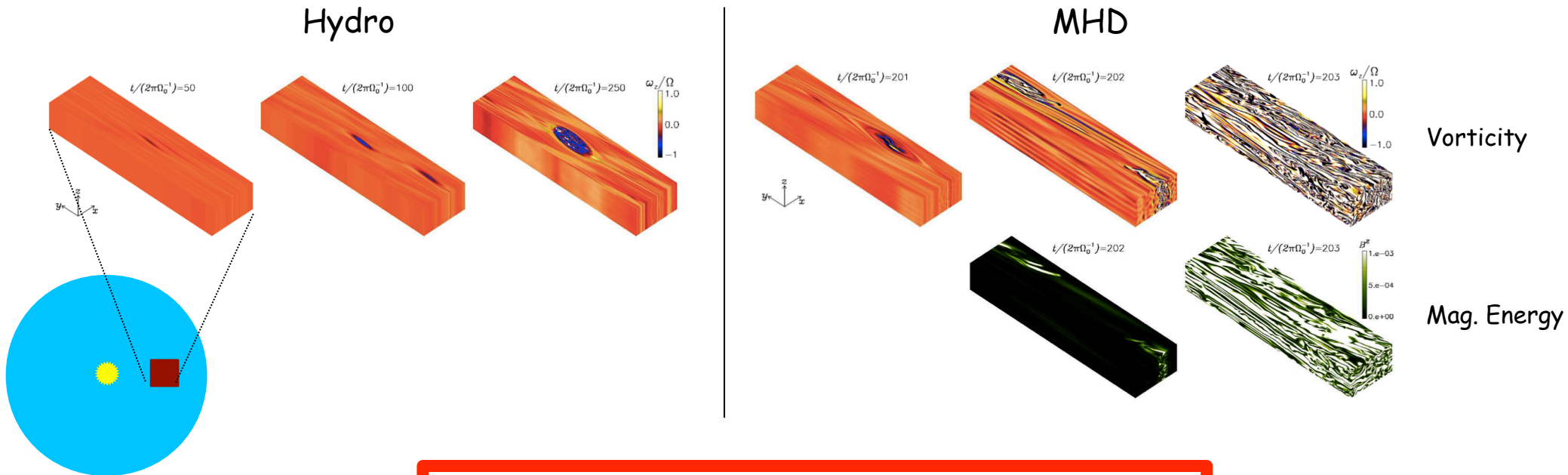


Vorticity

Mag. Energy

Baroclinic instability and layered accretion

What happens when the vortex is magnetized?

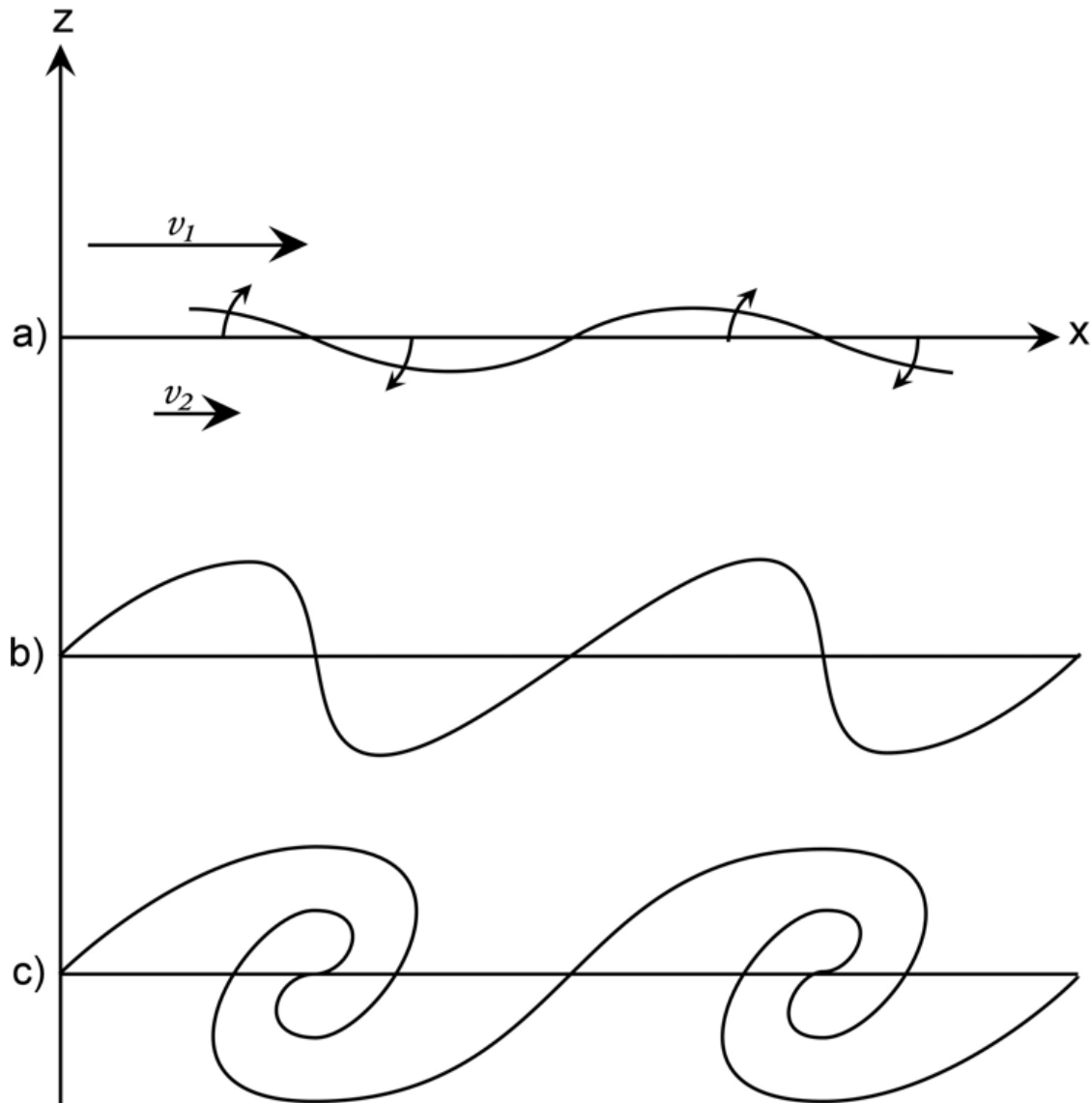


Baroclinic vortices
do *not* survive magnetization

Lyra & Klahr (2011)

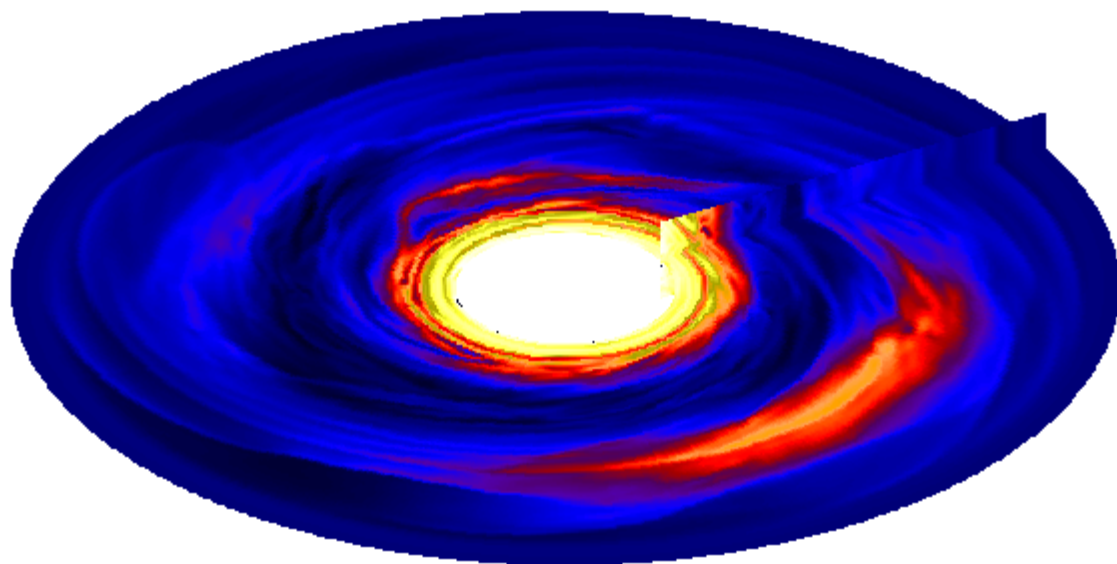
Rossby Wave Instability

(or.... Kelvin-Helmholtz in rotating disks)



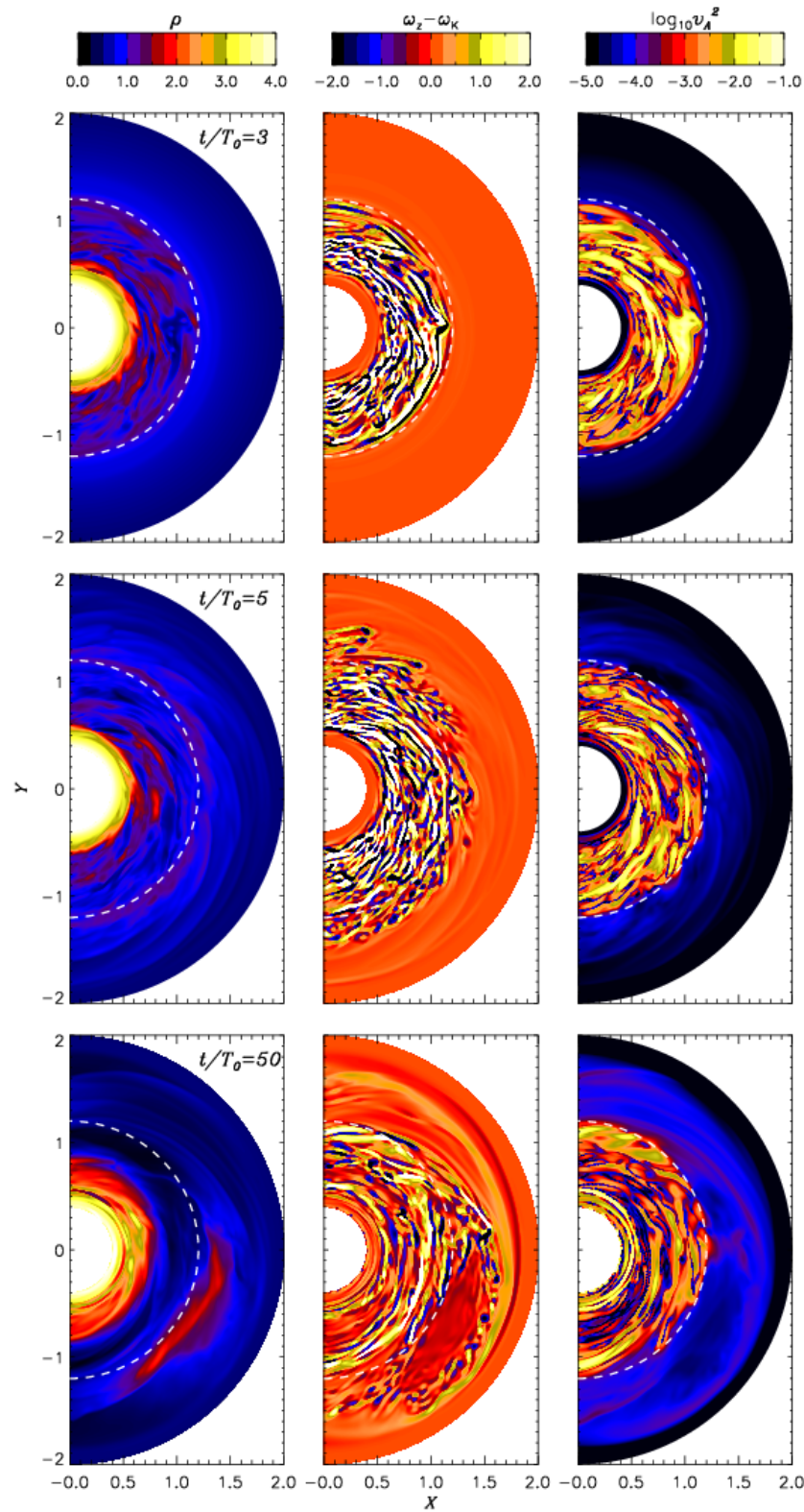
Active/dead zone boundary

$t = 22.28 \tau_0$



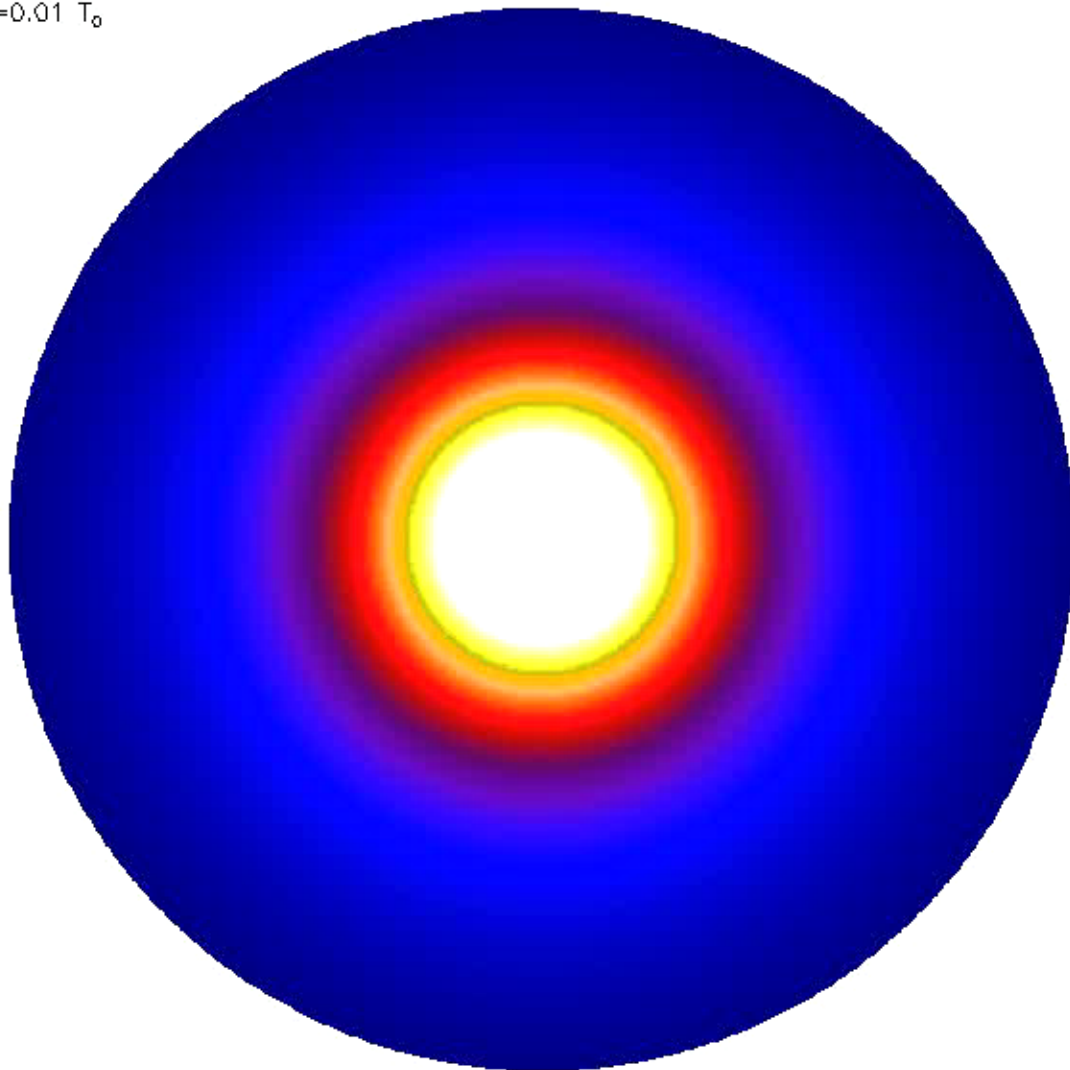
Magnetized inner disk + resistive outer disk

Lyra & Mac Low (2012)



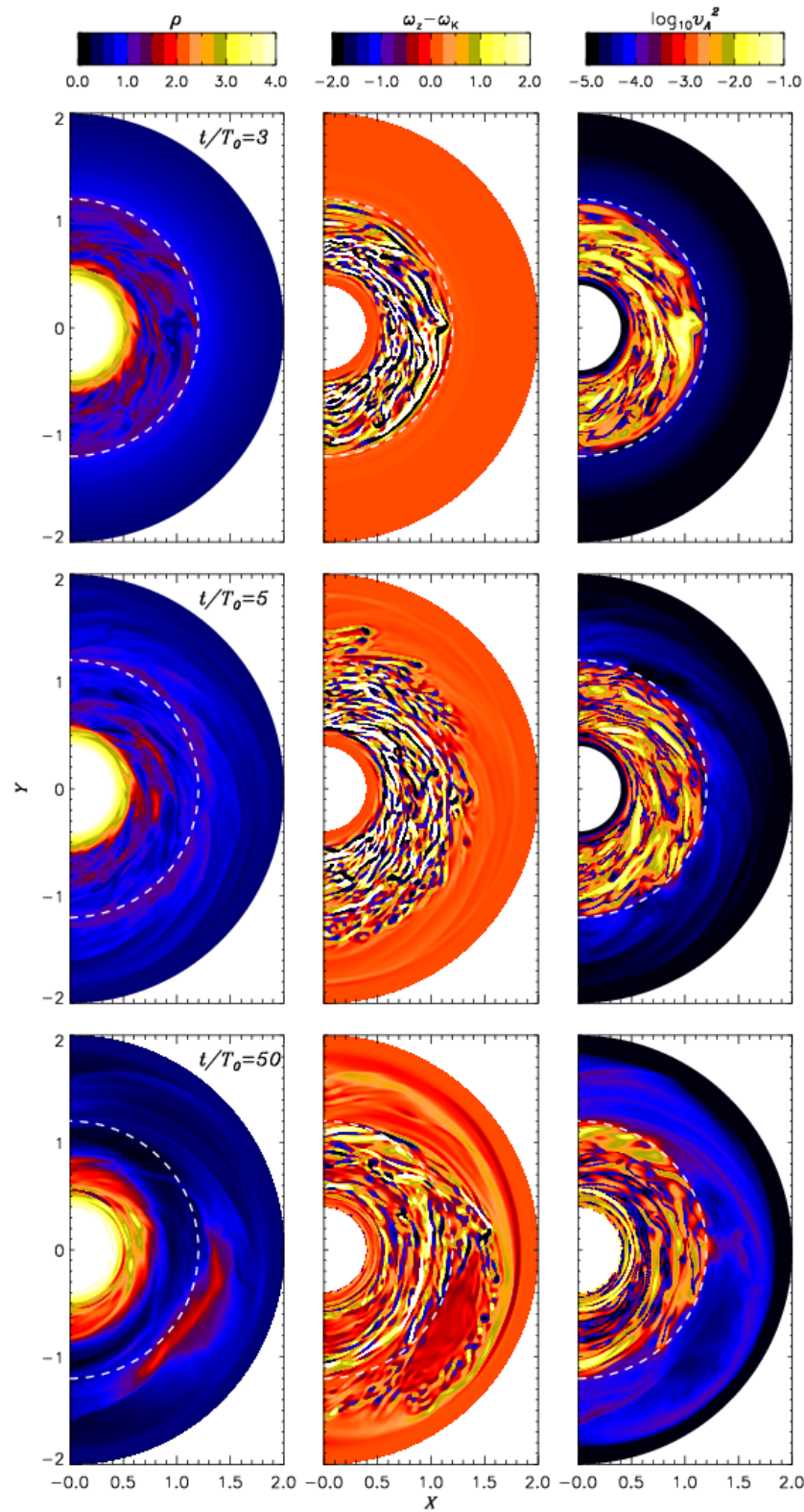
Active/dead zone boundary

$t = 0.01 T_0$



Magnetized inner disk + resistive outer disk

Lyra & Mac Low (2012)

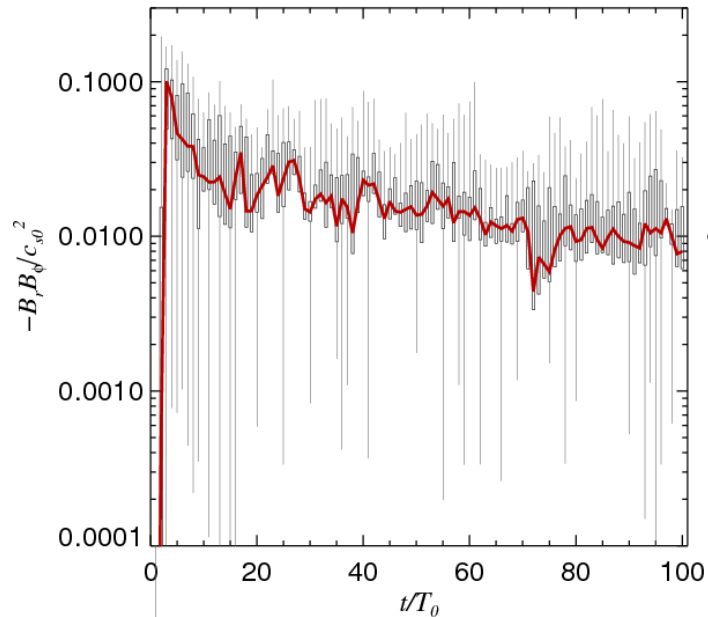


Significant angular momentum transport

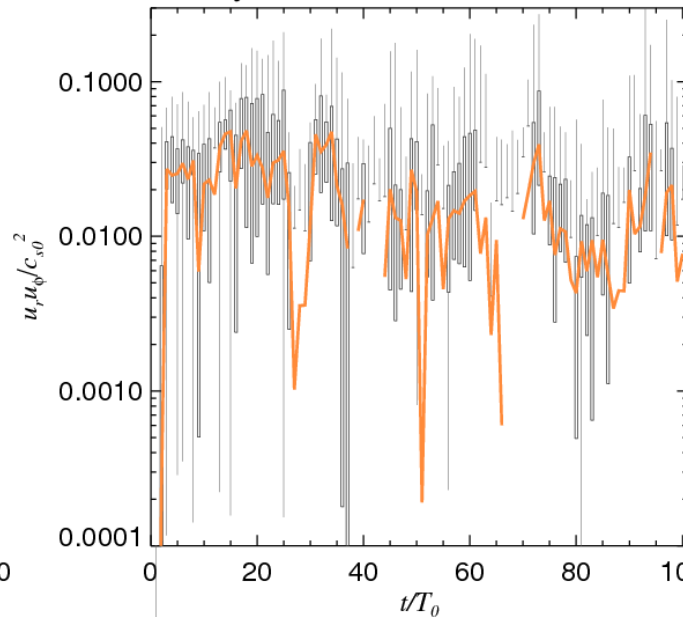
Active zone

Dead zone

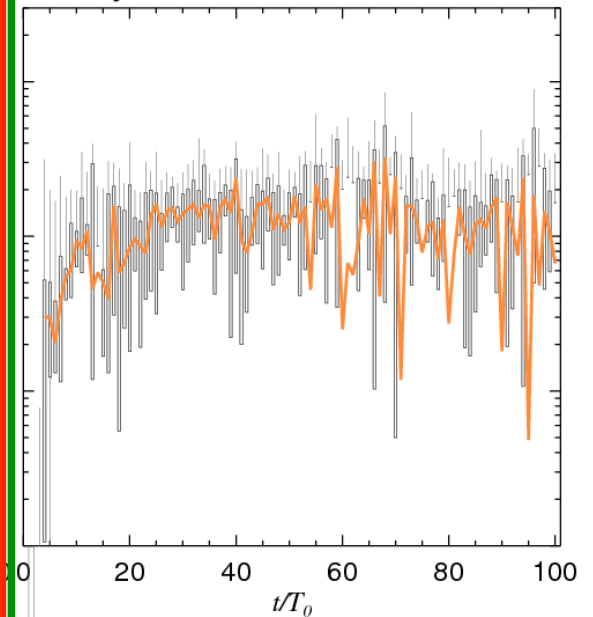
Maxwell stress – active zone



Reynolds stress – active zone



Reynolds stress – dead zone

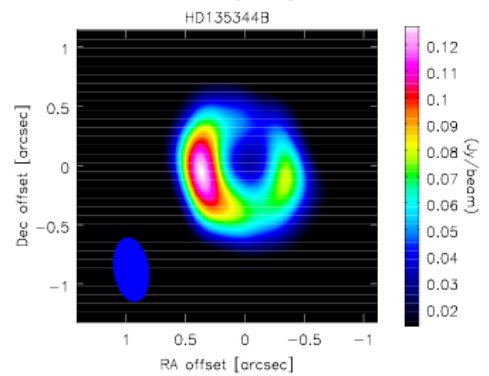
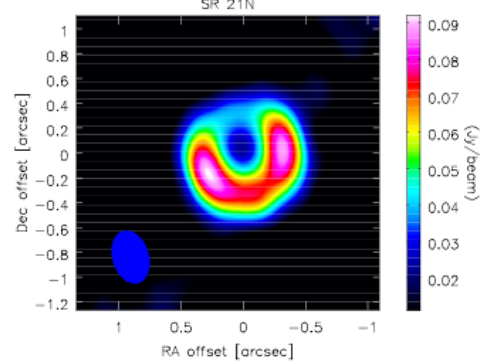
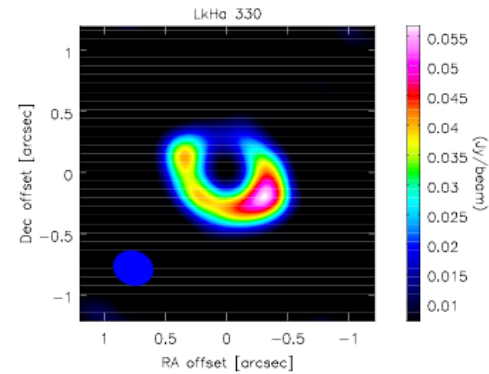
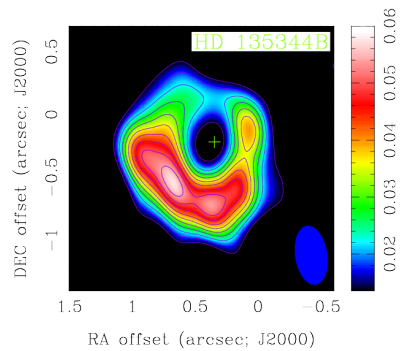
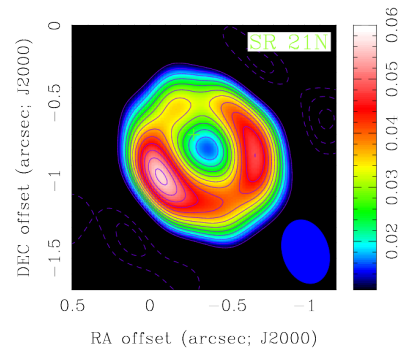
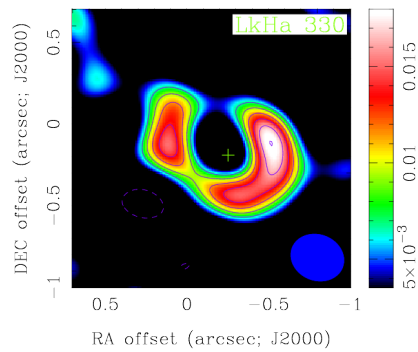


Large mass accretion rates in the **dead zone**,
comparable to the MRI in the **active zone**!

A possible detection of vortices in disks

Observations

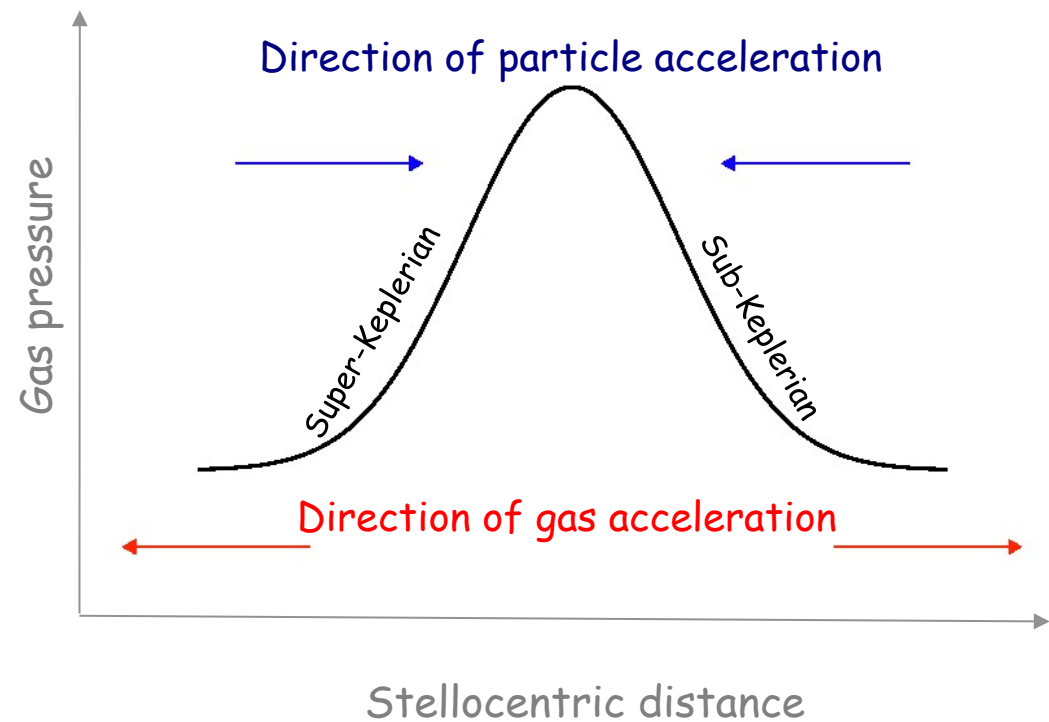
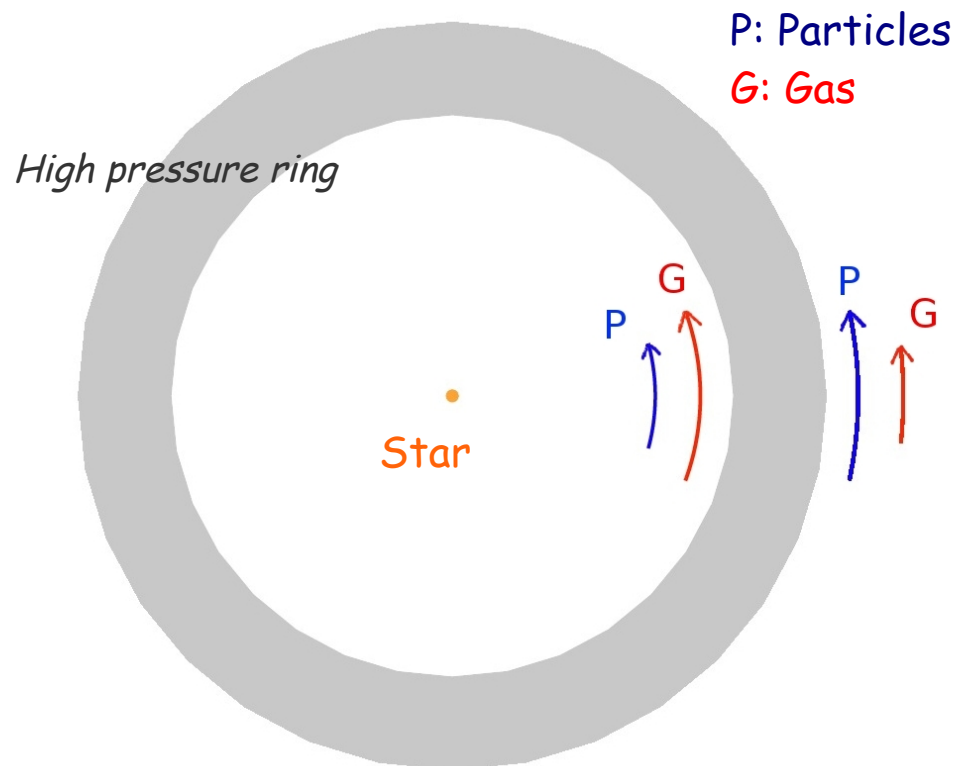
Brown et al. (2009)



Simulated observations
of Rossby vortices

Regaly et al. (2012)

Forming planets in turbulent disks



Forming planets in turbulent disks

Gas

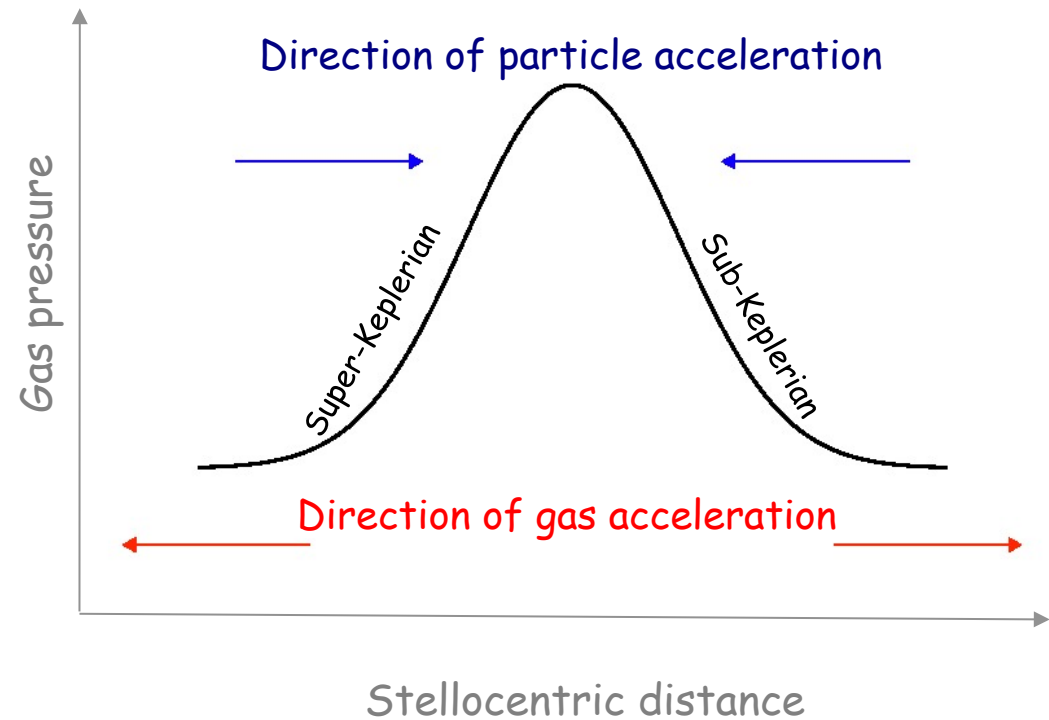
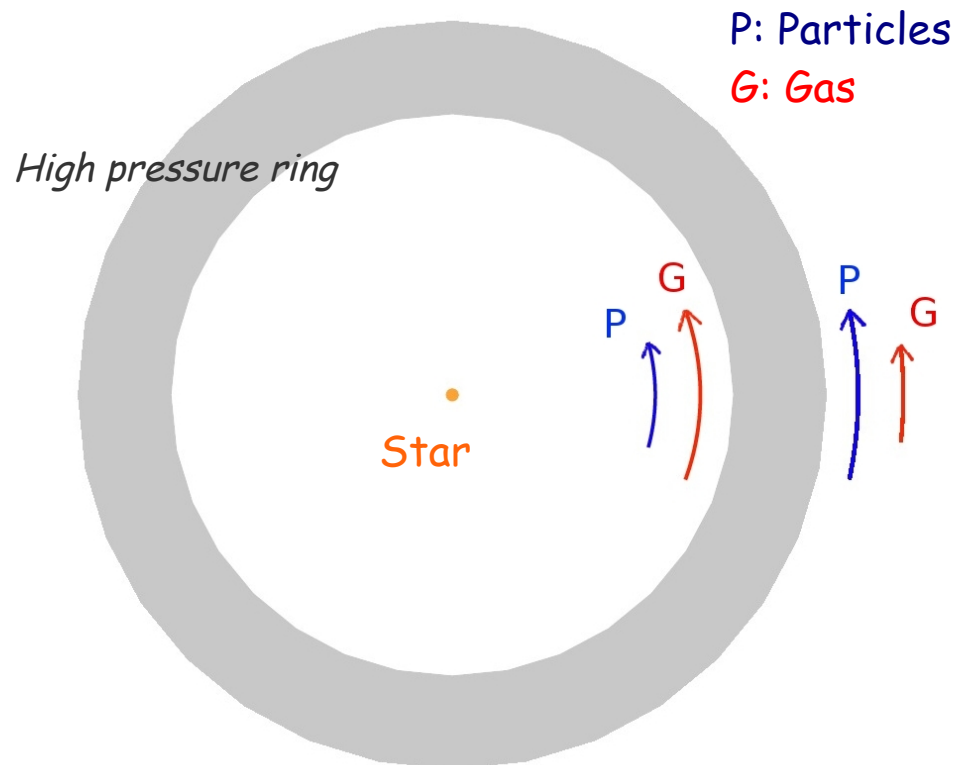
$$\frac{D \mathbf{u}}{Dt} = -\nabla \Phi - \rho^{-1} \nabla p$$

Particles

$$\frac{d \mathbf{w}}{dt} = -\nabla \Phi - \frac{(\mathbf{w} - \mathbf{u})}{\tau}$$

$$\mathbf{w} = \mathbf{u} + \tau \rho^{-1} \nabla p$$

The drag force pushes the particles *toward* the pressure gradient

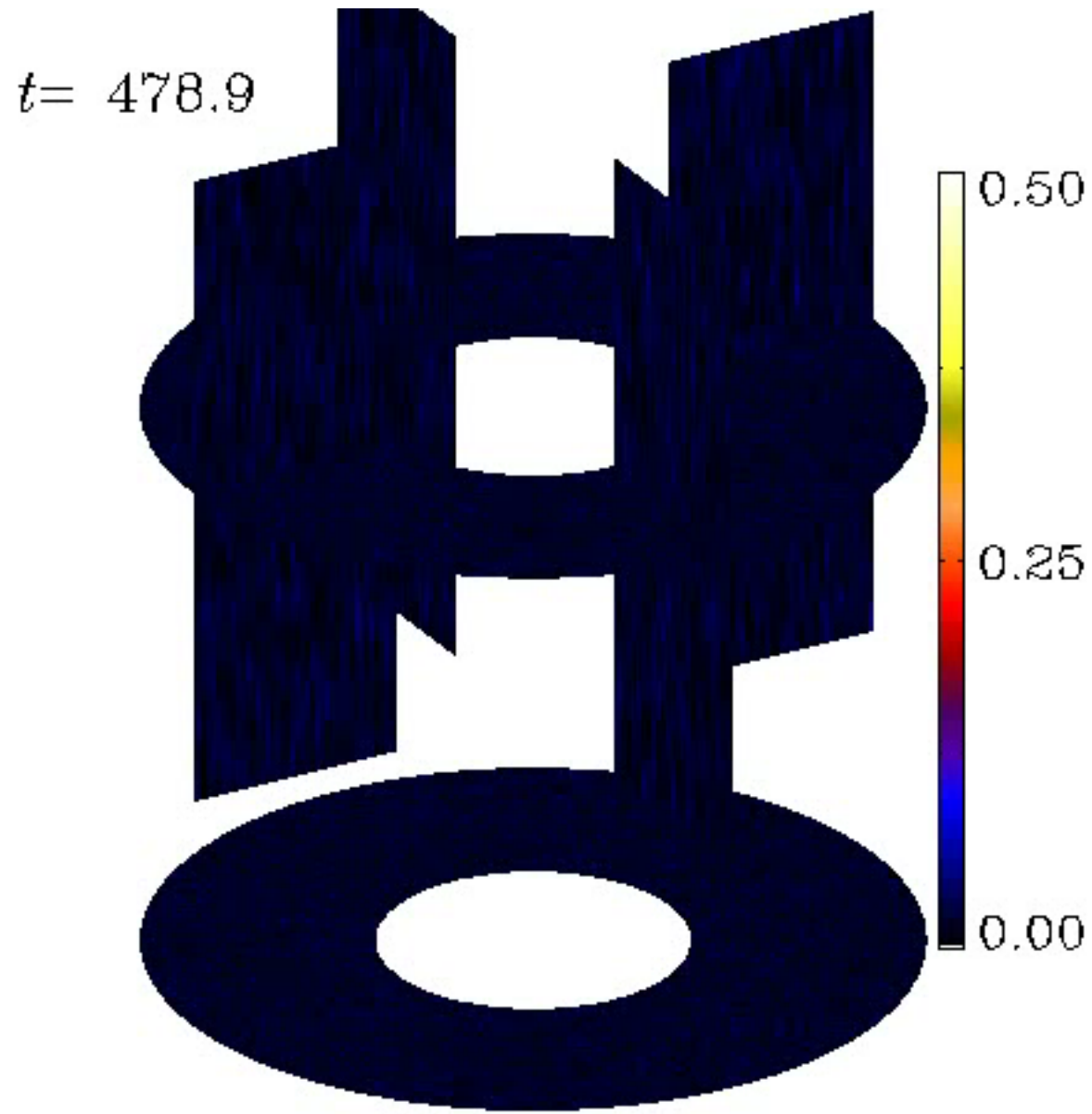


Solid particles

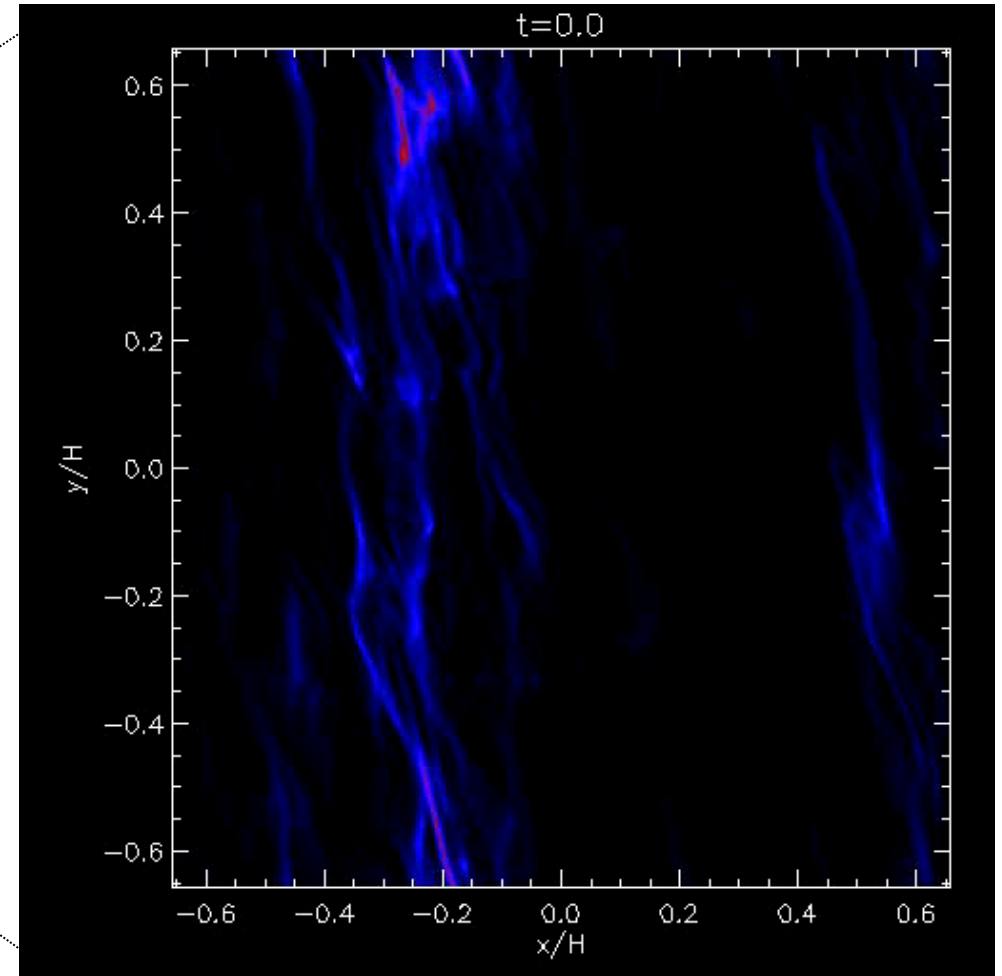
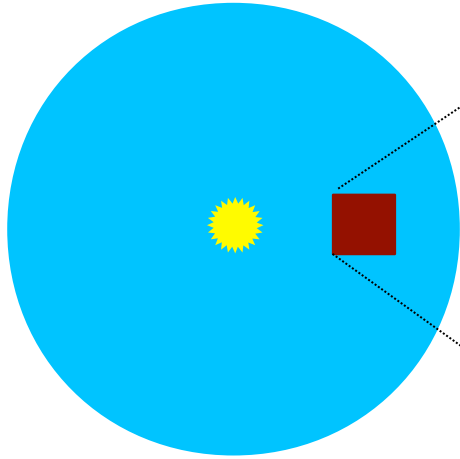
move toward

pressure maxima

Turbulence concentrates solids mechanically in pressure maxima



Gravitational collapse into planetesimals



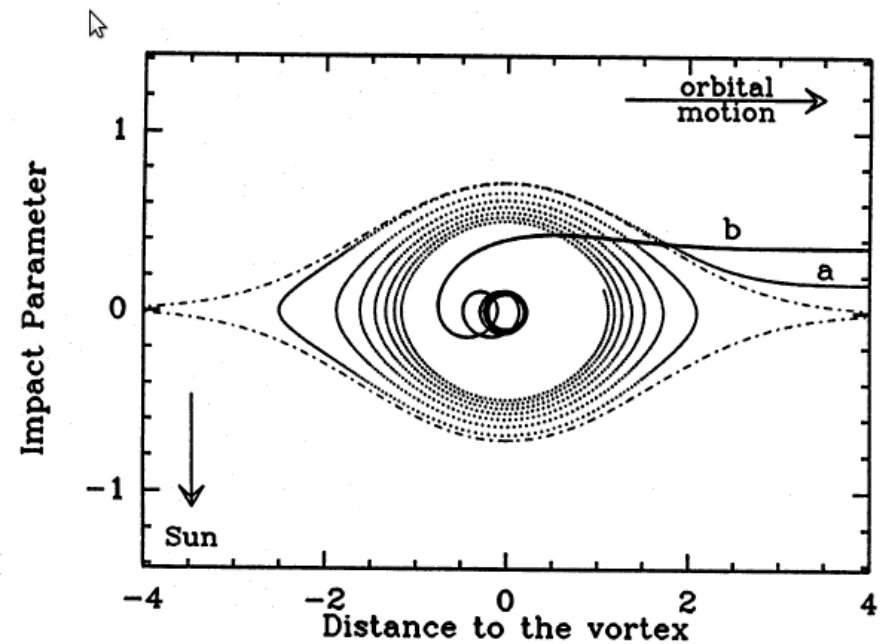
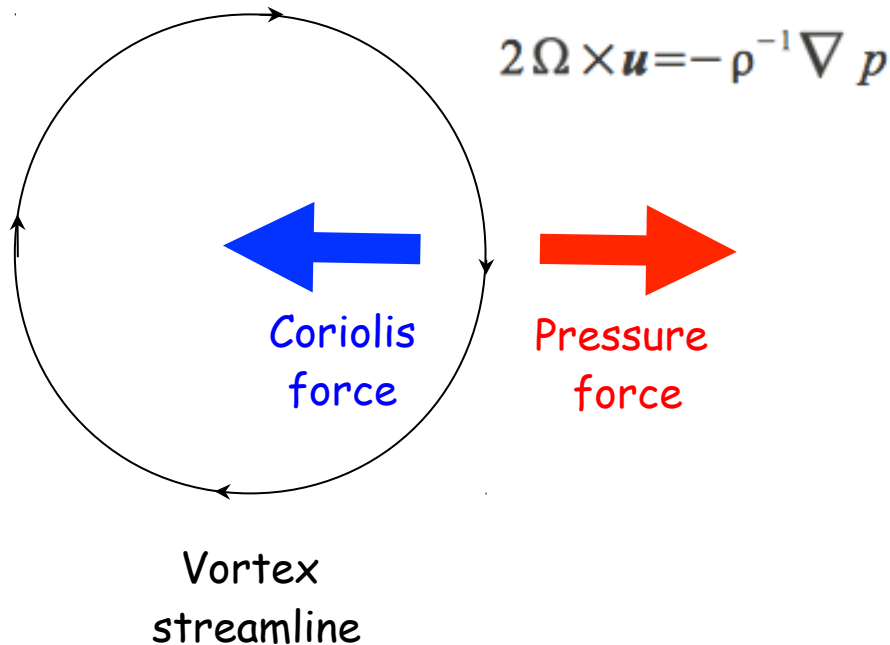
Johansen et al. (2007)

Turbulent eddies concentrate solids,
turning them into planetesimals...

...and vortices are **huge** eddies!

The Tea-Leaf effect

Geostrophic balance:



Barge & Sommeria (1995)

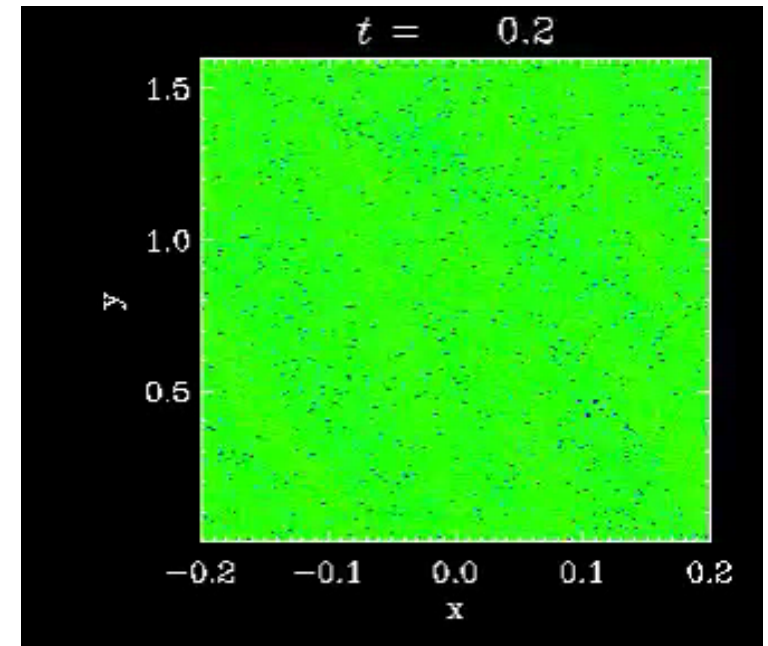
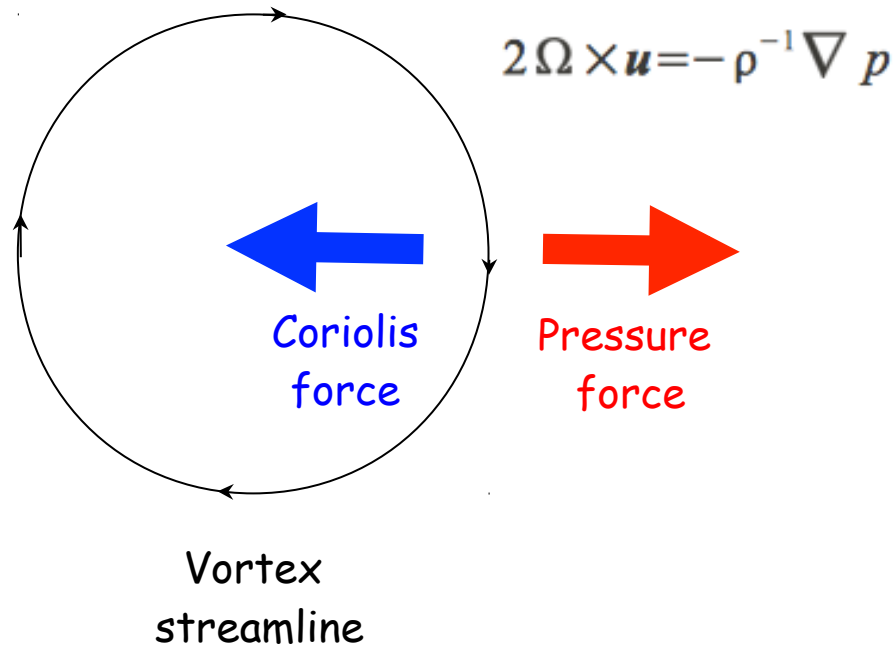
Particles do not feel the pressure gradient.
They sink towards the center, where they accumulate.

Aid to planet formation (Barge & Sommeria 1995)

Speed up planet formation enormously
(Lyra et al. 2008b, 2009a, 2009b, Raettig, Lyra & Klahr 2012)

The Tea-Leaf effect

Geostrophic balance:



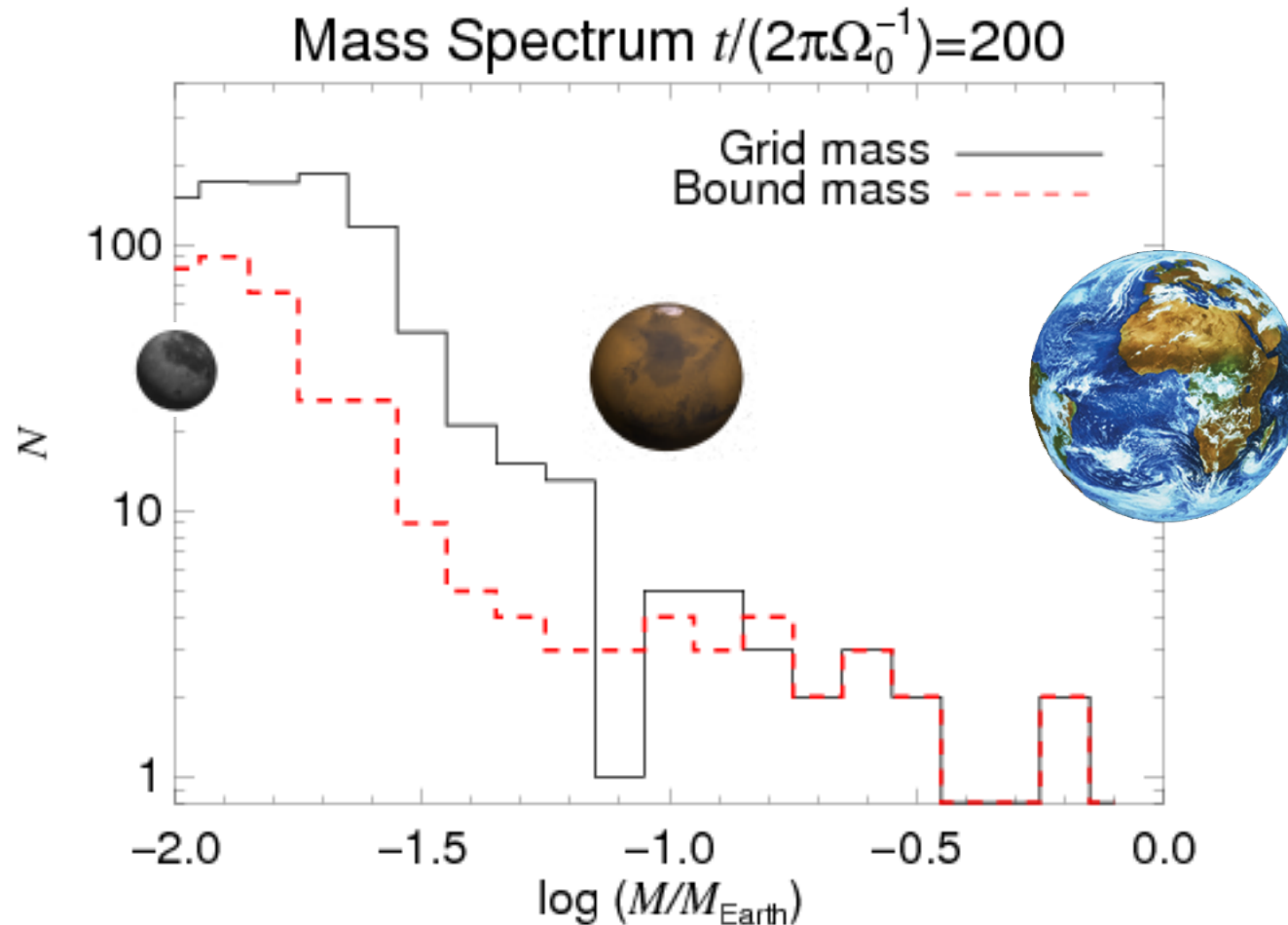
Raettig et al. (2012)

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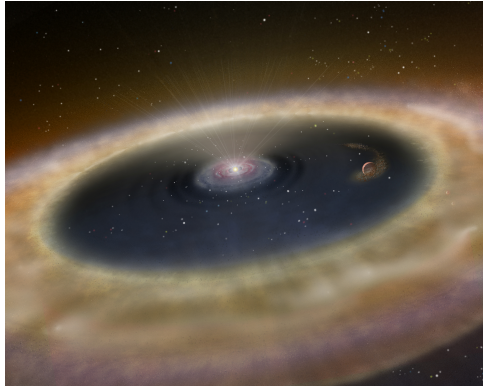
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The Initial Mass Function of planets

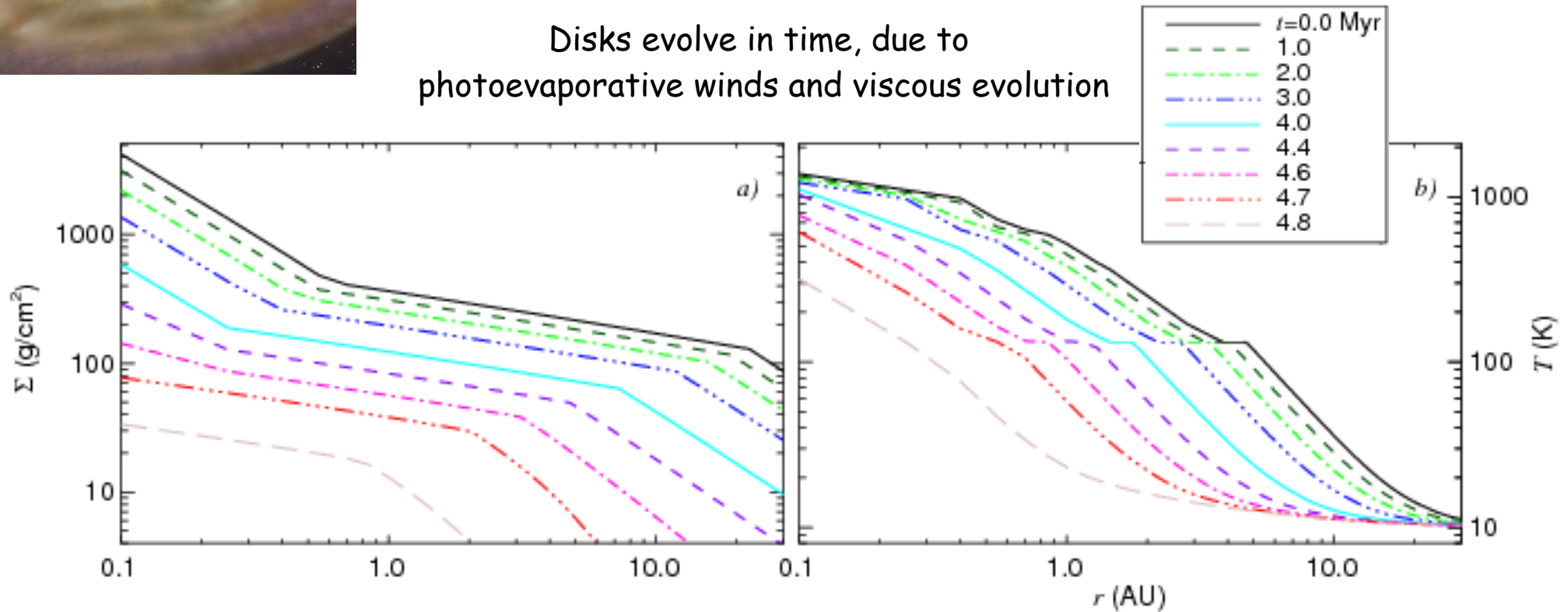


- Mass spectrum by the end of the simulation
 - 300 bound clumps were formed
- Power law $d(\log N)/d(\log M) = -2.3 \pm 0.2$
- 20 of these are more massive than Mars

Transitional disks - The thinning phase



Disks evolve in time, due to
photoevaporative winds and viscous evolution

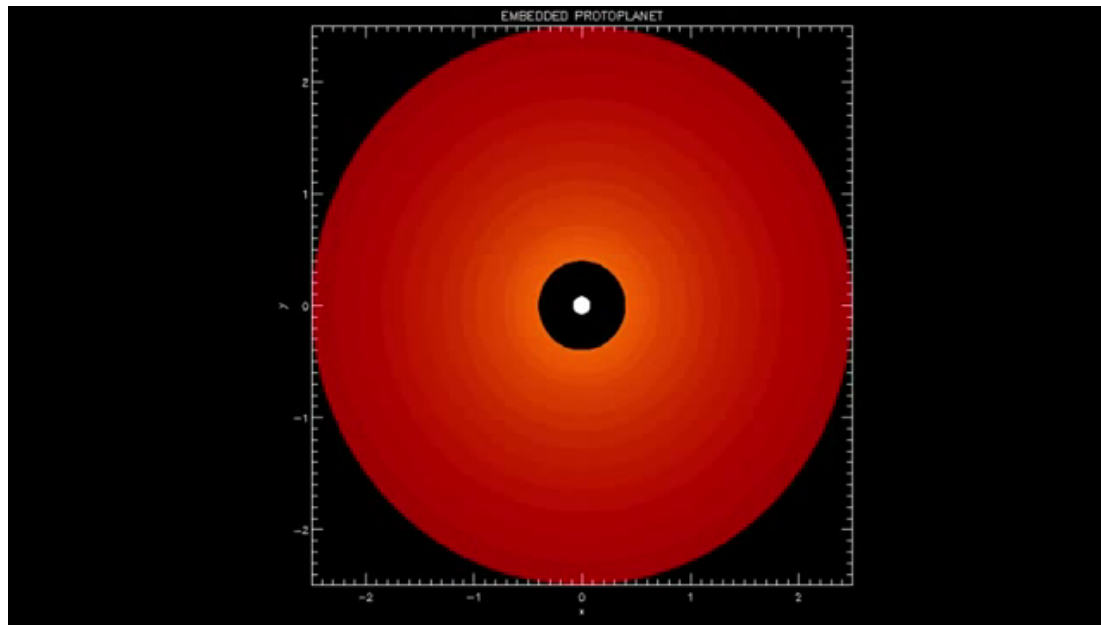


Lyra, Paardekooper, & Mac Low (2010)

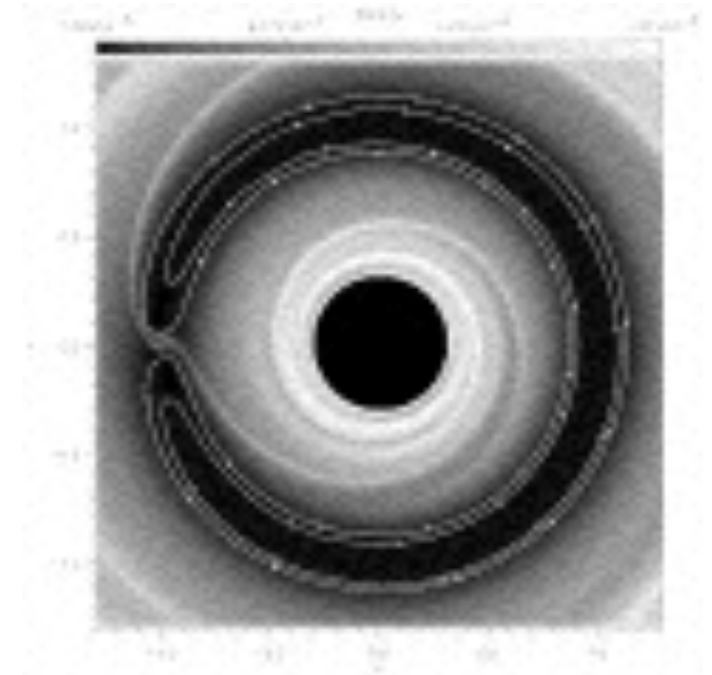
Planets form and start to migrate

Planet-disk interaction leads to **angular momentum exchange**

One armed spiral: Lindblad resonance
Horseshoe libration: Co-rotational torques



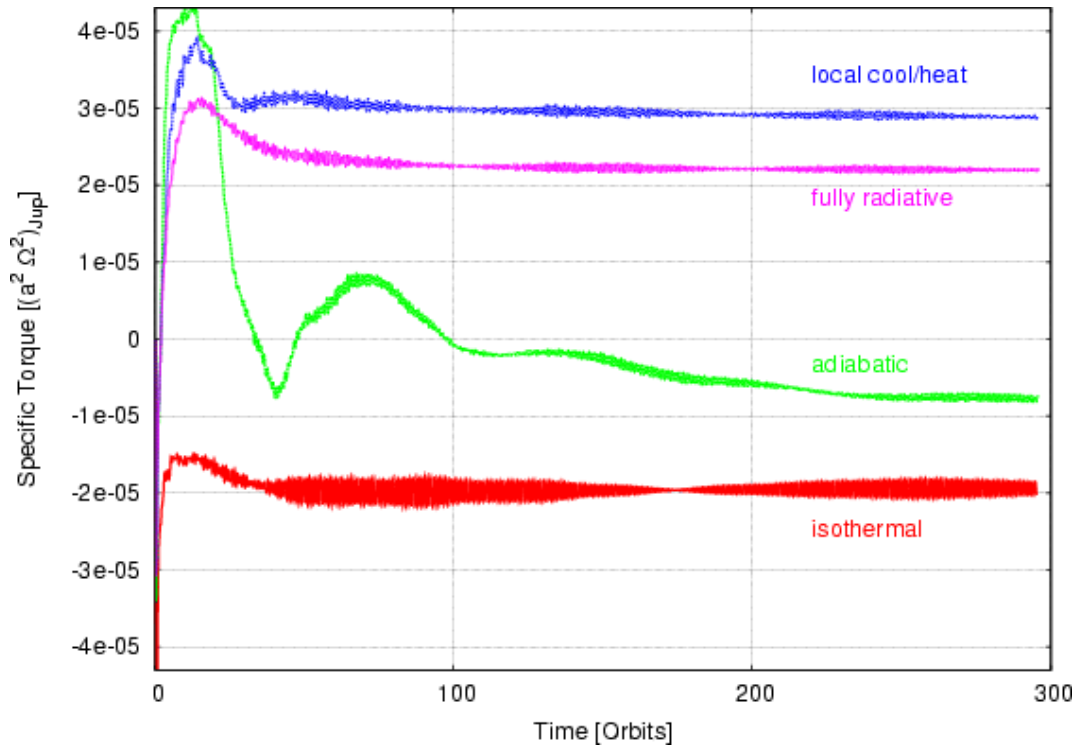
Kley & Nelson (2012, A&A Annual Review)



Lubow et al. (1999)

In **isothermal** disks,
the result is **inward** migration

Planets form and start to migrate



Kley & Crida (2008)

Rule of thumb: Migration is
outwards in
steep temperature gradients,
inwards in
isothermal regions.

Paardekooper & Mellema (2006)

Non-isothermal
co-rotational torque may lead
to outward migration

Hot topic!

Paardekooper & Mellema 2008

Baruteau & Masset 2008

Paardekooper & Papaloizou 2008

Kley & Crida 2009

Kley et al 2009

Paardekooper et al. 2010

Bitsch & Kley 2010

Lyra et al. 2010

Paardekooper et al. 2011

Ayliffe & Bate 2011

Yamada & Inaba 2011

Kley 2011

Bitsch et al. 2012

Nelson & Kley 2012

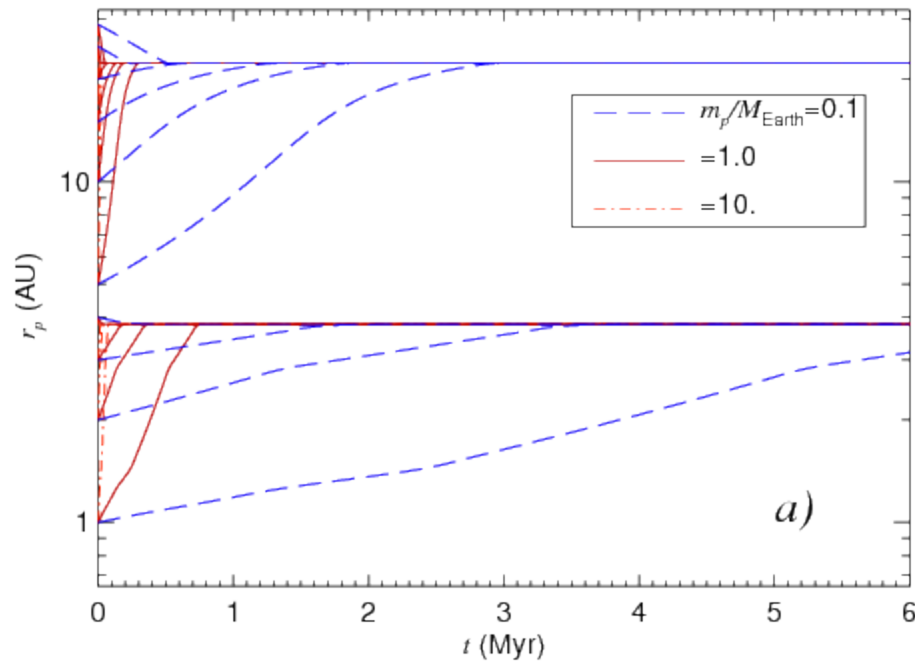
Pierens et al. 2012

Bitsch et al. 2013

Zhu et al. 2013

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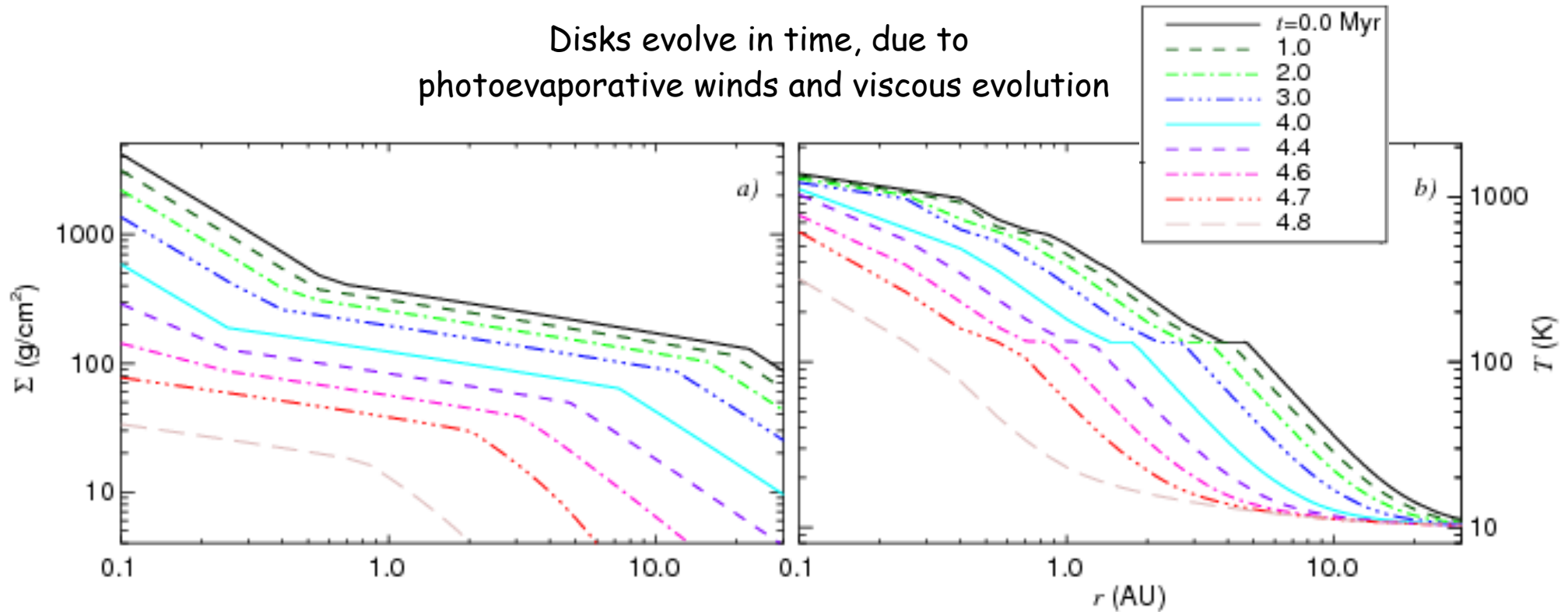


Lyra, Paardekooper, & Mac Low (2010)

Planet traps where migration
is **convergent**
($\tau=0$, $d\tau/dr < 0$).

Migration in Evolutionary Models

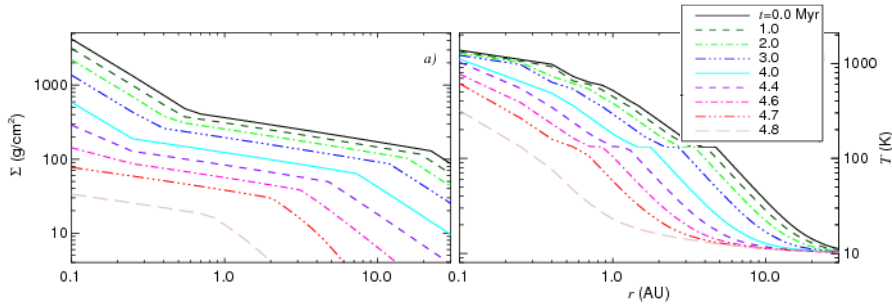
Disks evolve in time, due to
photoevaporative winds and viscous evolution



Lyra, Paardekooper, & Mac Low (2010)

Migration in Evolutionary Models

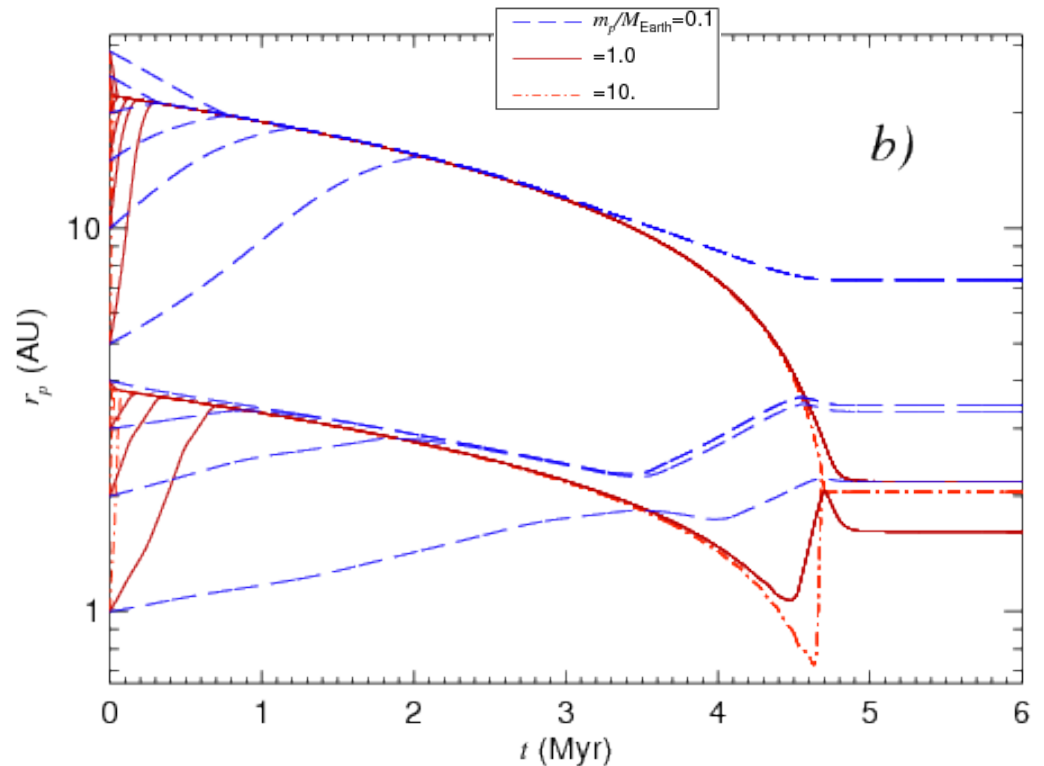
Disks evolve in time, due to
photoevaporative winds and viscous evolution



Single planets in a planetary trap
evolve in **lockstep with the gas** at the
accretion timescale.

At some point, the **disk becomes too thin**
to drive accretion. The **planet decouples**
and is **released** in a safe orbit.

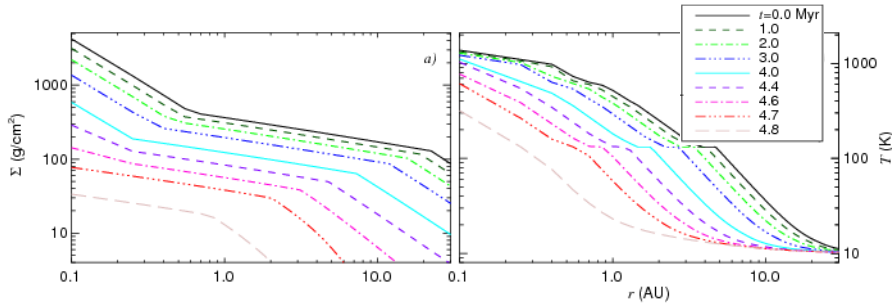
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Lyra, Paardekooper, & Mac Low (2010)

Migration in Evolutionary Models

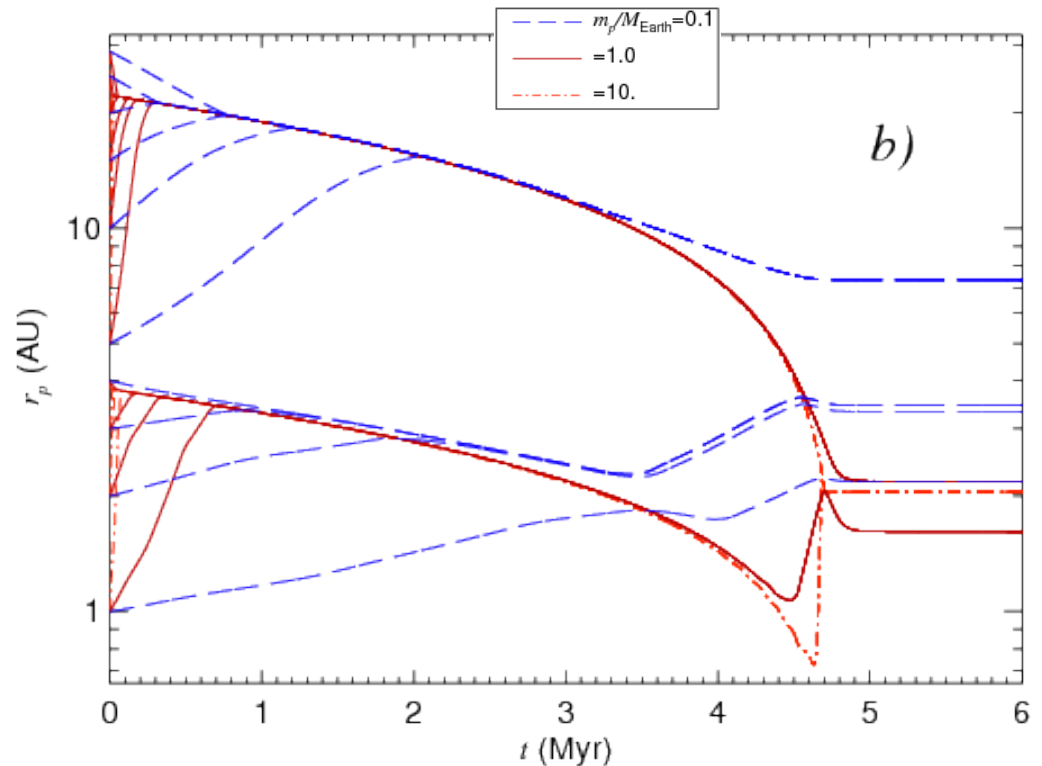
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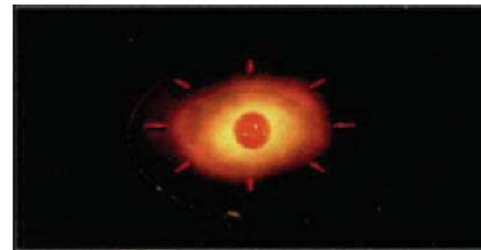
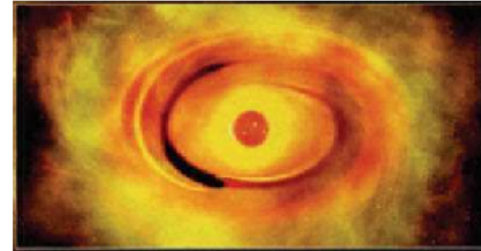
Lyra, Paardekooper, & Mac Low (2010)

Migration in Evolutionary Models

Single planets in a planetary trap evolve in lockstep with the gas at the accretion timescale.

At some point, the disk becomes too thin to drive accretion.

The planet decouples and is released in a safe orbit.



"La Terre sauvée"



Planètes

On sait pourquoi elles survivent à leur étoile

REPÈRES Jusqu'à récemment, la naissance du système solaire – et de tous les systèmes planétaires – posait un problème insurmontable : en effet, d'après les modélisations informatiques, les planètes auraient dû être précipitées vers le Soleil avant même d'avoir atteint leur taille définitive, il y a 4,6 milliards d'années. Mais un nouveau modèle semble résoudre définitivement ce paradoxe.

Par Román Ikonoff

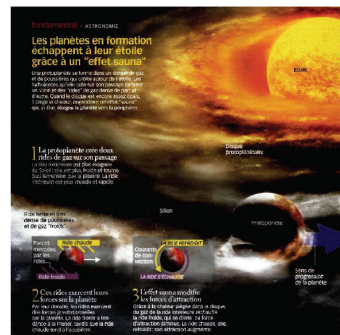
"La Terre sauvée". Le titre de la conférence donnée en janvier 2010 par le Brésilien Wladimir Lyra, le Néerlandais Sjoen Jas Pasado-Kooper et l'Américain Michael Mac Low, lors du 215^e meeting de la Société astronomique américaine (AAS), était un bon malice.

LA TERRE NE DEVRAIT PAS EXISTER

Les trois chercheurs annonçaient ni plus ni moins avoir sauvé la Terre – et toutes les autres planètes du système solaire – d'une chute inéluctable sur le Soleil. Date prévue de ce cataclysme : 4,6 milliards d'années... en arrière ! Autrement dit, notre planète bleue aurait échappé à une catastrophe qui

n'a tout bonnement pas eu lieu – et nous sommes là pour en attester ! Au vrai, ce n'est donc pas la Terre que les trois scientifiques ont sauvée de la chute fatale... mais la communauté astronomique. Car il faut savoir que depuis une vingtaine d'années, tous les modèles informatiques simulant la naissance du système solaire aboutissaient au même scénario catastrophe : toutes les planètes étaient précipitées dans la fournaise solaire bien avant d'atteindre l'âge de raison. Coucherait Mars, Vénus, Saturne ou la Terre ne devraient pas exister. Pas plus que les "exoplanètes", ces centaines de planètes lointaines que les télescopes et satellites ont découvertes autour d'autres étoiles →

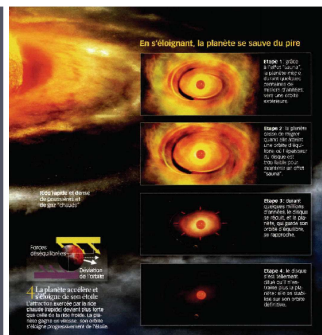
2010 > MAI > SCIENCE & VIE 77



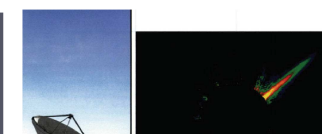
de la Terre. C'est, en effet, ce que les simulations informatiques ont montré : les planètes seraient précipitées vers le Soleil avant même d'avoir atteint leur taille définitive. Mais un nouveau modèle semble résoudre définitivement ce paradoxe.



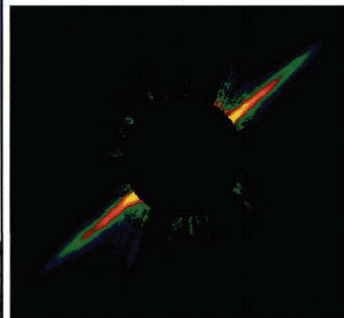
2010 > MAI > SCIENCE & VIE 77



de la Terre. C'est, en effet, ce que les simulations informatiques ont montré : les planètes seraient précipitées vers le Soleil avant même d'avoir atteint leur taille définitive. Mais un nouveau modèle semble résoudre définitivement ce paradoxe.



2010 > MAI > SCIENCE & VIE 77



"Nous avons obtenu un modèle sur cinq millions d'années où la Terre ne tombe pas sur le Soleil"

Wladimir Lyra, astronome brésilien

disque se dilate et écarte l'effet solaire avant de se faire sentir. Inaugure la planète dans la configuration classique : un disque – dilaté mais encore "entraîné" – qui la pousse vers le Soleil...

DONNÉES CONCRÈTES DÈS 2011

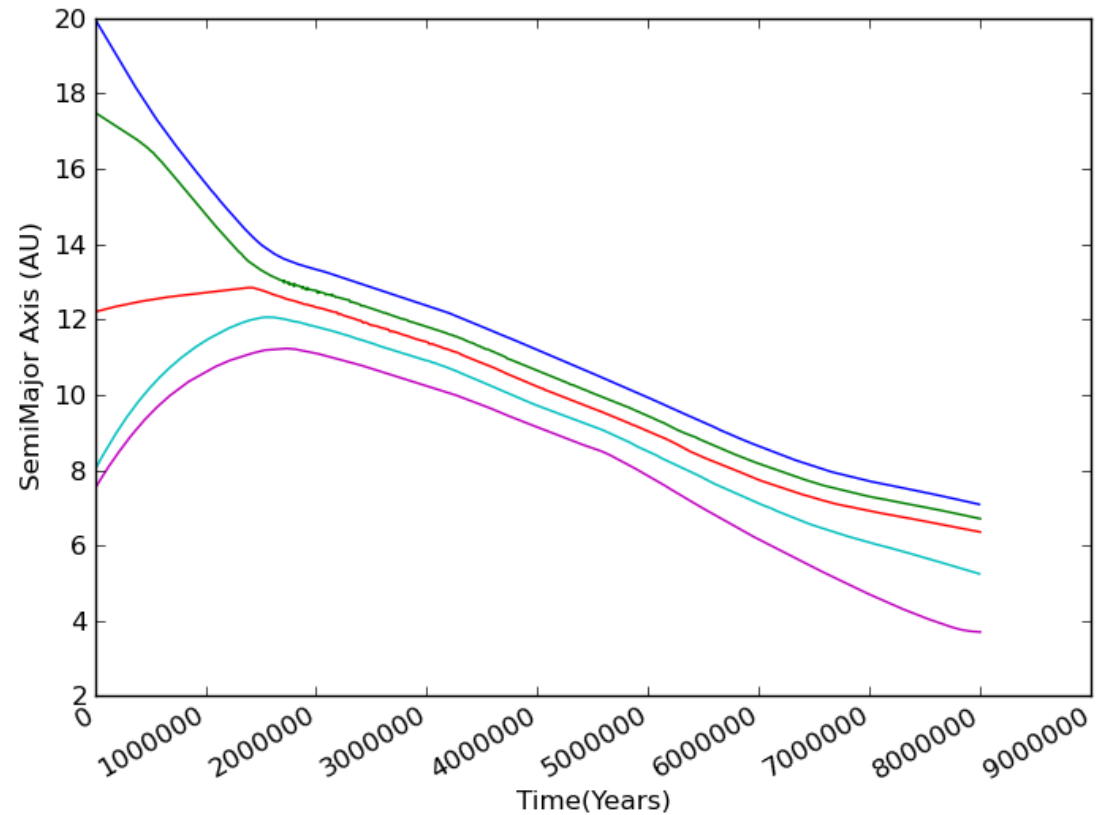
"En mai 2011, raconte Wladimir Lyra, Michael Mac Low m'a proposé d'élaborer ce effort sur cinq millions d'années où la Terre ne tombe pas sur le Soleil." Ainsi testé sur des planètes de masses égales à celle de la Terre, des fois

plus légères ou des fois plus lourdes, et d'éloignement au Soleil variant entre 0,1 et 20 UA (1 UA est la distance moyenne Terre-Soleil, soit 150 millions de kilomètres), le modèle a été formel : grâce à l'effet solaire, la Terre, mais aussi toutes les planètes des systèmes solaires de l'Univers, ont été sauvées ! "On tient là un filon très intéressant, confirme Frédéric Masset, mais il reste des questions copieusement ouvertes, notamment des effets que leur modèle a mis de côté mais qui pourraient modifier le rapport des forces en jeu... Néanmoins, cela se devait pas remettre en question

EN SAVOIR PLUS
Système solaire, systèmes stellaires, de Thérèse ENCRENAZ, éd. Dunod, 2005

2010 > MAI > SCIENCE & VIE 151

Migration + N-Body in Evolutionary Models



Migration in resonance !

See also
Sandor, Lyra & Dullemond (2011)
Hellary & Nelson (2012)

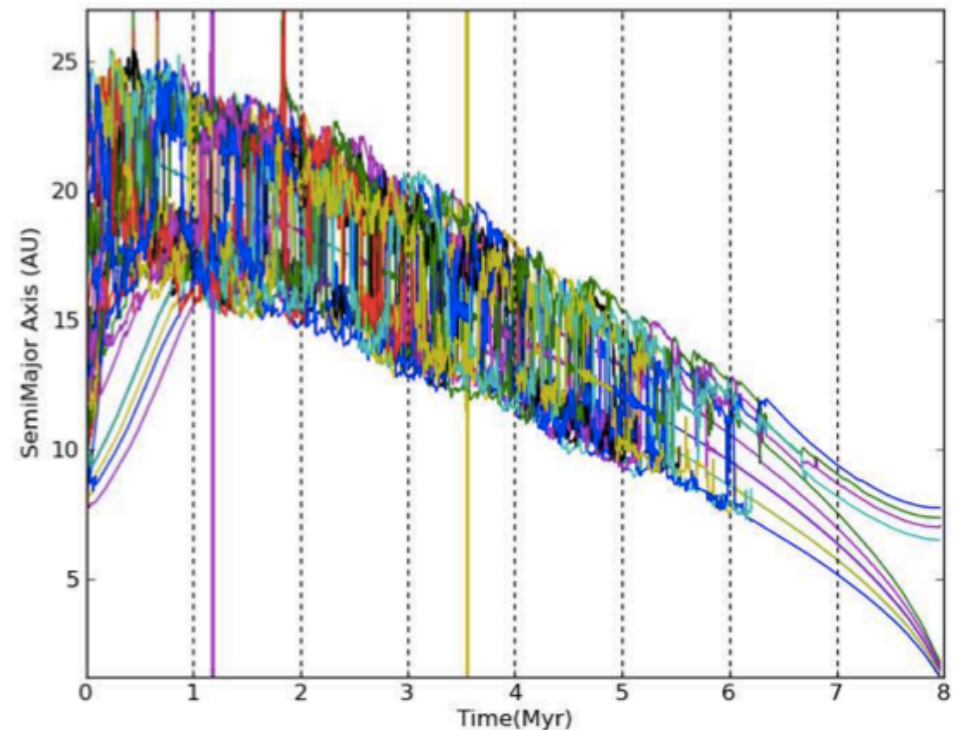
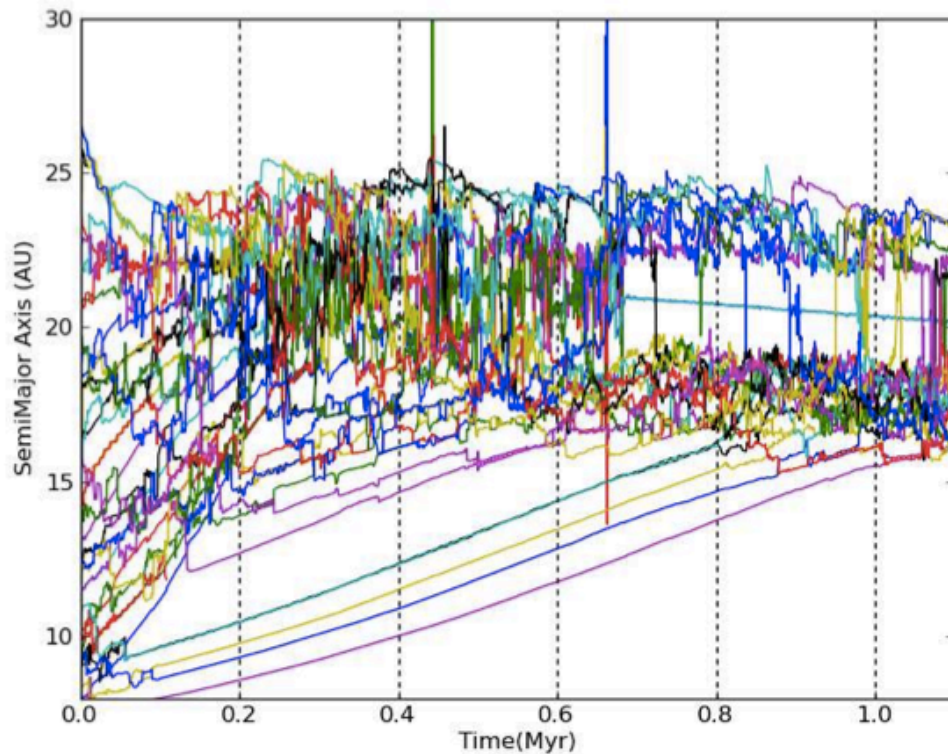
Orbital migration of interacting planets in a radiative evolutionary model

Combines

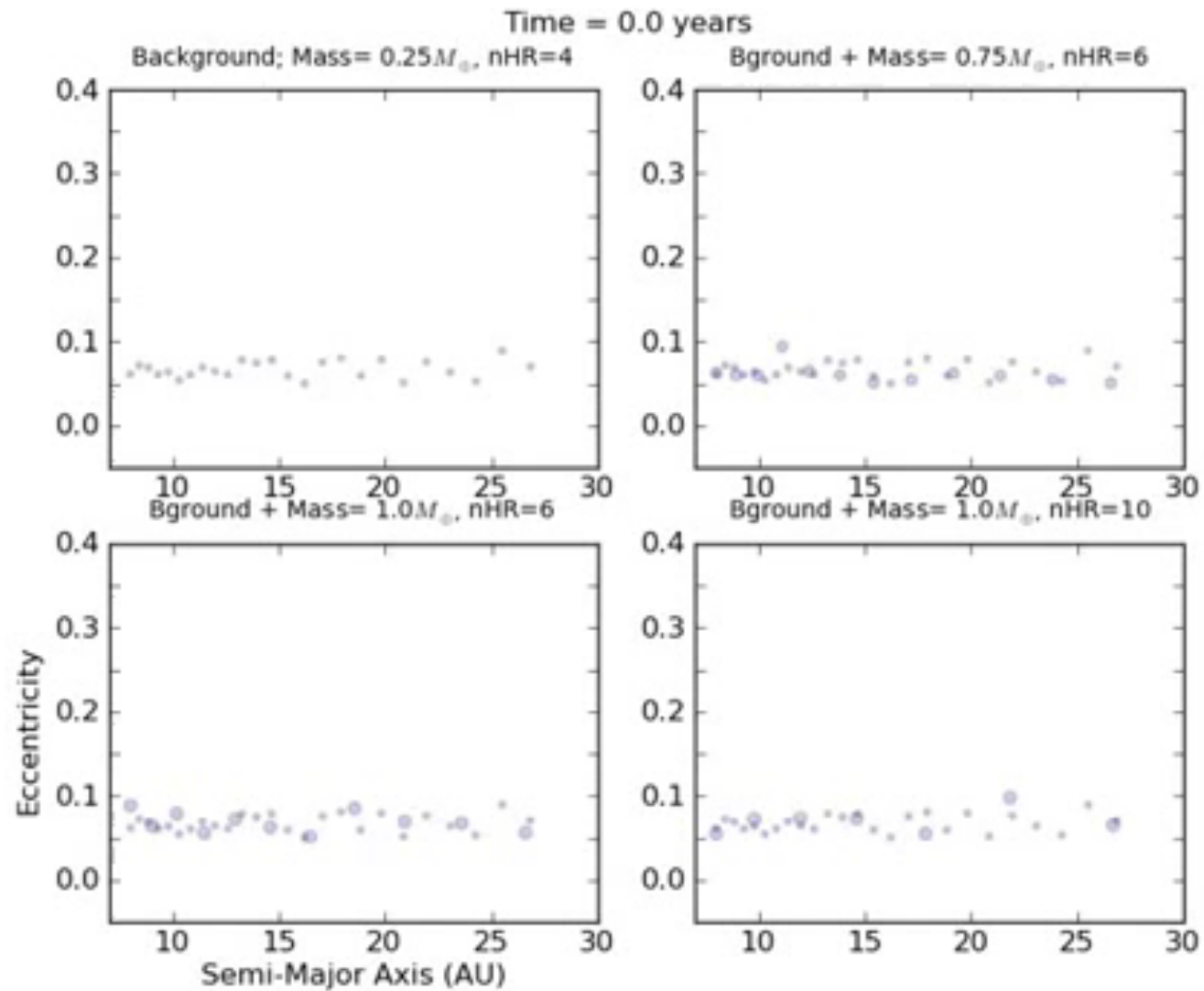
migration + N-body + photoevaporation + turbulence

modelled as stochastic forcing

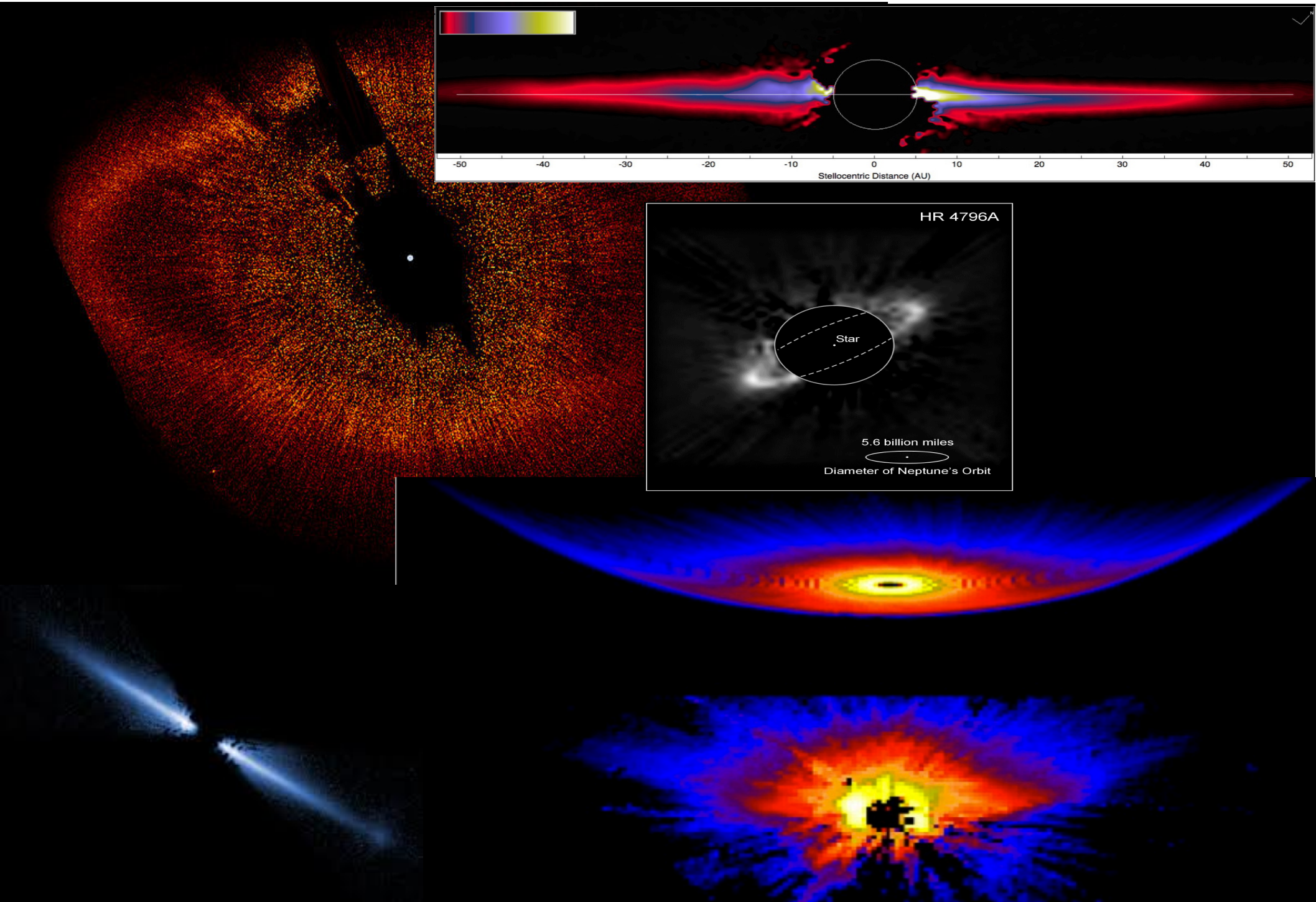
(Laughlin et al. 2004, Ogiwara et al. 2007)



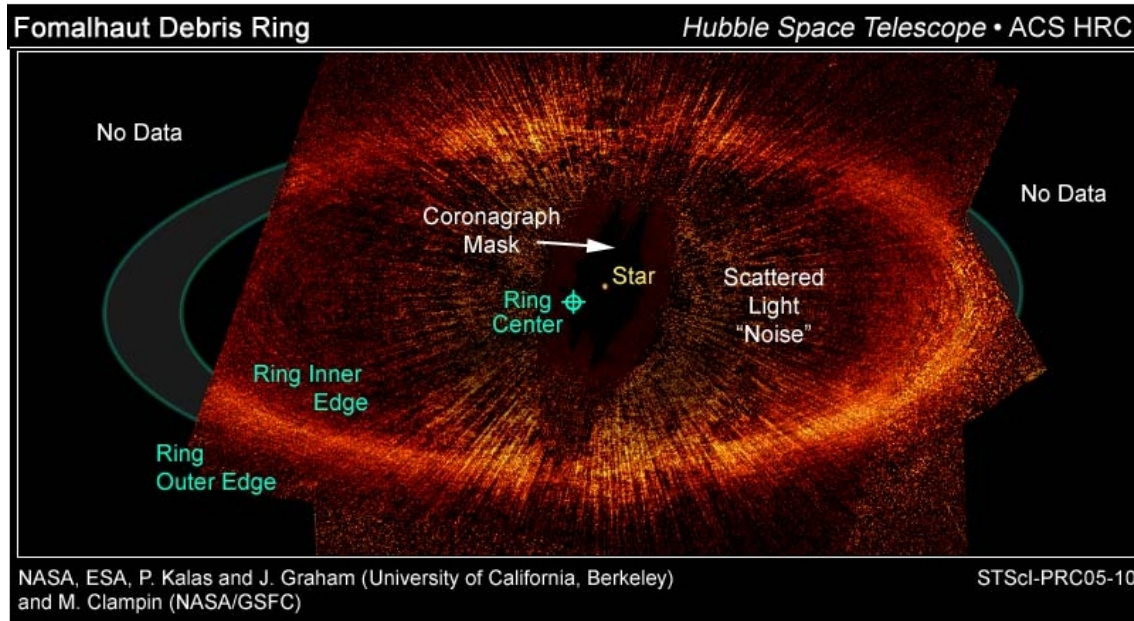
Orbital migration of interacting planets in a radiative evolutionary model



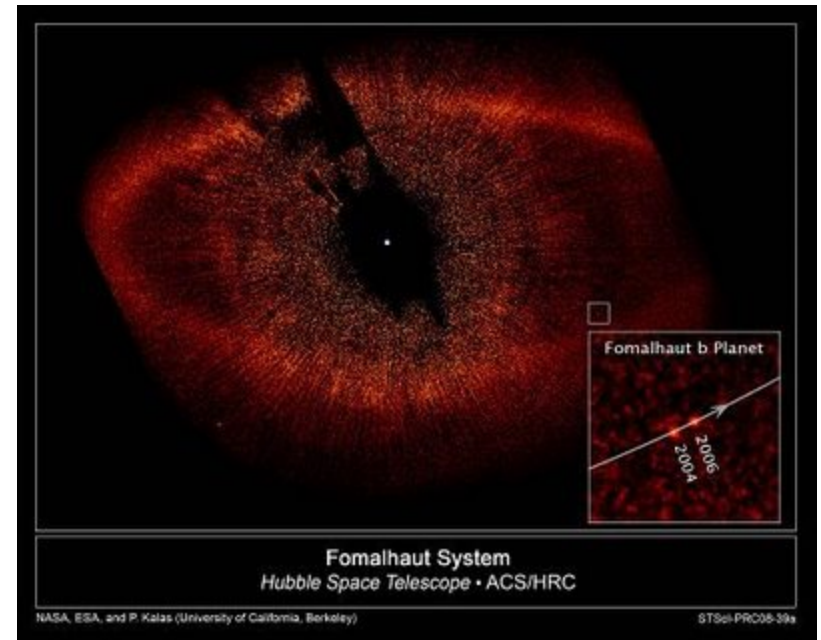
Debris disks - The gas-poor phase



Sharp and eccentric rings in debris disks: Signposts of planets



Narrow sharp eccentric ring



Detection of a source
quickly heralded as a planet
Fomalhaut b

Sharp and eccentric rings in debris disks: Signposts of planets

However.....

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doi:10.1088/0004-637X/747/2/116

INFRARED NON-DETECTION OF FOMALHAUT b: IMPLICATIONS FOR THE PLANET INTERPRETATION

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ABSTRACT

The nearby A4-type star Fomalhaut hosts a debris belt in the form of an eccentric ring, which is thought to be caused by dynamical influence from a giant planet companion. In 2008, a detection of a point source inside the inner edge of the ring was reported and was interpreted as a direct image of the planet, named Fomalhaut b. The detection was made at $\sim 600\text{--}800\text{ nm}$, but no corresponding signatures were found in the near-infrared range, where the bulk emission of such a planet should be expected. Here, we present deep observations of Fomalhaut with *Spitzer*/IRAC at $4.5\text{ }\mu\text{m}$, using a novel point-spread function subtraction technique based on angular differential imaging and Locally Optimized Combination of Images, in order to substantially improve the *Spitzer* contrast at small separations. The results provide more than an order of magnitude improvement in the upper flux limit of Fomalhaut b and exclude the possibility that any flux from a giant planet surface contributes to the observed flux at visible wavelengths. This renders any direct connection between the observed light source and the dynamically inferred giant planet highly unlikely. We discuss several possible interpretations of the total body of observations of the Fomalhaut system and find that the interpretation that best matches the available data for the observed source is scattered light from a transient or semi-transient dust cloud.

Key words: circumstellar matter – planetary systems – stars: early-type

Online-only material: color figures

Planet not detected in infrared

Sharp and eccentric rings in debris disks: Signposts of planets

However (take two)....

Direct Imaging Confirmation and Characterization of a Dust-Enshrouded Candidate Exoplanet Orbiting Fomalhaut

Thayne Currie^{1,2}, John Debes³, Timothy J. Rodigas⁴, Adam Burrows⁵, Yoichi Itoh⁶,
Misato Fukagawa⁷, Scott J. Kenyon⁸, Marc Kuchner², Soko Matsumura⁹

currie@astro.utoronto.ca

ABSTRACT

We present Subaru/IRCS J band data for Fomalhaut and a (re)reduction of archival 2004–2006 HST/ACS data first presented by Kalas et al. (2008). We confirm the existence of a candidate exoplanet, Fomalhaut b, in both the 2004 and 2006 F606W data sets at a high signal-to-noise. Additionally, we confirm

h.EP] 24 Oct 2012

Fomalhaut

No

NASA, ESA,
and M. Clarm

Fomalhaut b Planet

STScI-PRC08-39a

It should not have been detected anyway...

Some of the Fom b controversy

Janson et al. 2012

Variability by 0.7-0.8 mag in F606W band

Astrometric orbit not apsidally aligned with the ring

No infrared emission

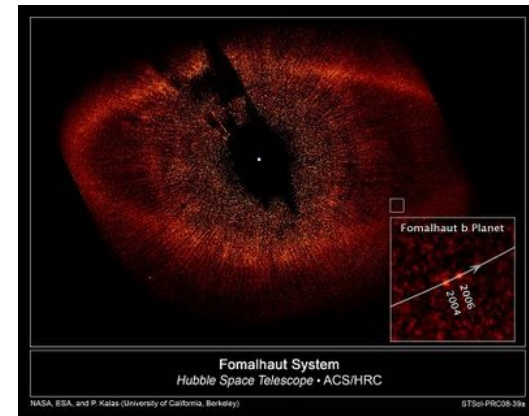
Currie et al. 2012

No variability found within 0.15 mag in the same band

Consistent with apsidal alignment

Thermal emission from 0.5 MJ would not be detectable.

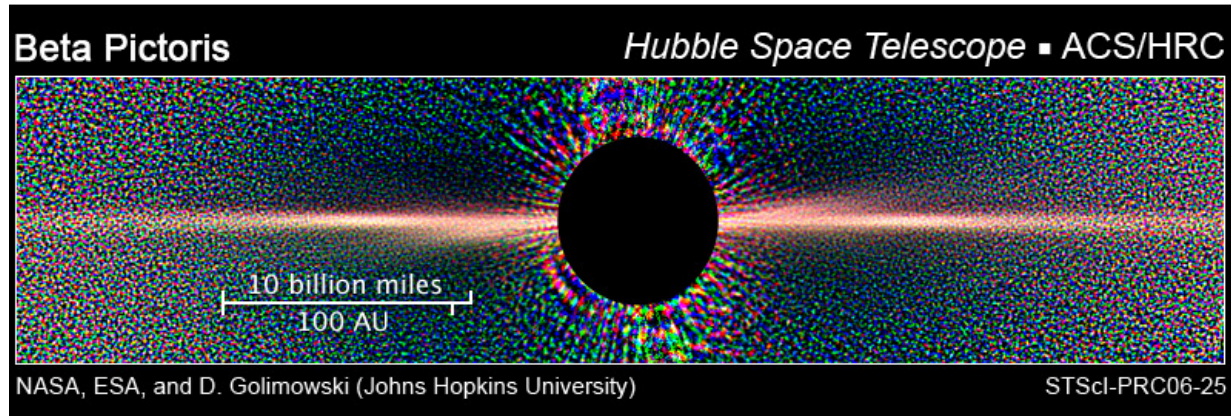
Observed optical emission requires reflection by *something* of several Jupiter radii



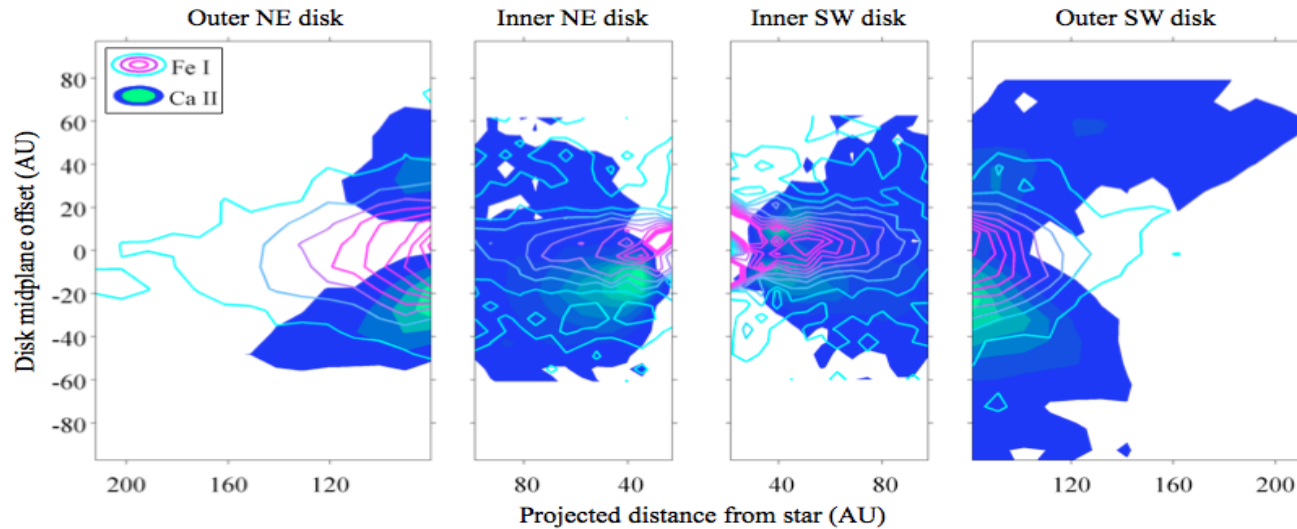
Are there
alternative explanations?

Debris disks are not completely gas-free

Dust



Gas



VLT imaging by
Nilsson et al. (2012)

Gas in debris disks

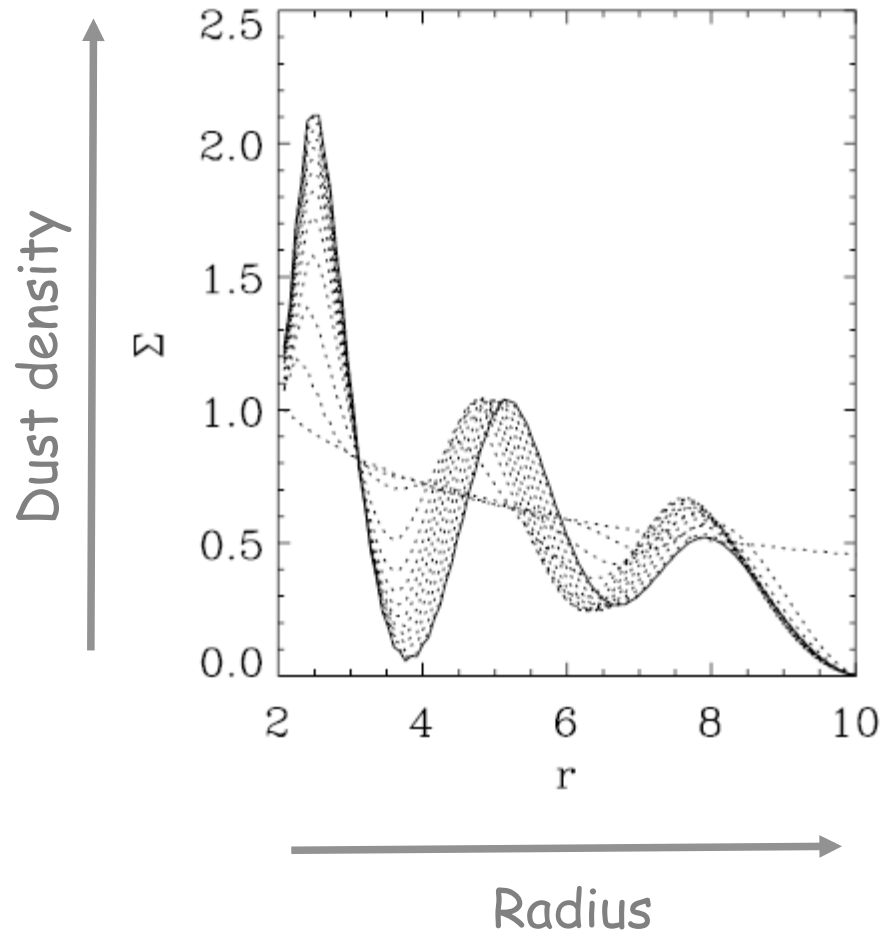
Detections

β Pictoris	many species	Lagrange et al. (1998), ...
51 Ophiuchi	many species	Roberge et al. (2002)
σ Herculis	C II, N II	Chen & Jura (2003)
HD 32297	Na I, CII	Redfield (2007), Donaldson et al. (2012)
HD 135344	H ₂ , CO	Thi et al. (2001), Pontoppidan et al. (2008)
49 Ceti	H ₂ , CO	Dent et al. (2005), Roberge et al. (2012)
AU Mic	H ₂	France et al. (2007)
HD172555	SiO	Lisse et al. (2009)

Source of gas: Outgassing processes

Infalling comets	Beust & Valiron (2007)
Grain sublimation	e.g. Rafikov (2012)
Grain-Grain collisions	Czechowski & Mann (2007)
Photo-stimulated desorption	Chen et al. (2007)
Planet-Planet collisions	Van den Ancker (2001), Lisse (2008)
Primordial?	

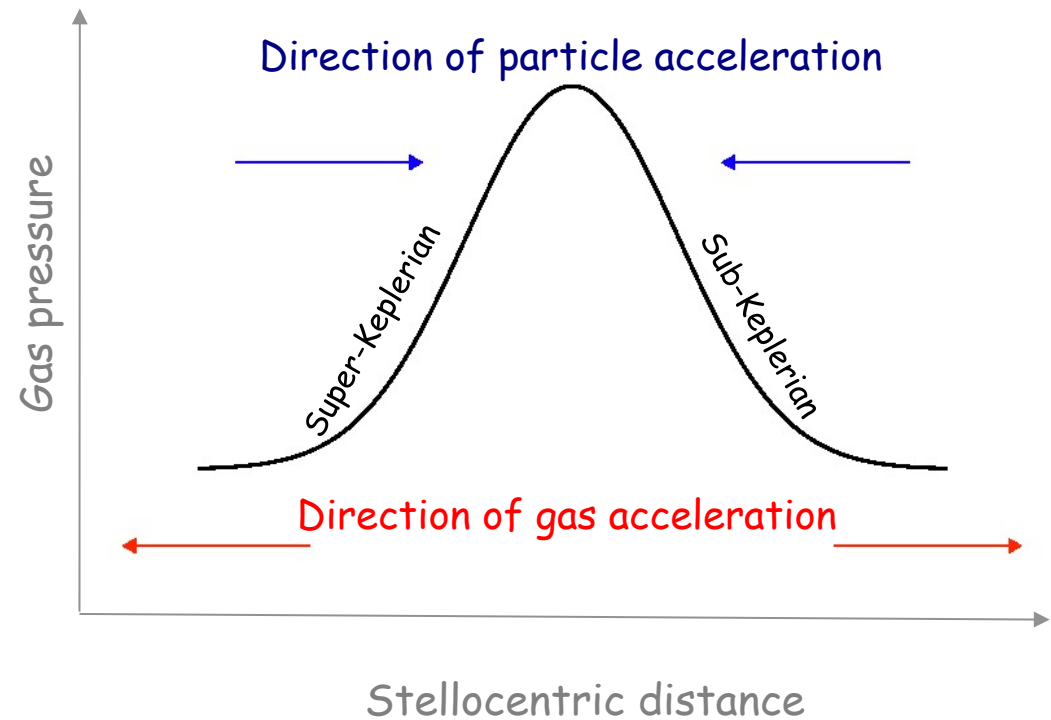
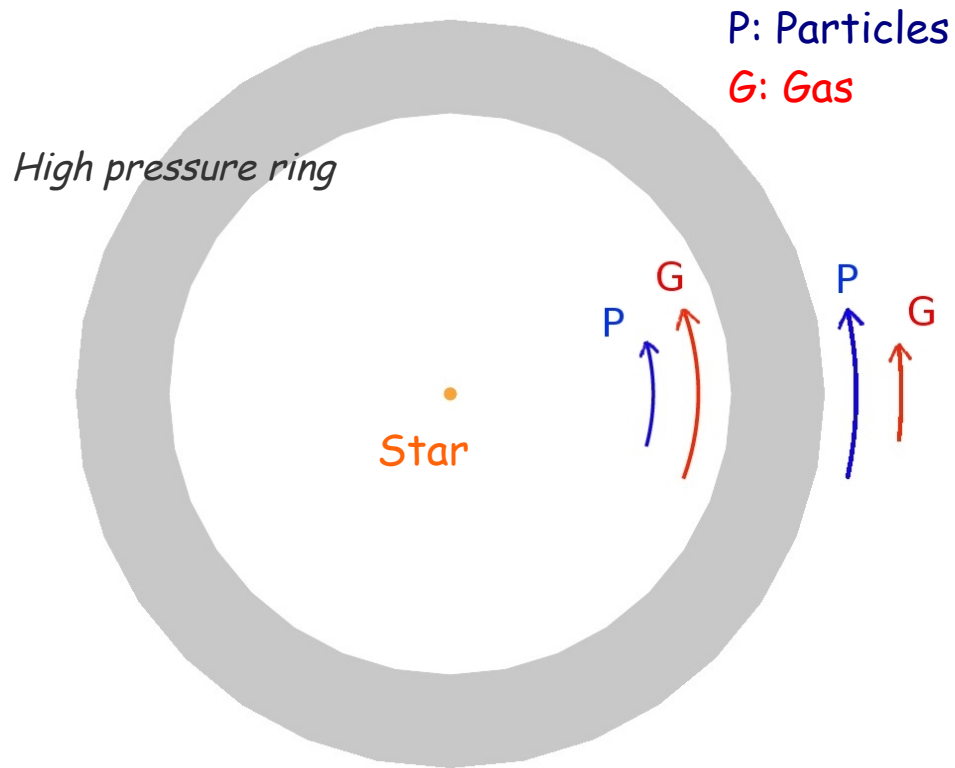
Dust and gas together leads to instability...



Klahr & Lin (2005)

Suggested that an instability might cause dust in debris disks to clump together.

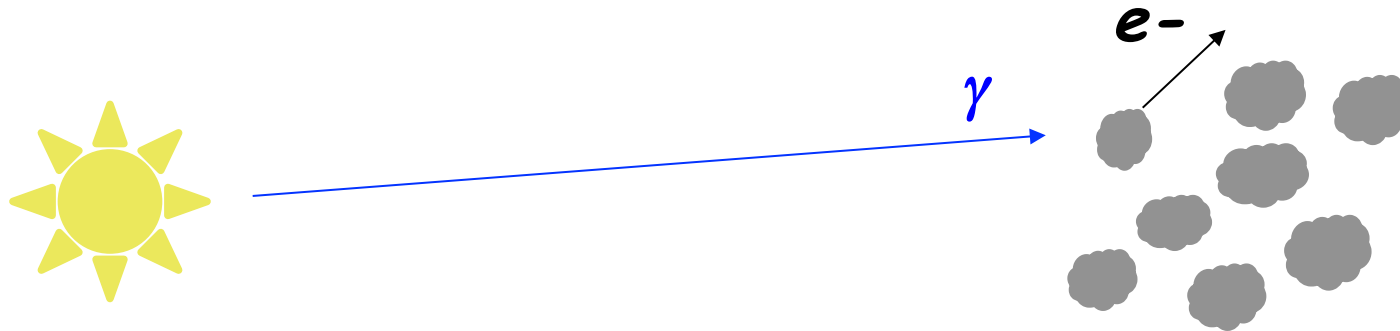
Particles move toward pressure maxima



Adapted from Whipple (1972)

Photoelectric heating

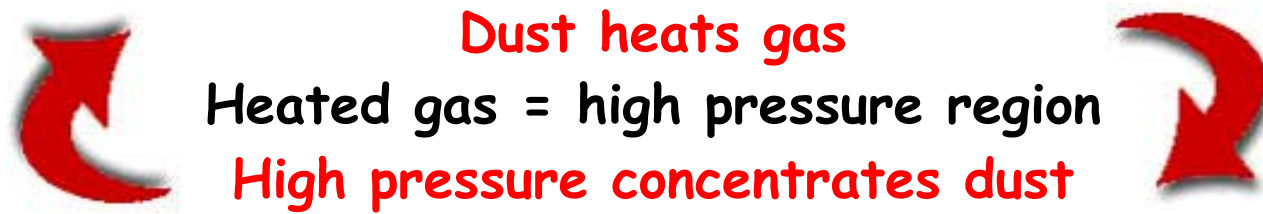
In optically thin debris disks,
the **dust** is the **main heating agent** for the gas.



Dust intercepts starlight directly,
emits electron, that heats the gas.

Gas is photoelectrically heated by the dust itself

Runaway process: instability



Model equations

Klahr & Lin (2005) used a simplified, 1-D model.

$$\frac{\partial}{\partial t} \Sigma_d + \frac{1}{r} \frac{\partial}{\partial r} r \Sigma_d v_r = 0.$$

Continuity equation

$$V_\phi = \Omega r + \frac{1}{2\Omega \Sigma_g} \frac{\partial}{\partial r} P$$

Terminal velocity

$$T_g = T_0 \left(\frac{\Sigma_d}{\Sigma_0} \right)^\beta,$$

Equation of state

Model equations

Our simulation adds much more physics, and works in 2D.

Klahr & Lin (2005)

1D

$$\frac{\partial}{\partial t} \Sigma_d + \frac{1}{r} \frac{\partial}{\partial r} r \Sigma_d v_r = 0.$$

$$V_\phi = \Omega r + \frac{1}{2\Omega \Sigma_g} \frac{\partial}{\partial r} P$$

$$T_g = T_0 \left(\frac{\Sigma_d}{\Sigma_0} \right)^\beta,$$

Inertia for both gas and dust

Energy equation

*Drag force and
drag force backreaction*

Lyra & Kuchner (2012)

2D

$$\frac{\partial \Sigma_g}{\partial t} = -(\mathbf{u} \cdot \nabla) \Sigma_g - \Sigma_g \nabla \cdot \mathbf{u}$$

$$\frac{\partial \mathbf{u}}{\partial t} = -(\mathbf{u} \cdot \nabla) \mathbf{u} - \frac{1}{\Sigma_g} \nabla P - \nabla \Phi - \frac{\Sigma_d}{\Sigma_g} \mathbf{f}_d$$

$$\frac{\partial S}{\partial t} = -(\mathbf{u} \cdot \nabla) S - \frac{c_v}{T} \frac{(T - T_p)}{\tau_T}.$$

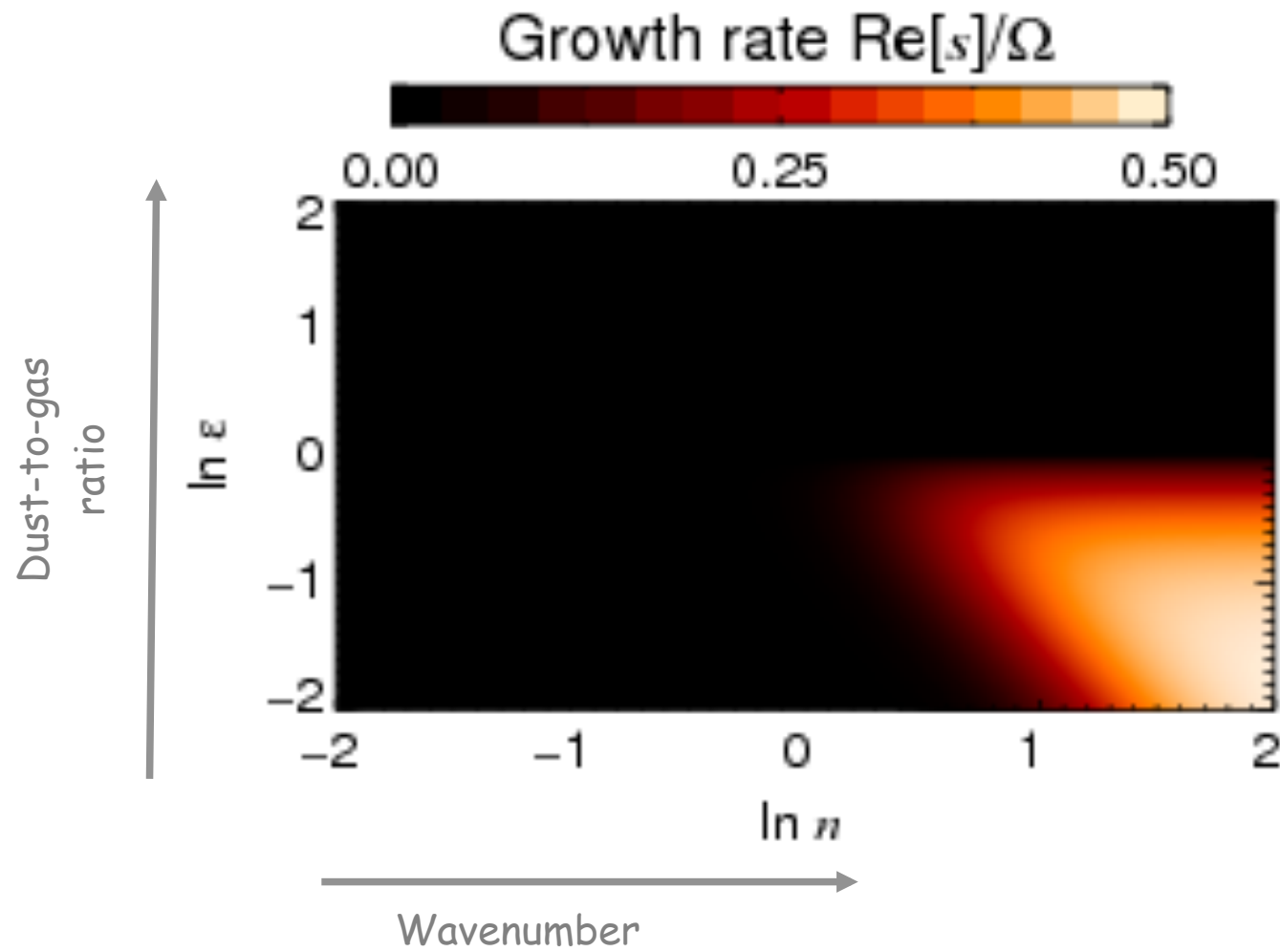
$$\frac{d\mathbf{x}}{dt} = \mathbf{v}$$

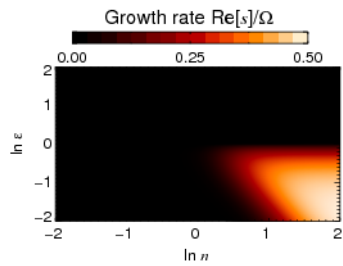
$$\frac{d\mathbf{v}}{dt} = -\nabla \Phi + \mathbf{f}_d$$

$$\mathbf{f}_d = -\frac{(\mathbf{v} - \mathbf{u})}{\tau_f}$$

$$T_p = T_0 \frac{\Sigma_d}{\Sigma_0}.$$

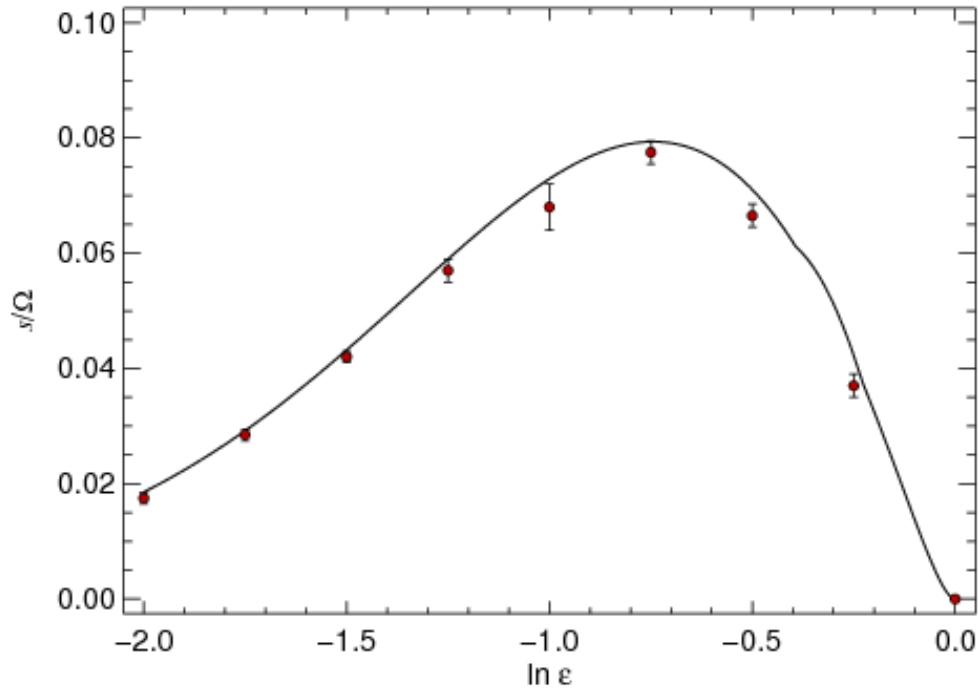
Linear Analysis



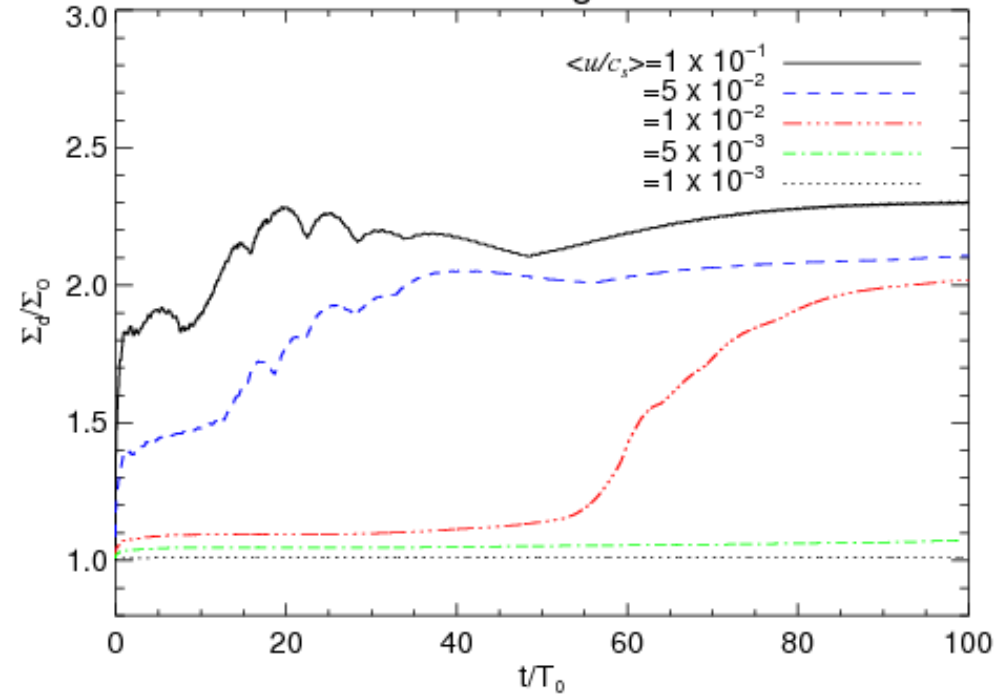


Linear and nonlinear growth

Growth rates $\alpha=10^{-2}$



Nonlinear growth



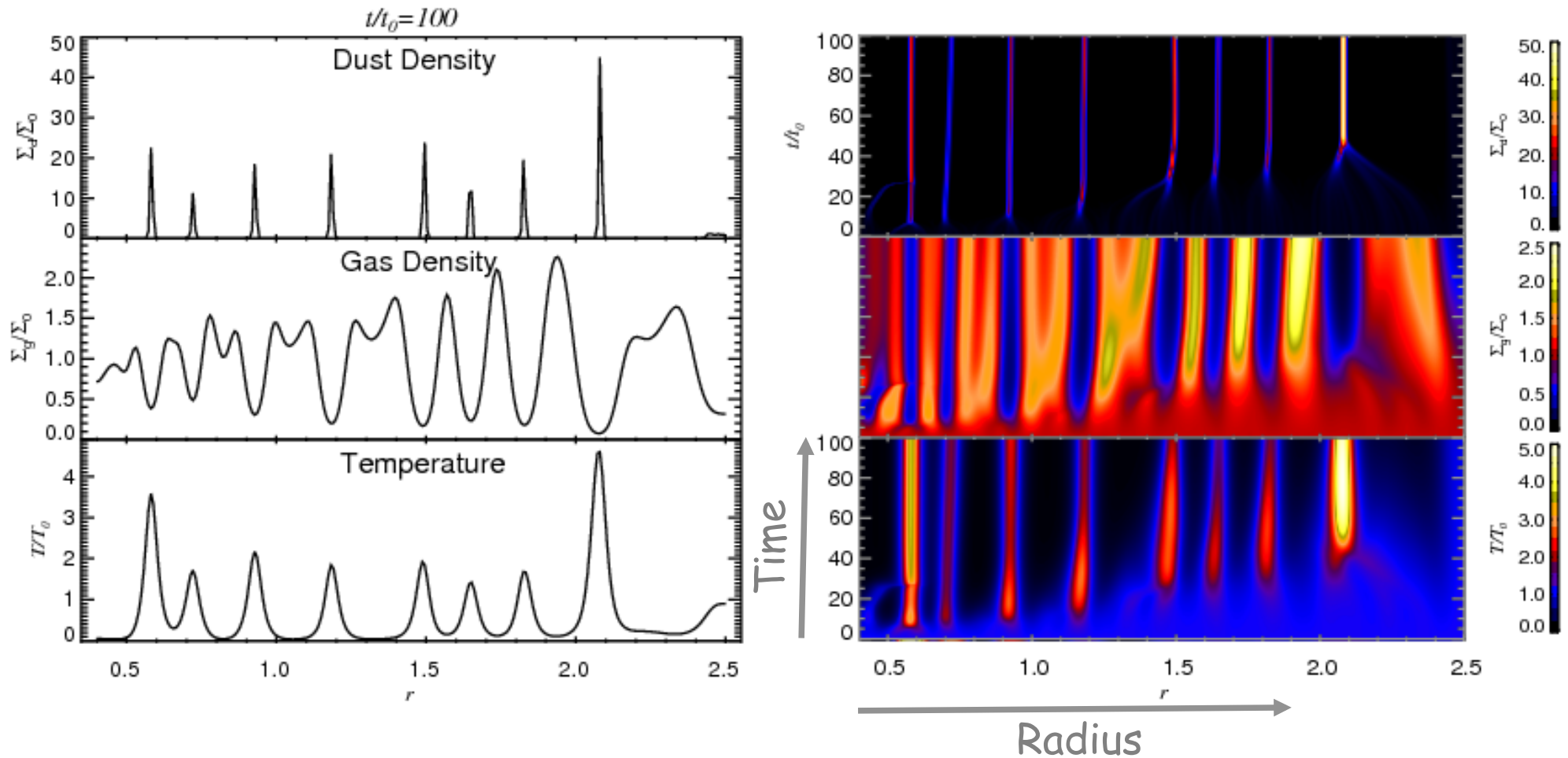
Linear growth only exists for $\epsilon < 1$

But there is
nonlinear growth
beyond !

Instability



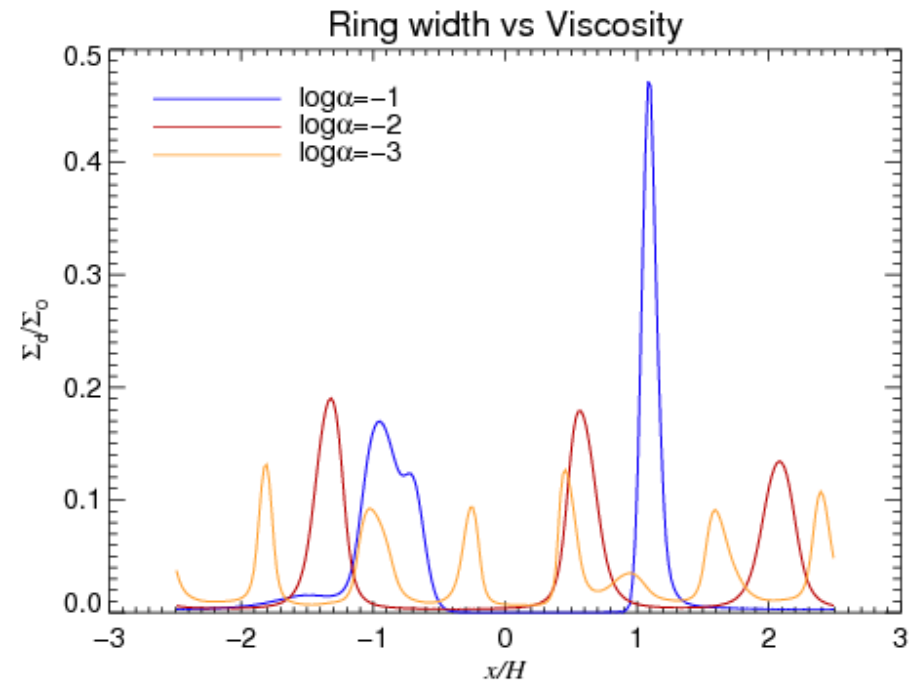
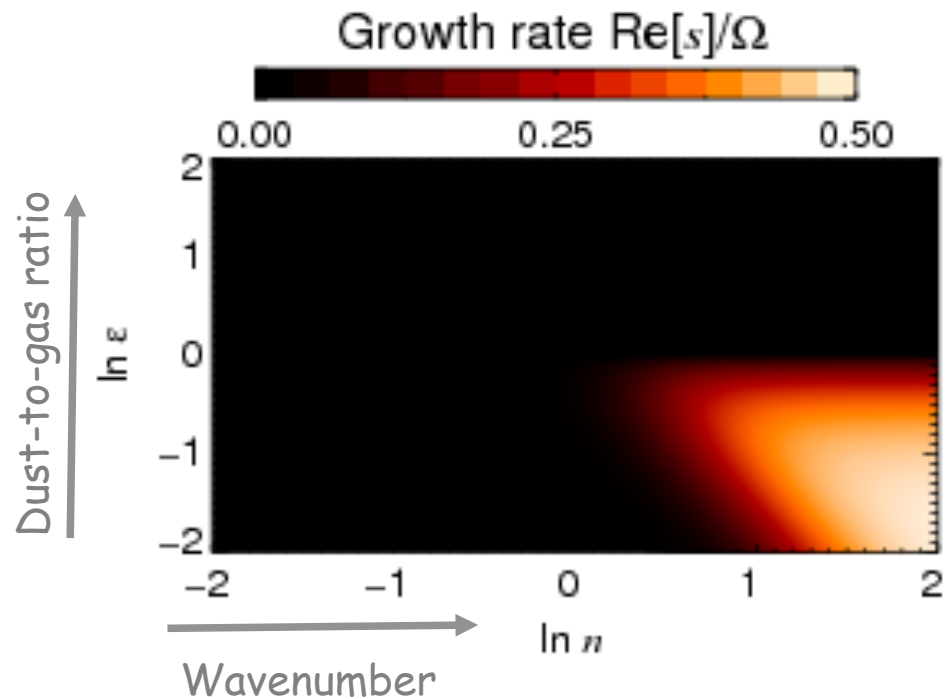
Dust heats gas
Heated gas = high pressure region
High pressure concentrates dust



Narrow hot dust rings
Cold gas collects between rings

Ring width

Ring spacing and width is determined by the wavelength of maximum growth.

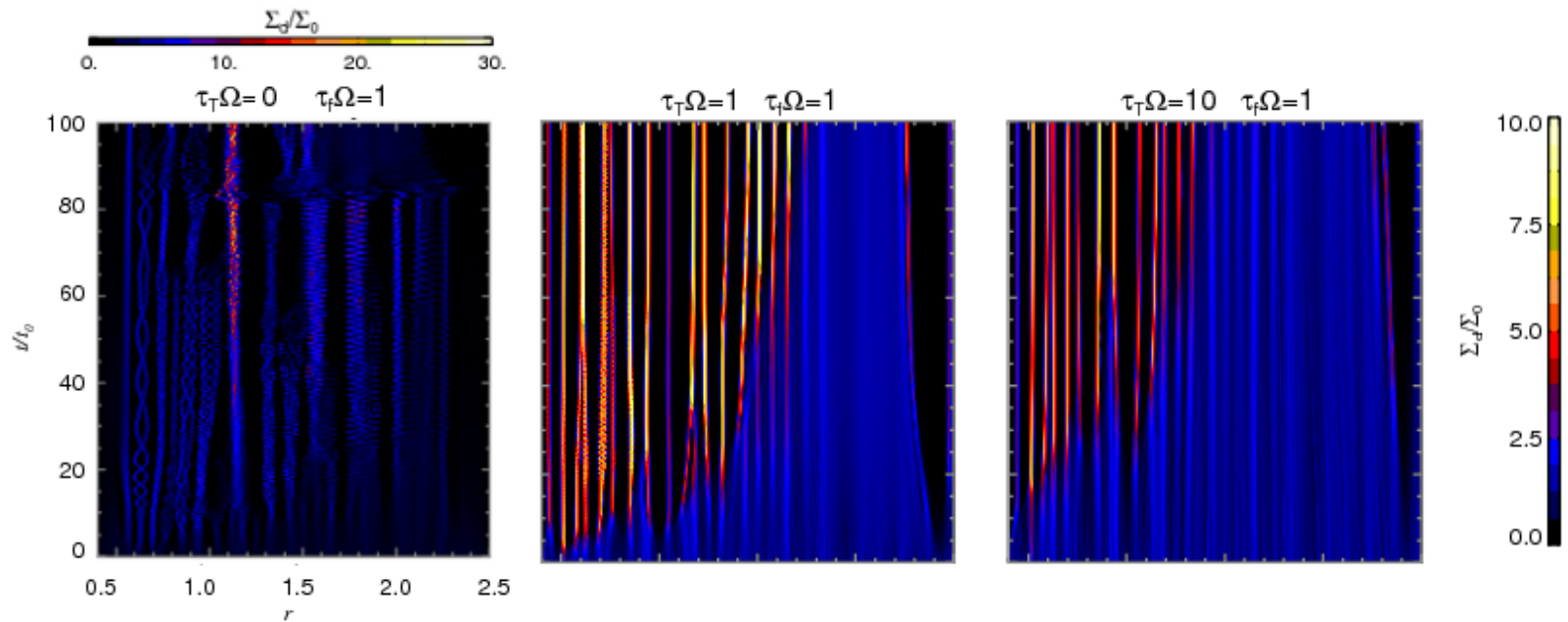


Which in turn is determined by viscosity

Ring width ~ 10 Kolmogorov lengths

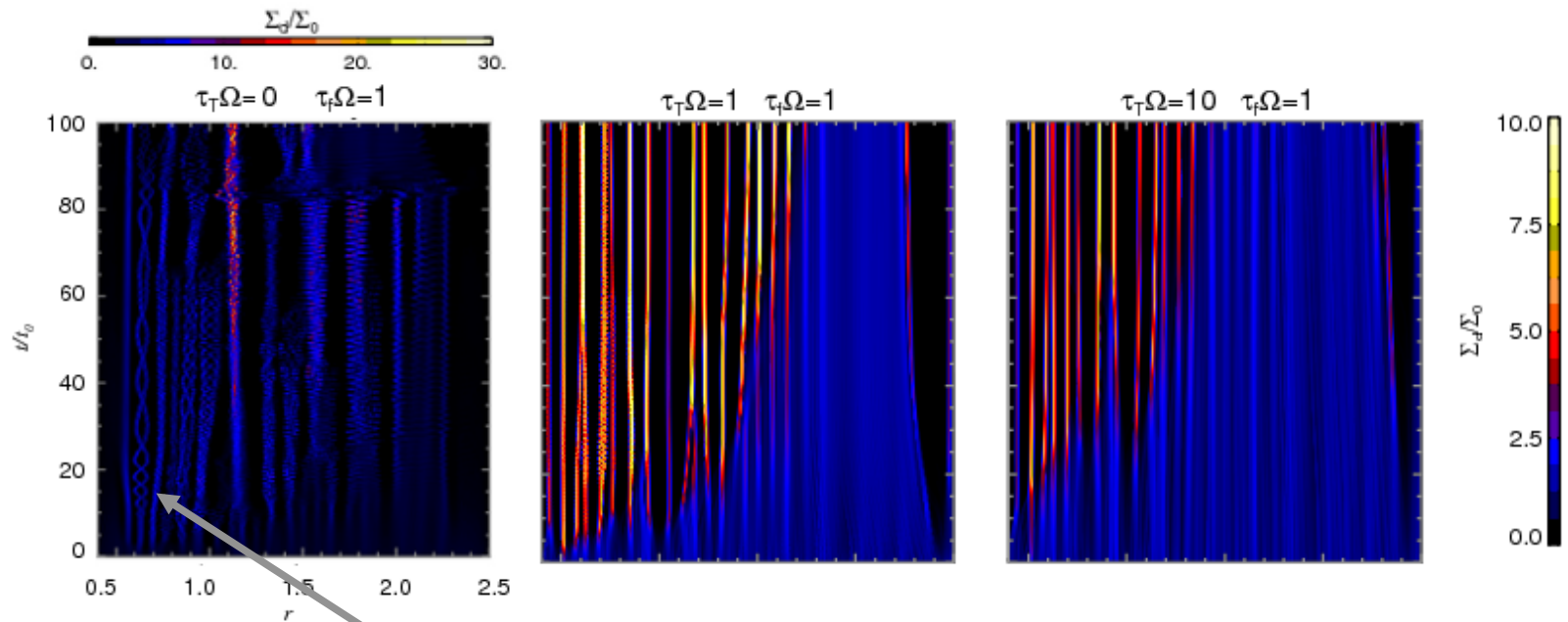
Oscillations

Thermal coupling time



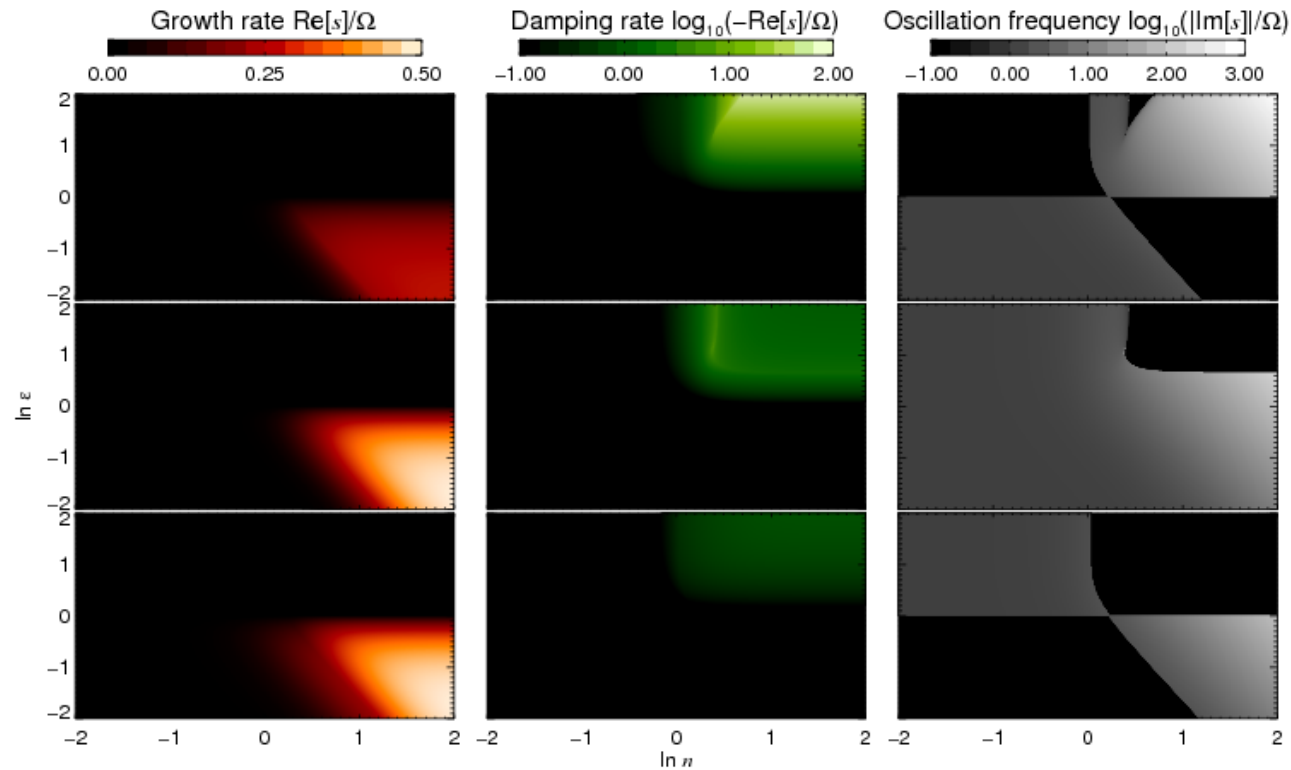
Oscillations

Thermal coupling time



Oscillations appear
with decreasing thermal time.

Solutions



Dispersion relation

$$A\omega^5 + B\omega^4 + C\omega^3 + D\omega^2 + E\omega + F = 0$$

$$A=1$$

$$B=2\epsilon + 2$$

$$C=\epsilon^2 + \epsilon(n^2+2) + 3$$

$$D=\epsilon^2 n^2 + \epsilon(3n^2+2) + 2$$

$$E=\epsilon^2(2n^2+1) + \epsilon(3n^2+2) + 2$$

$$F=\epsilon^2 n^2 - \epsilon n^2$$

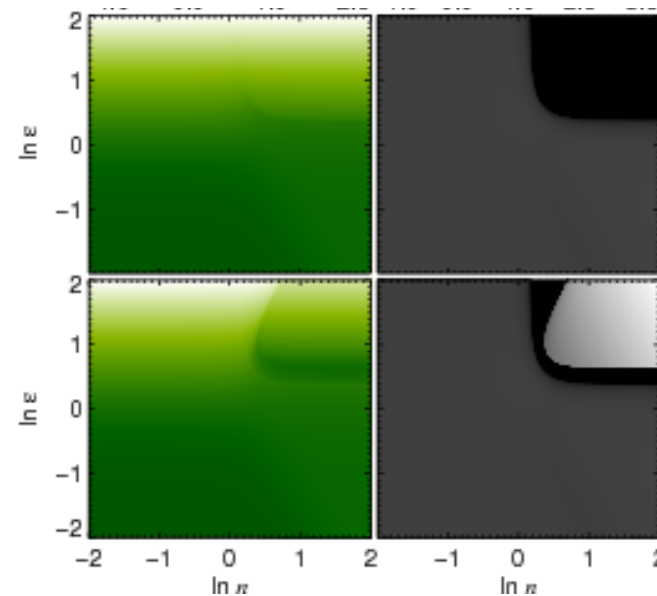
$$\epsilon = \Sigma_d / \Sigma_g$$

$$n = kH$$

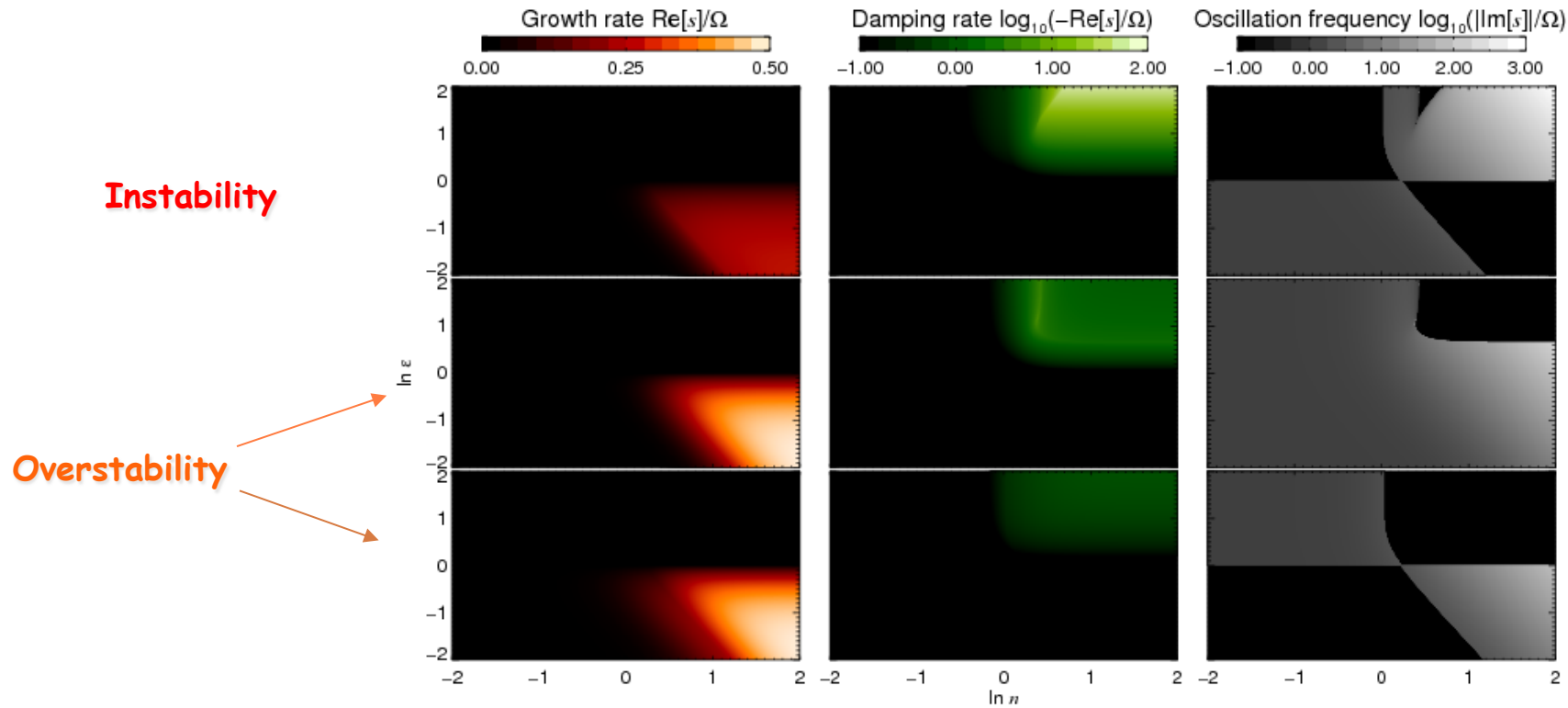
$$\omega = s/\Omega$$

Dust-to-gas
ratio

Wavenumber



Solutions



Dispersion relation

$$A\omega^5 + B\omega^4 + C\omega^3 + D\omega^2 + E\omega + F = 0$$

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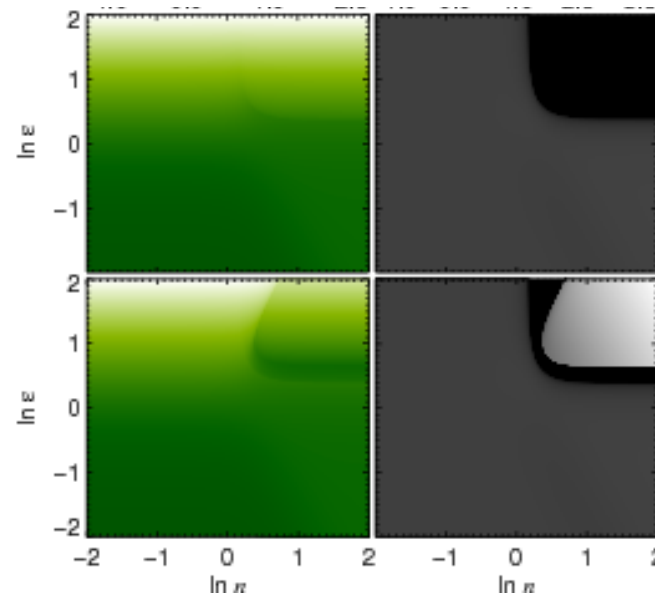
$$\epsilon = \Sigma_d / \Sigma_g$$

$$n = kH$$

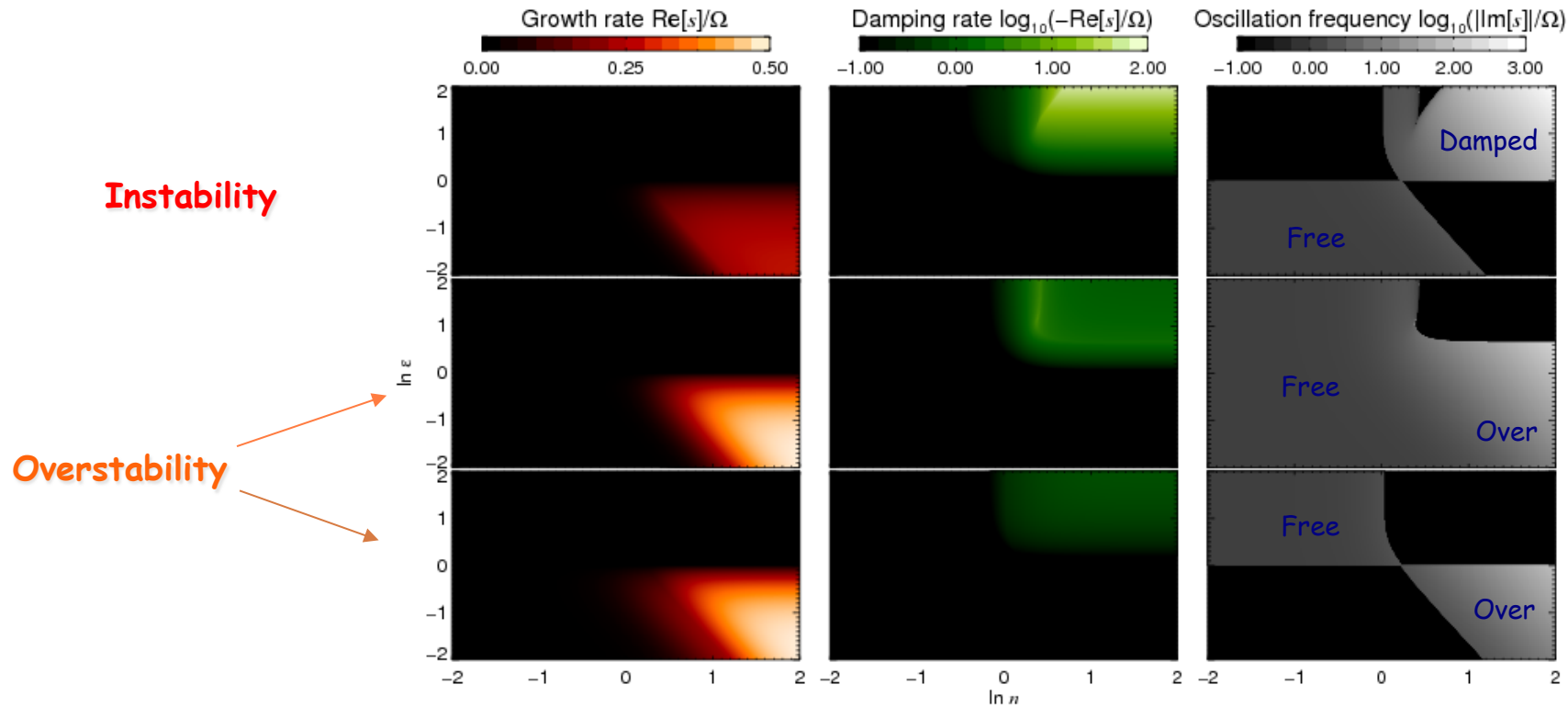
$$\omega = s/\Omega$$

Dust-to-gas
ratio

Wavenumber



Solutions



Dispersion relation

$$A\omega^5 + B\omega^4 + C\omega^3 + D\omega^2 + E\omega + F = 0$$

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$$E=\epsilon^2(2n^2+1) + \epsilon(3n^2+2) + 2$$

$$F=\epsilon^2 n^2 - \epsilon n^2$$

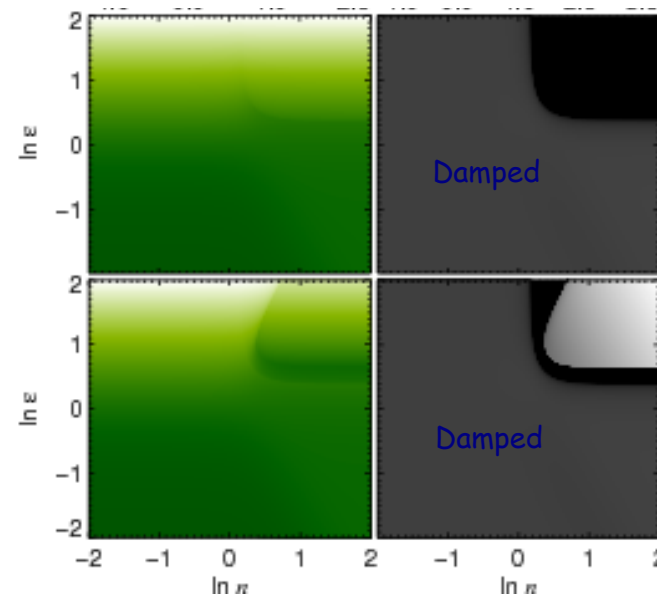
$$\epsilon = \Sigma_d / \Sigma_g$$

$$n = kH$$

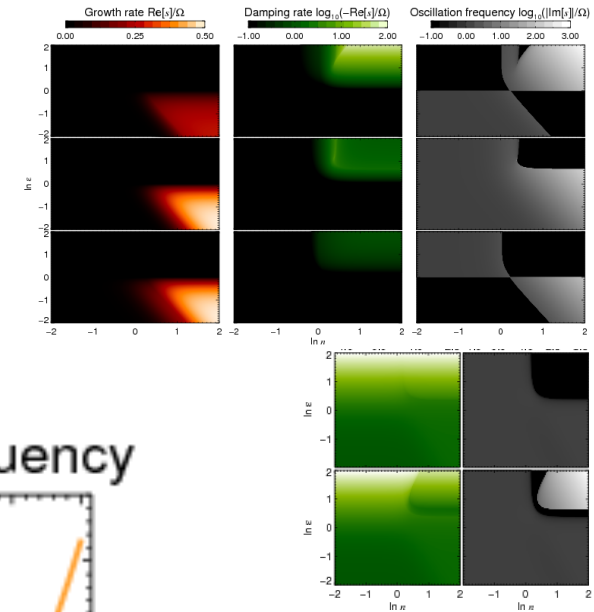
$$\omega = s/\Omega$$

Dust-to-gas
ratio

Wavenumber



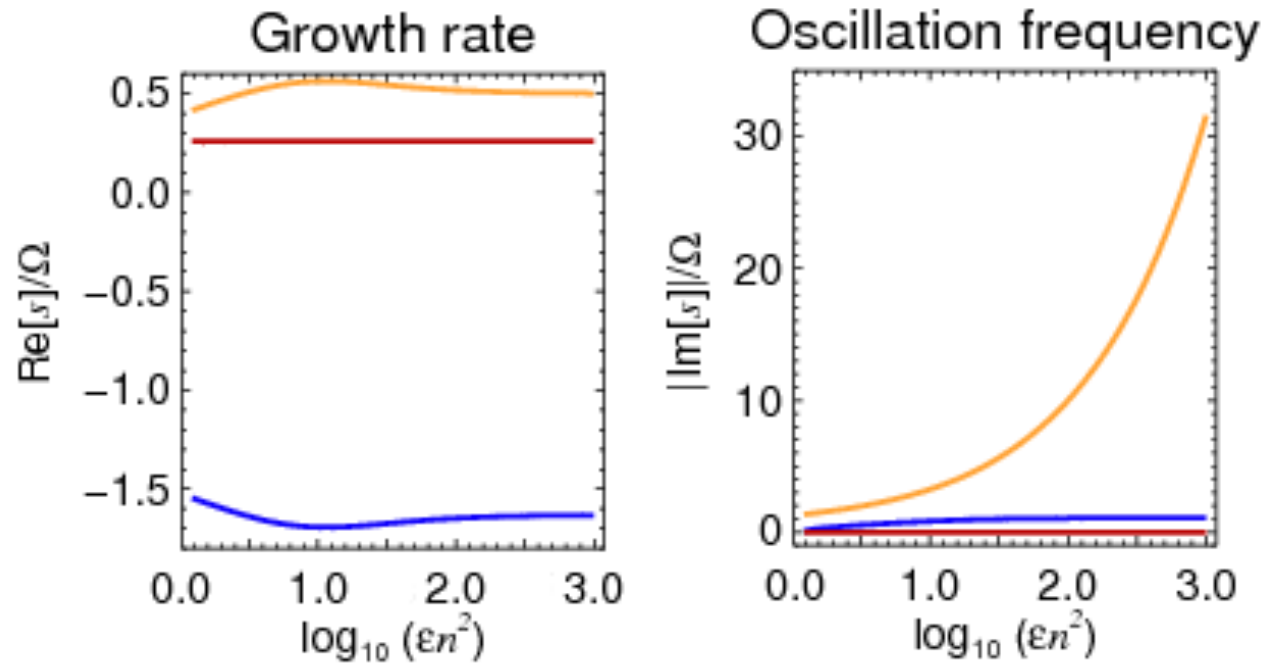
Solutions



Overstability

Instability

Oscillations

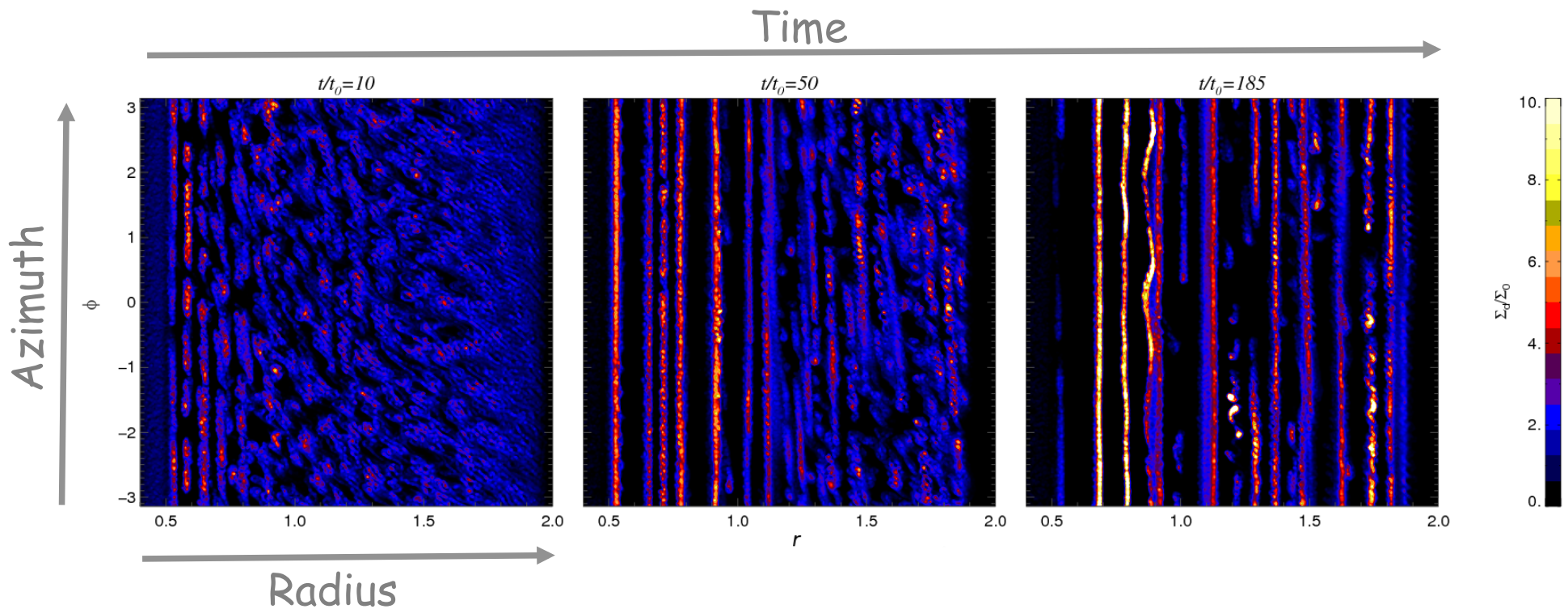


Max growth rate: $\Omega/2$.
Million-fold amplification in five orbits!

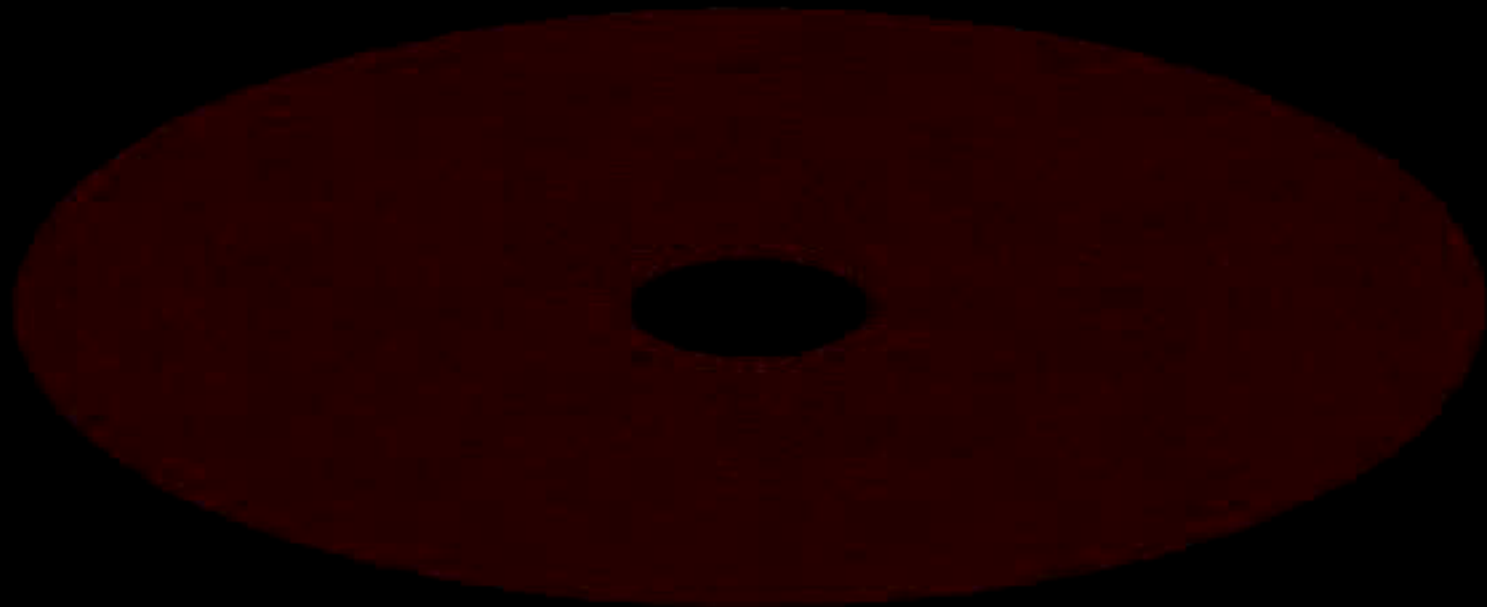
A very powerful instability.

The model in 2D: Eccentric rings

Growth of axisymmetric modes
+
Damping of nonaxisymmetric modes.
= Rings !!!

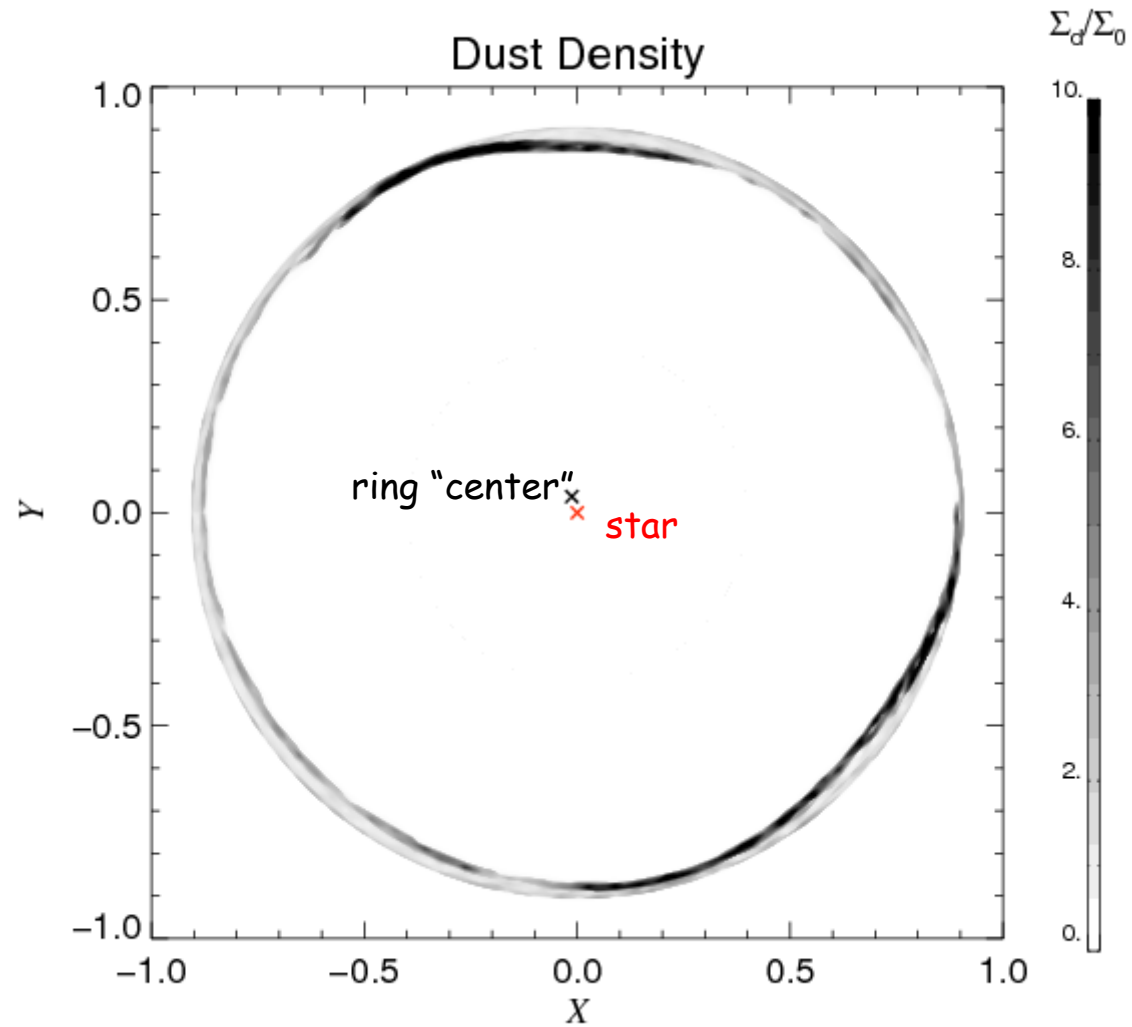


Epicyclic oscillations
make the ring appear *eccentric* !!!



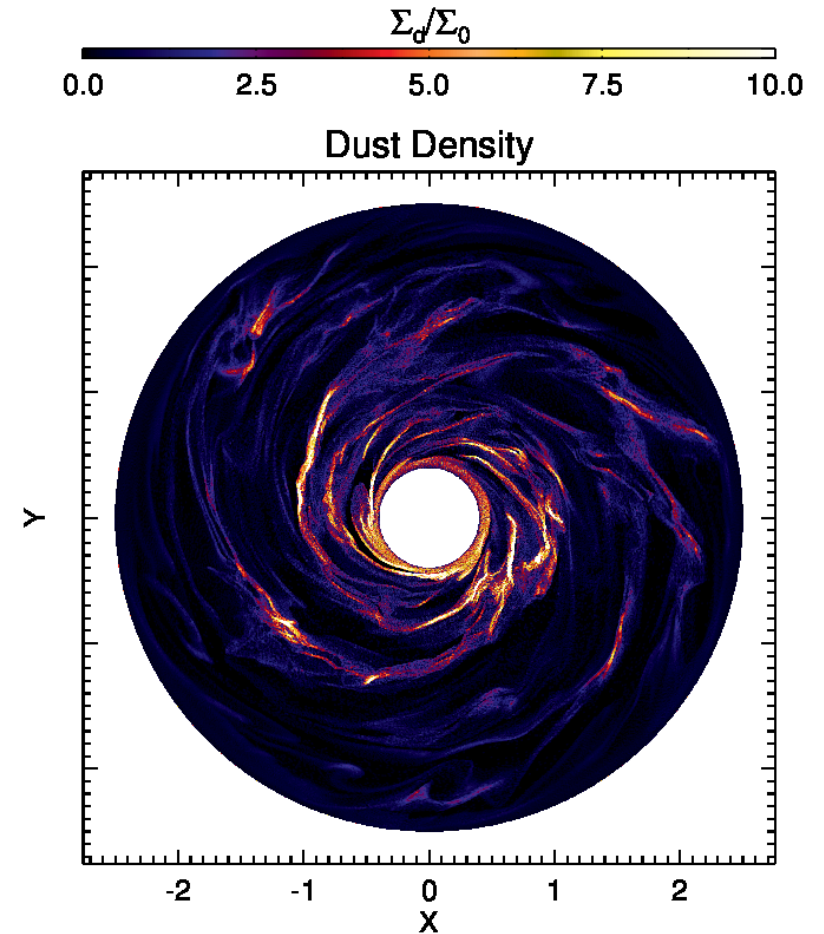
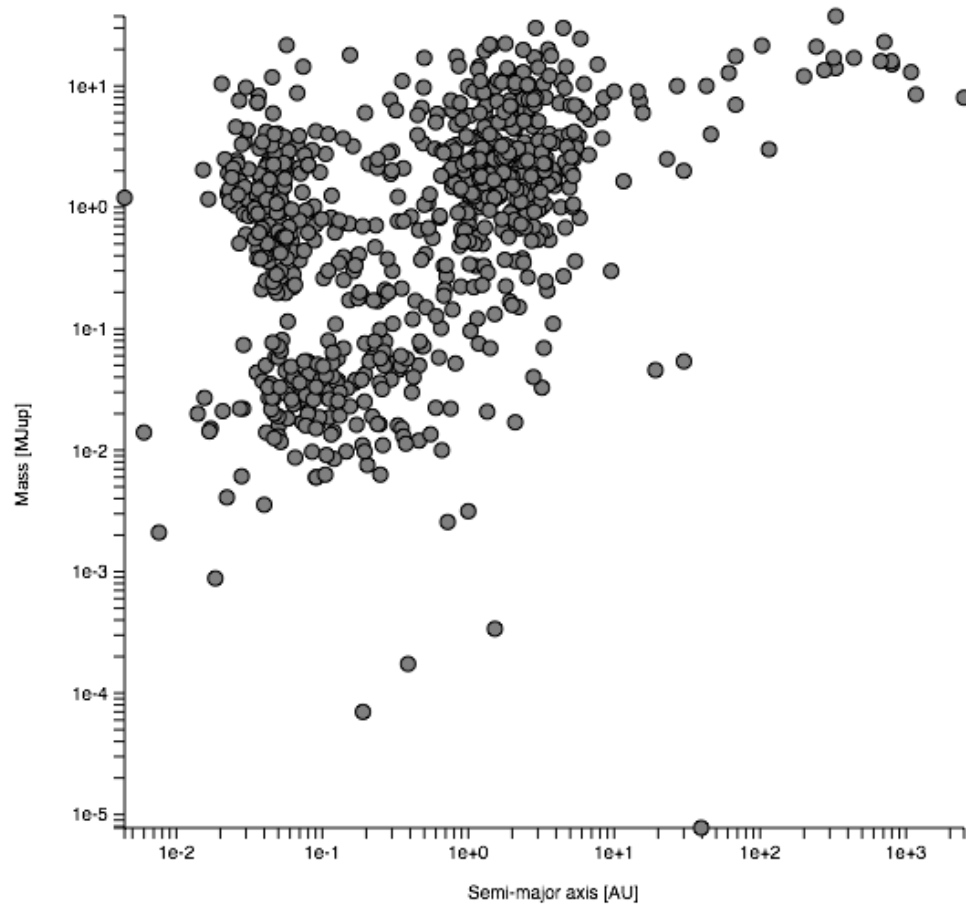
Lyra & Kuchner (2013, Nature, in review)

Ring eccentricity



Eccentricity $e=0.04$

Streaming instability in global models

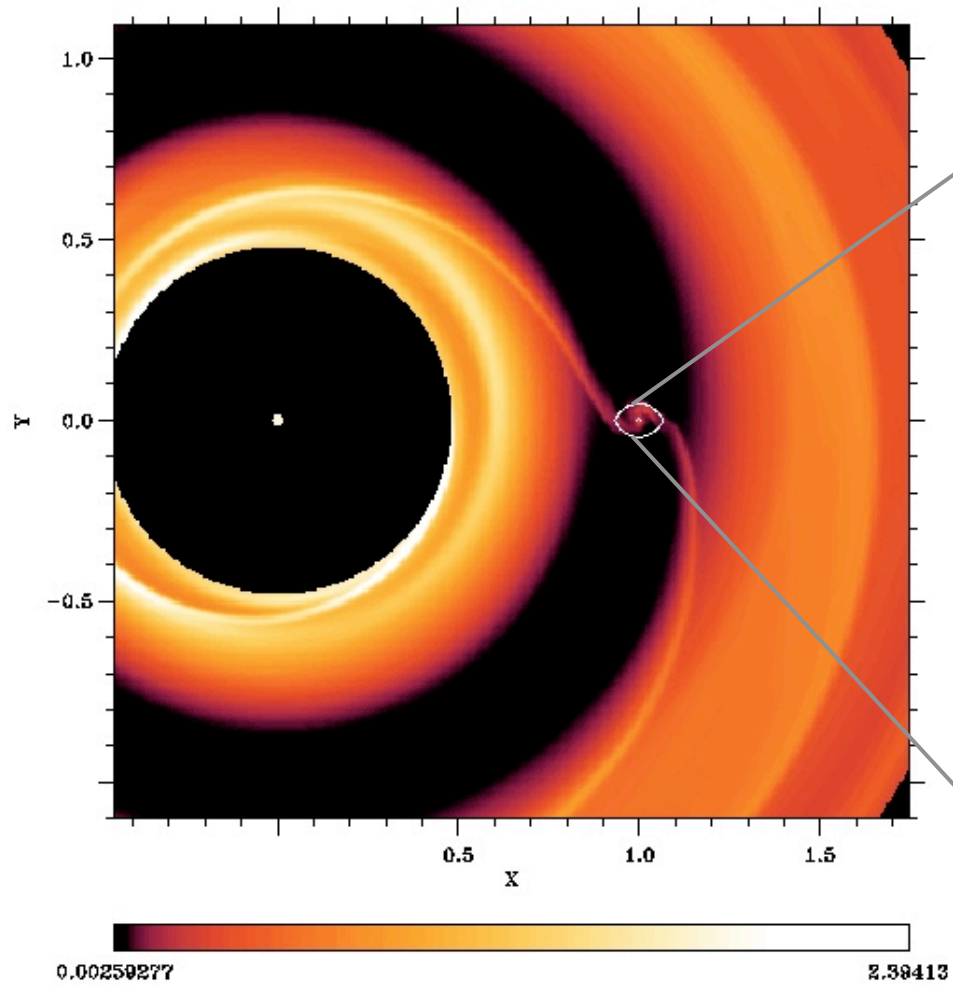


Lyra & Turner (2013, in prep)

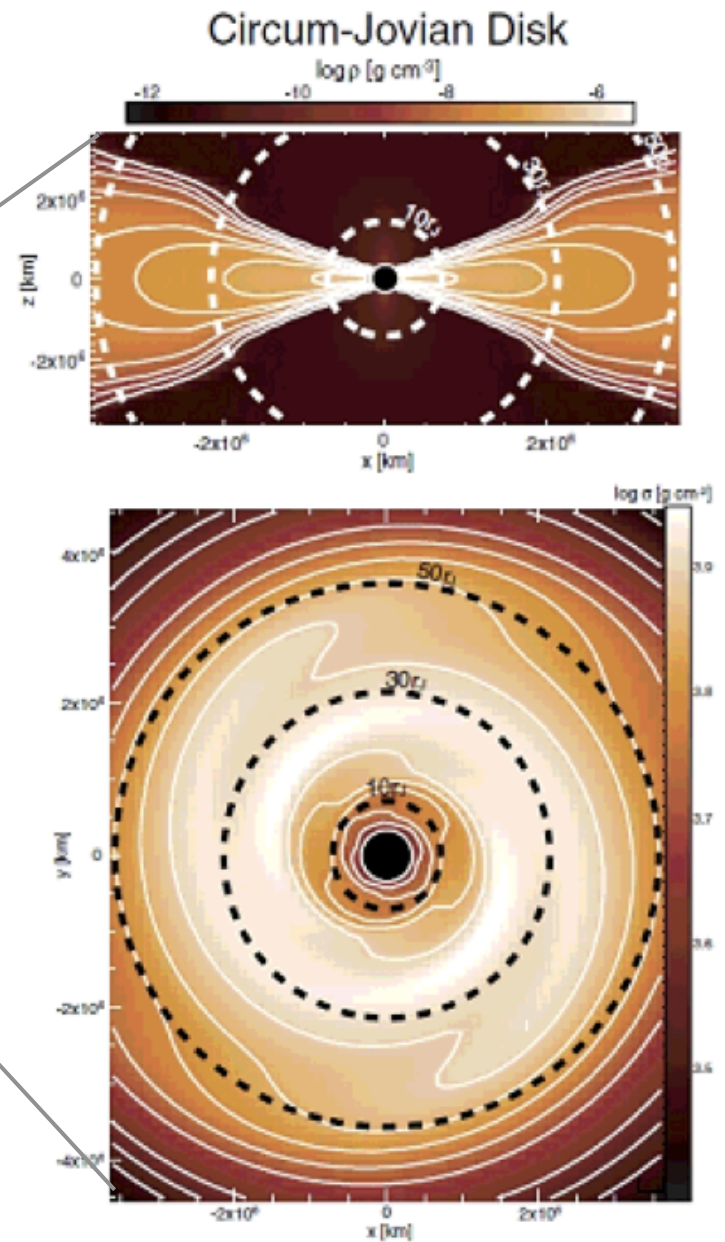
What is the mass spectrum?

Is there a strong mass-distance relationship?

Formation of Satellites



Credit: Pawel Artymowicz



Machida (2008)

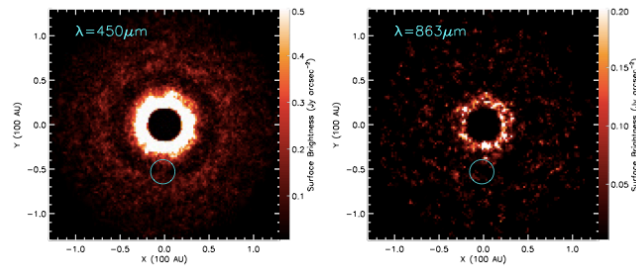
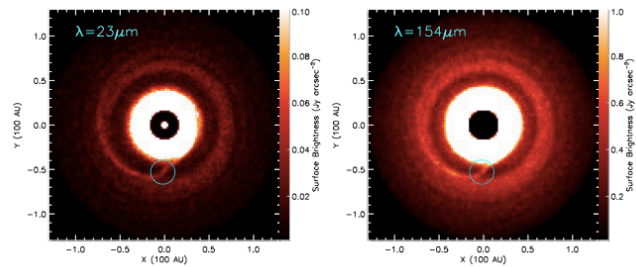
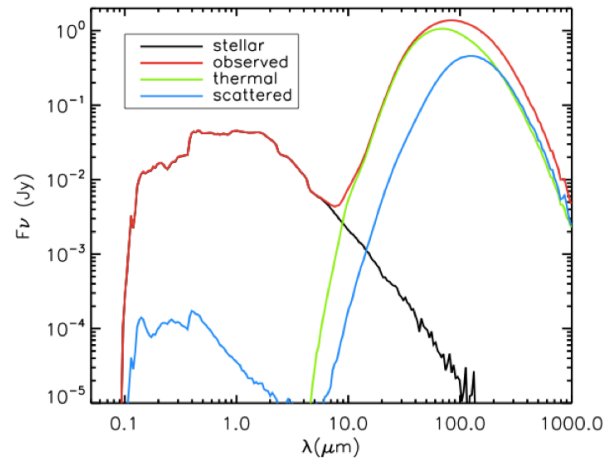
Formation of Satellites



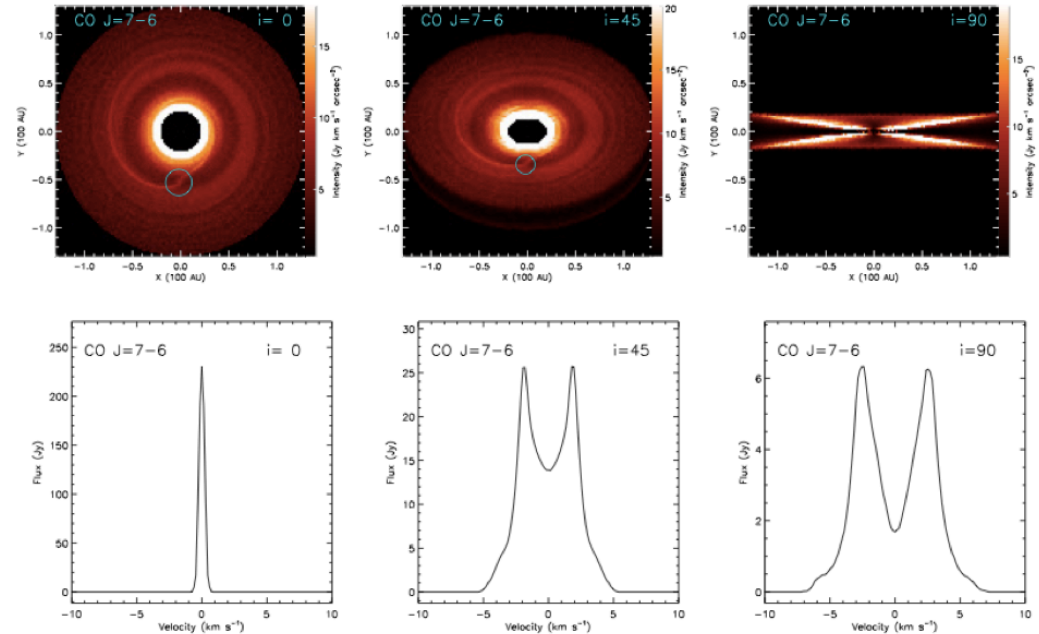
Credit: Pawel Ar

Machida (2008)

Combine full disk hydro models with Radiative Transfer for comparison with observations



Simulated observations
in far-IR and mm



Rovibrational lines of CO

(very) Preliminary results

Alex Richert (grad student, Penn State),
Yuexing Li (Penn State)

Summarizing

Gravitational collapse of an interstellar cloud

Outward transport of angular momentum through turbulence generated by the MRI. Dust coagulates into pebbles and boulders, sedimenting towards the midplane.

Rocks in the turbulent medium are trapped in transient pressure maxima and undergo collapse into planetesimals and dwarf planets.

Vortices may be excited in the dead zone. Inside them, the first dozens of Mars-mass embryos are formed. IMF ~ -2

Opacity transitions develop into regions of convergent migration. Low mass planets converge to these zones by inward/outward migration.

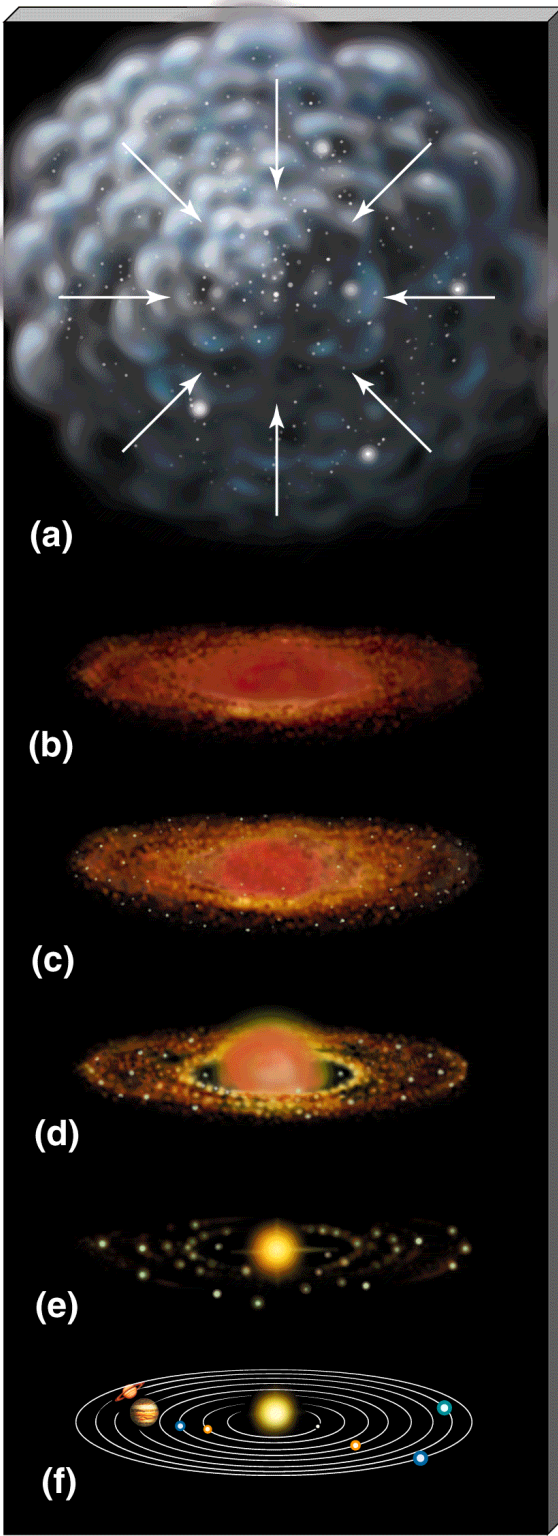
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N-body interactions and stochastic forcing during disk evaporation produce the system's final architecture.

Debris disks with gas are subject to a thermo-centrifugal instability

The instability generates sharp eccentric rings. Caution before shouting "planet!". Not all that glitters is gold.



Summarizing

Gravitational collapse of an interstellar cloud

Outward transport of angular momentum through turbulence generated by the MRI. Dust coagulates into pebbles and boulders, sedimenting towards the midplane.

Rocks in the turbulent medium are trapped in transient pressure maxima and undergo collapse into planetesimals and dwarf planets.

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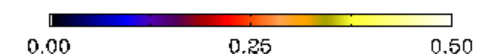
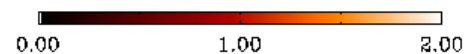
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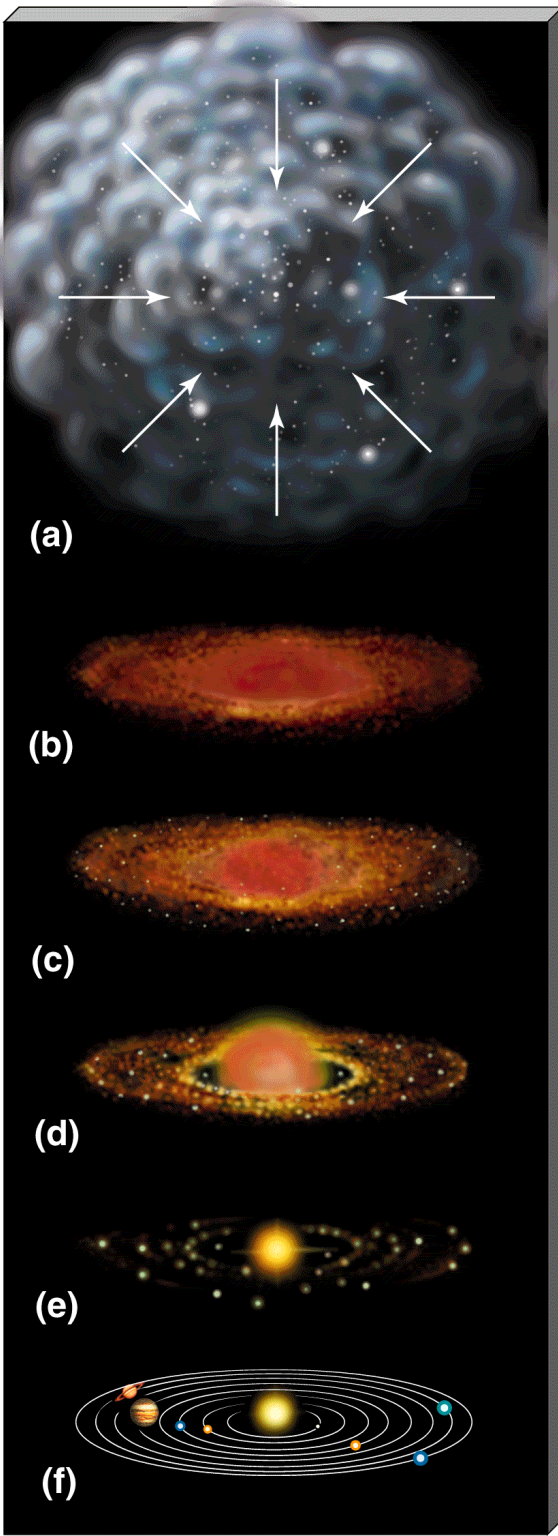
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Vortices may be excited in the dead zone. Inside them, the first dozens of Mars-mass embryos are formed. IMF ~ -2

Opacity transition
converge to the

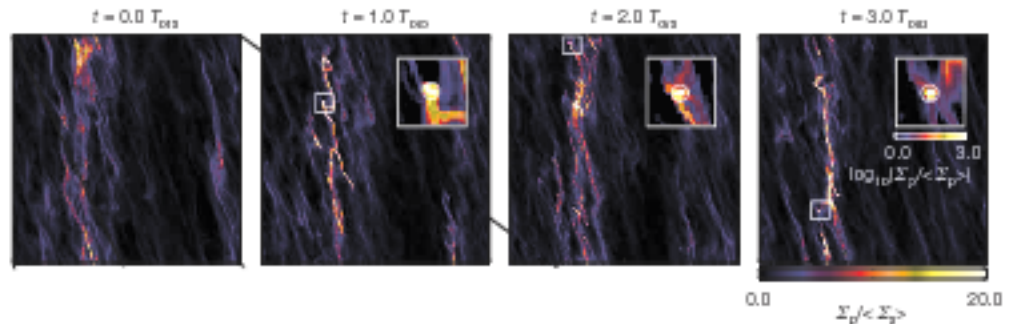
Convergent migration
forcing. Collisions

The disk thins due

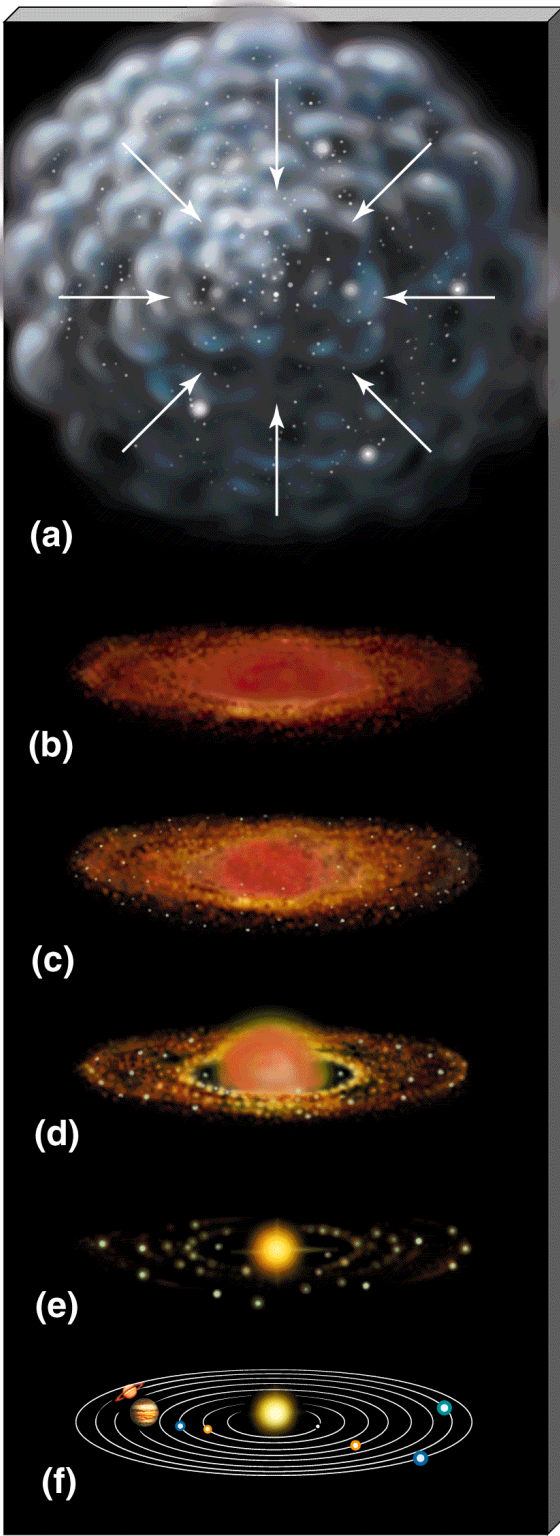
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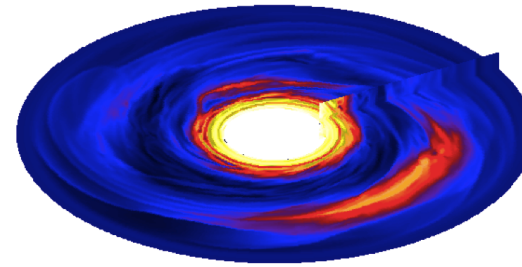
$t=22.28 \tau_0$

Summarizing

Gravitational collapse of an

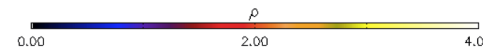
Outward transport of angular momentum. Dust coagulates into particles.

Rocks in the turbulent medium undergo collapse into planetesimals.



are generated by the differential rotation. Gas is drawn towards the midplane.

Pressure maxima and minima trap dust.



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Opacity transitions develop into regions of convergent migration. Low mass planets converge to the inner disk.

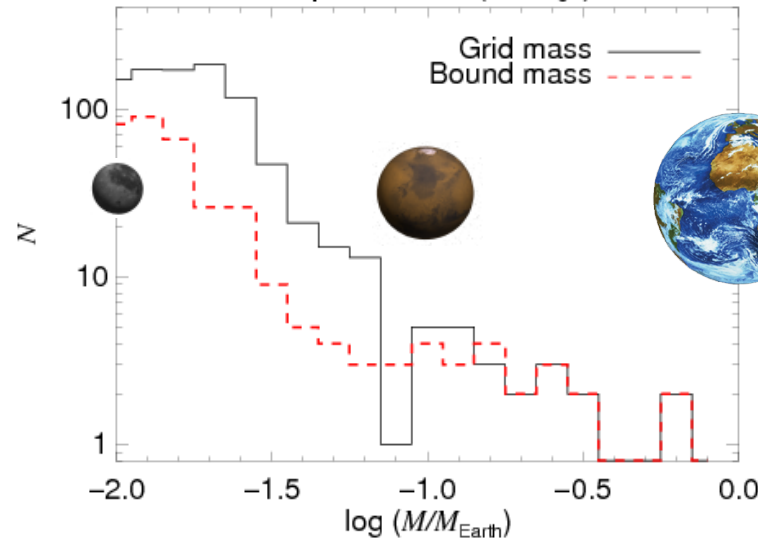
Convergent migration and collisional growth.

The disk thins and the gas is depleted.

N-body interactions sculpt the system's final configuration.

Debris disks with scattered planetesimals.

Mass Spectrum $t/(2\pi\Omega_0^{-1})=200$



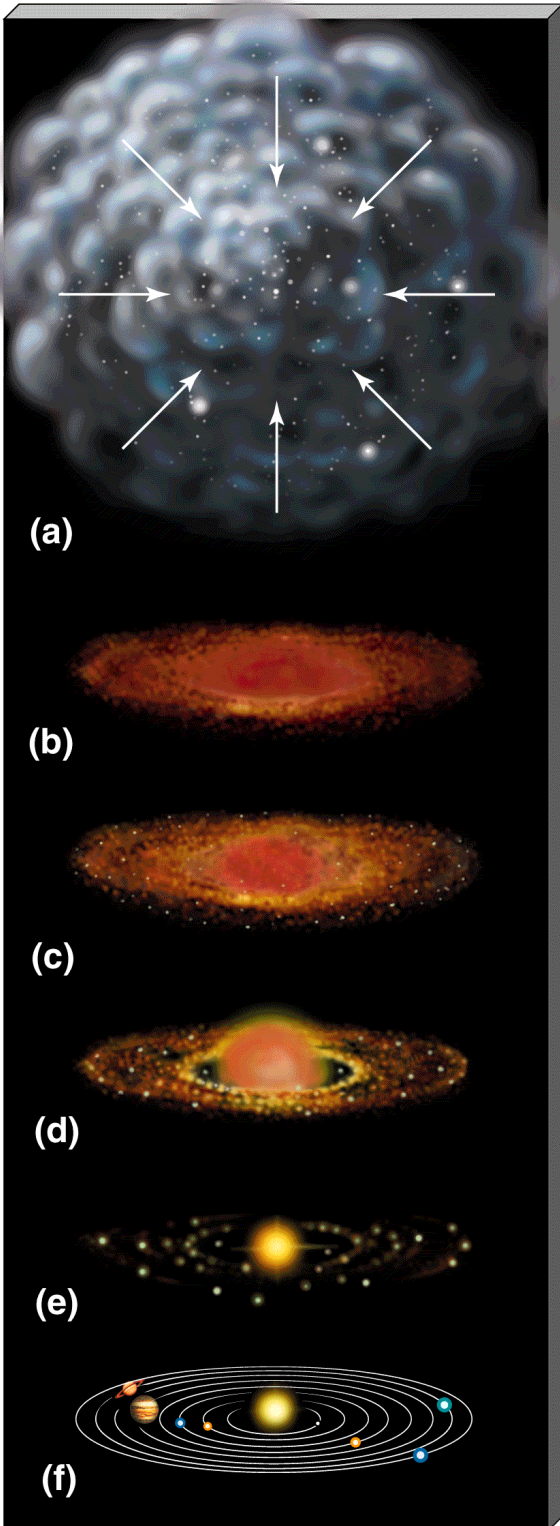
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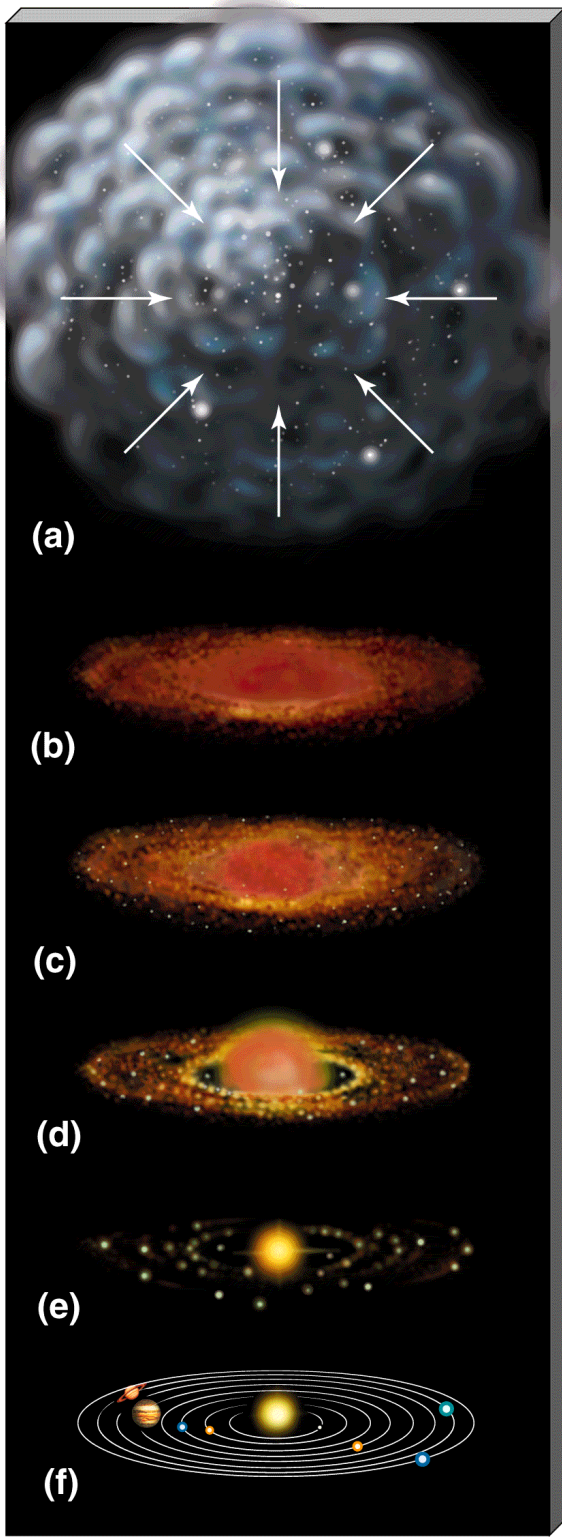
stable orbits.

migration produce the

stability

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Summarizing

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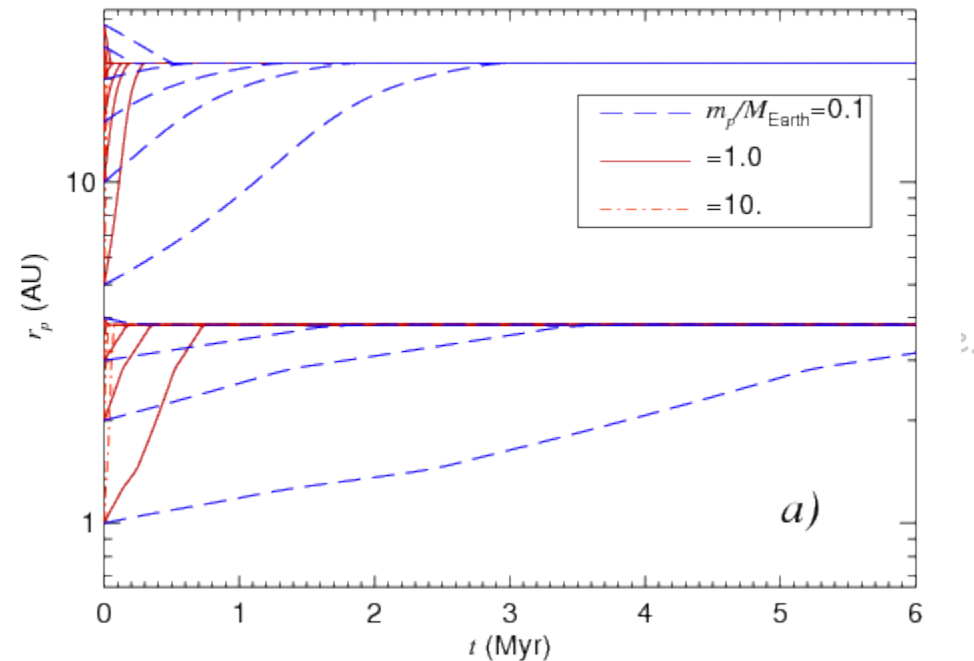
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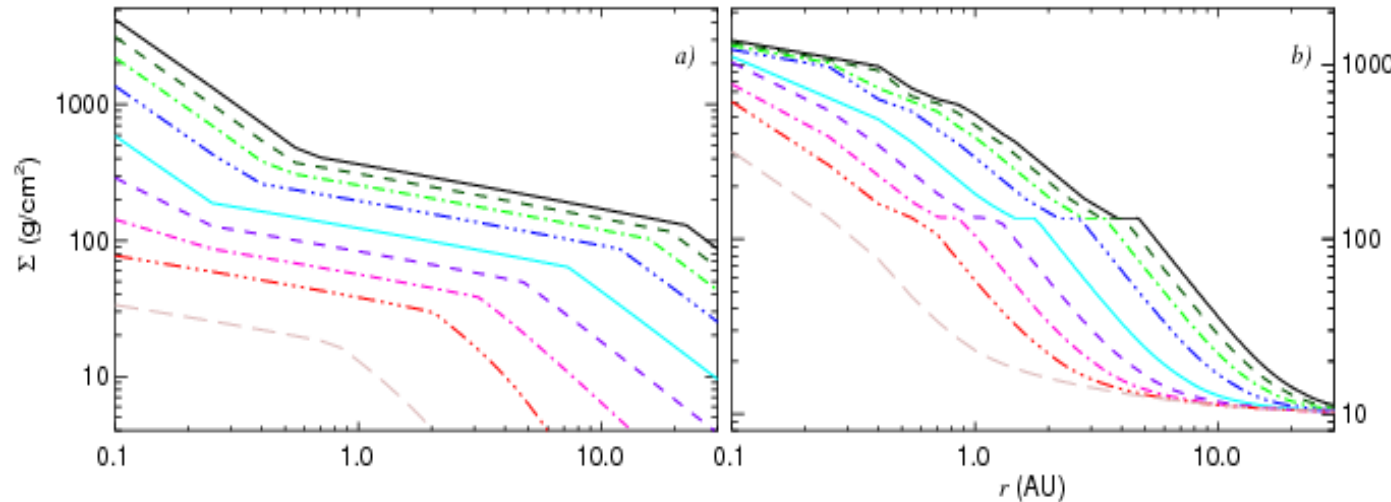
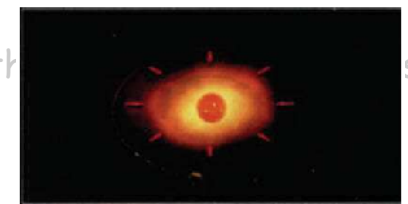
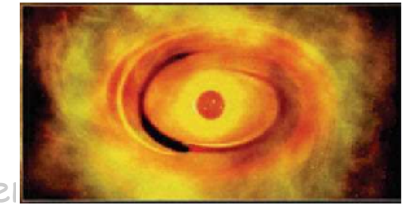
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Summarizing

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Outward transport of angular momentum through turbulence, MRI. Dust coagulates into pebbles and boulders, sediment



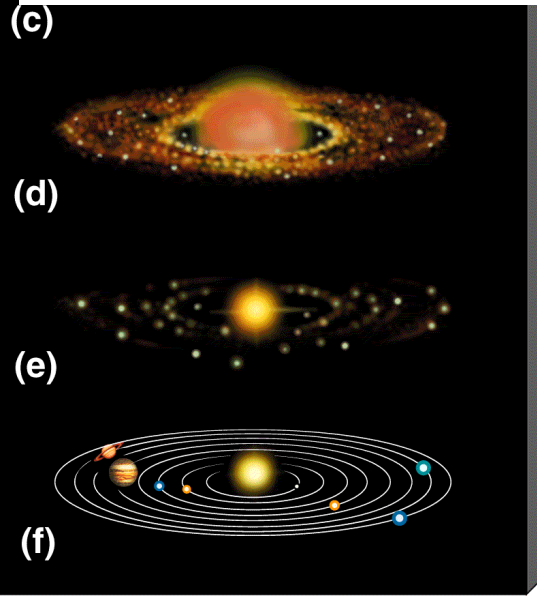
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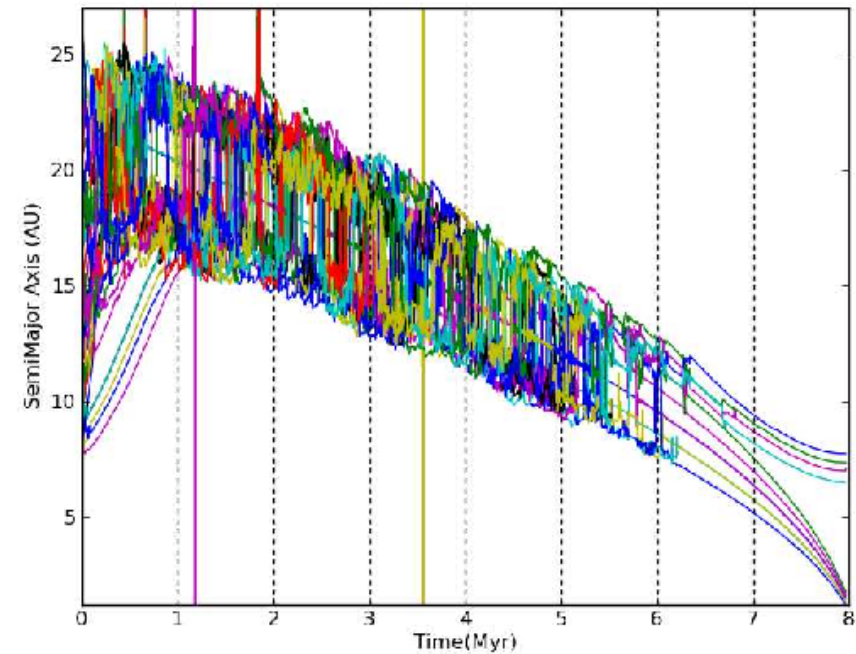
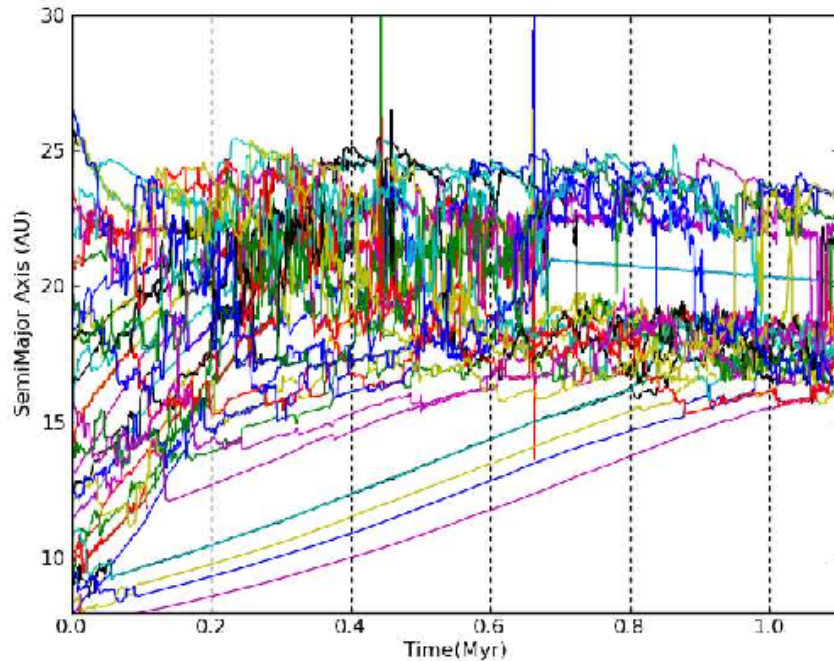
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Summarizing

Gravitational collapse of an interstellar cloud



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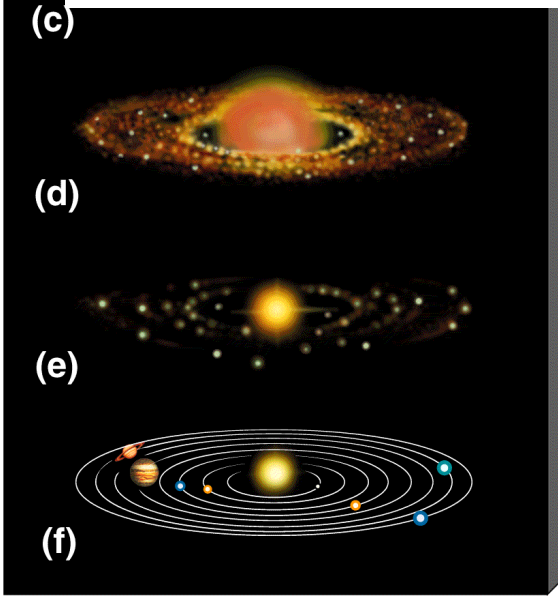
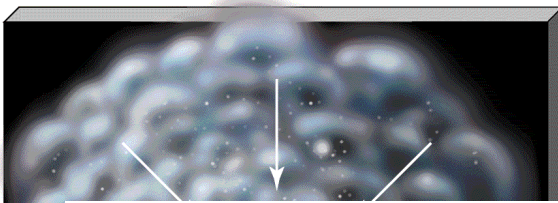
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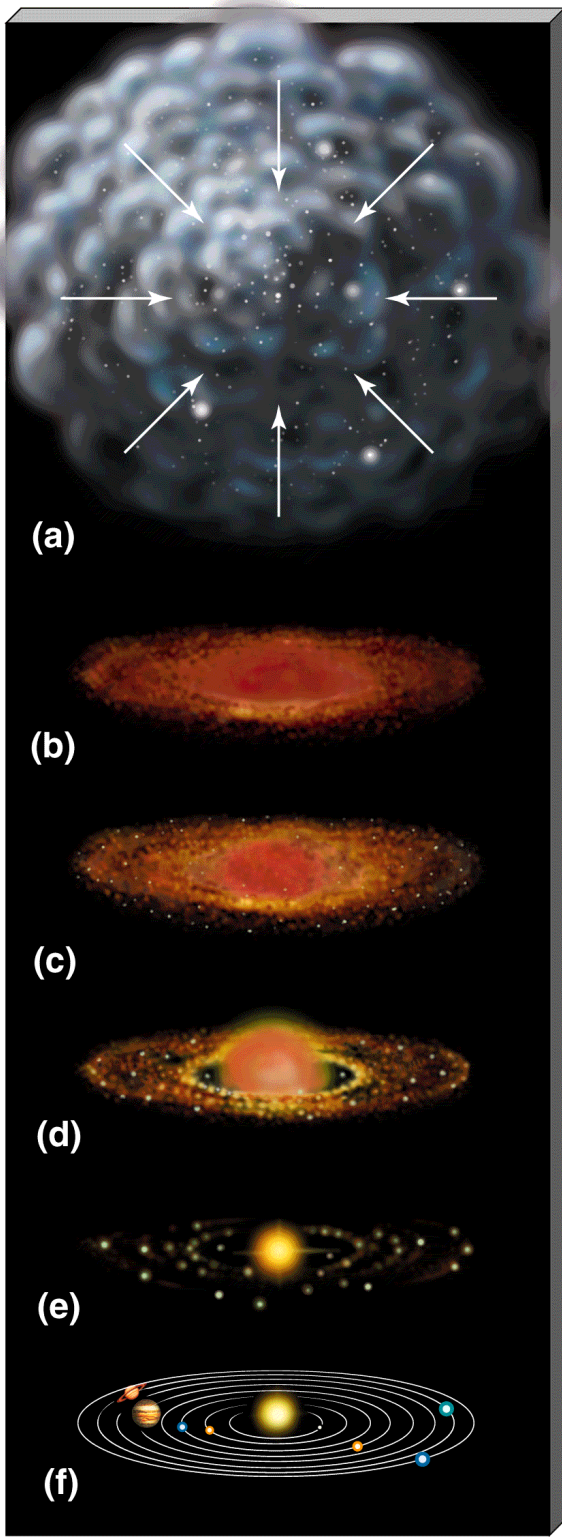
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Summariz

Gravitational collapse

Outward transport
MRI. Dust coagulation

Rocks in the turbulent
undergo collapse in

Vortices may be ex
mass embryos are

Dust heats gas
Heated gas = high pressure region
High pressure concentrates dust

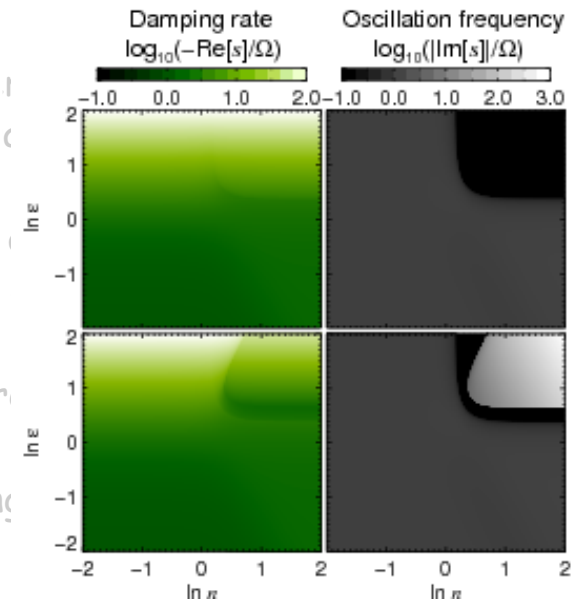
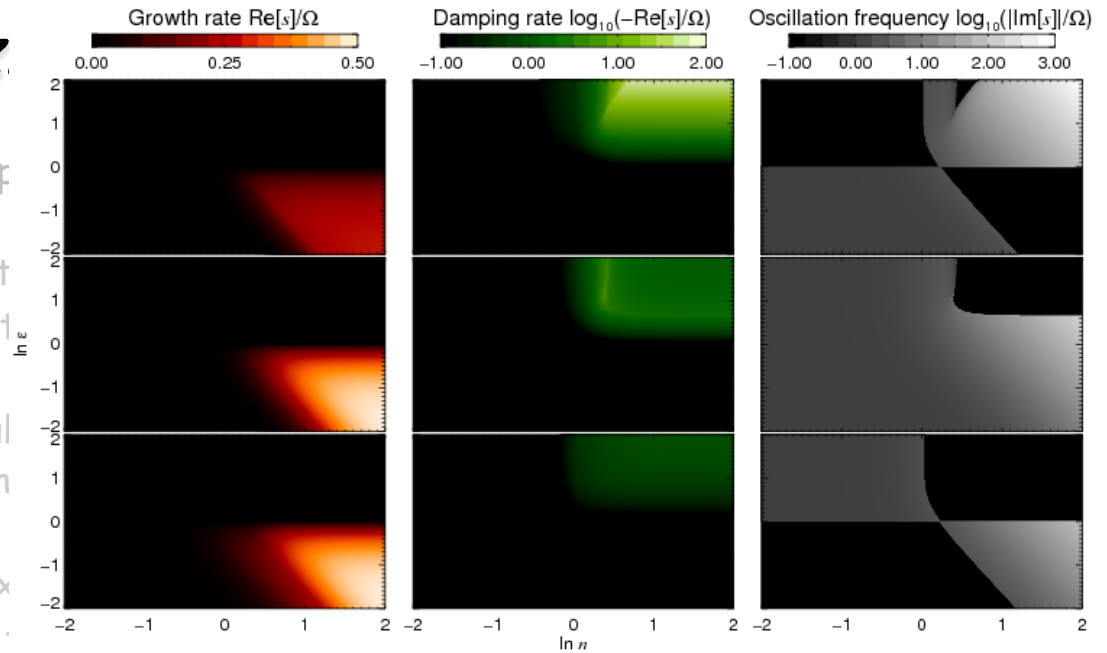
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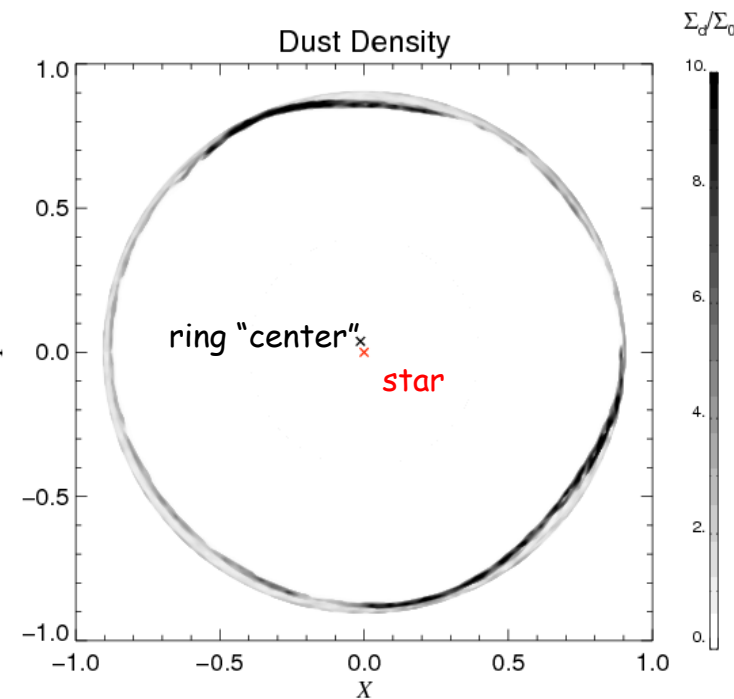
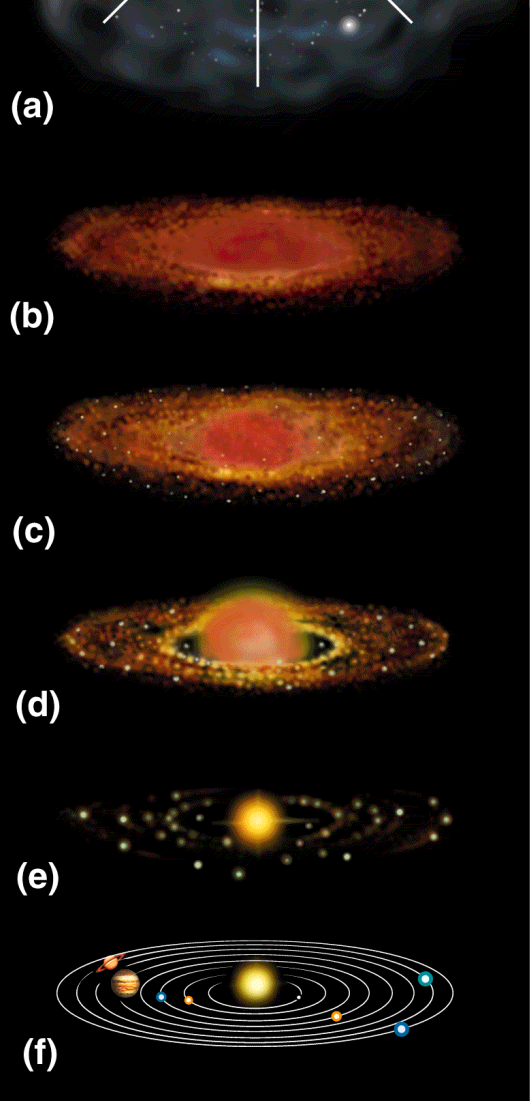
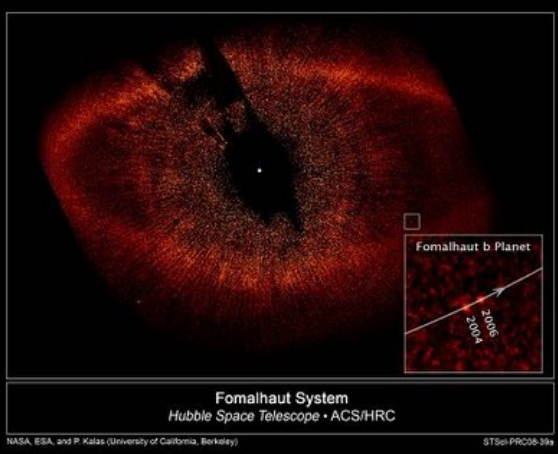
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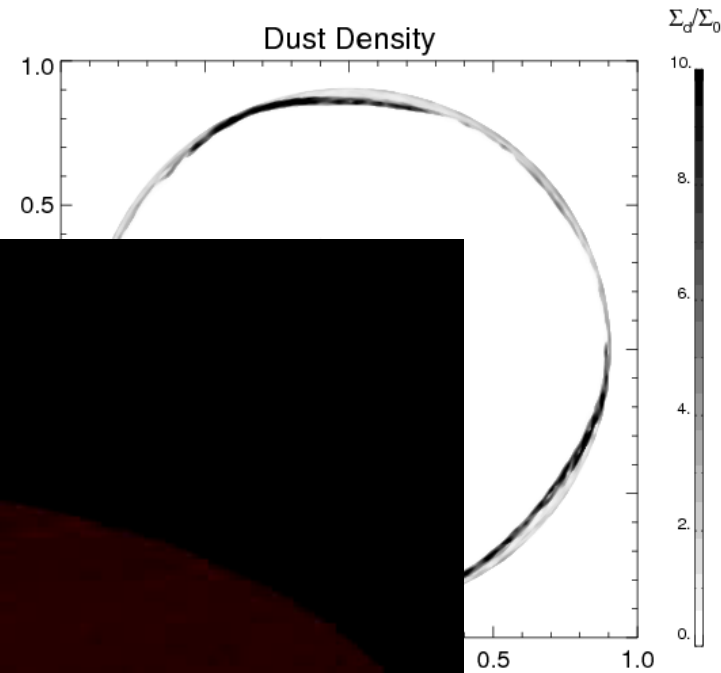
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Gravitational collapse of an interstellar c



(a)

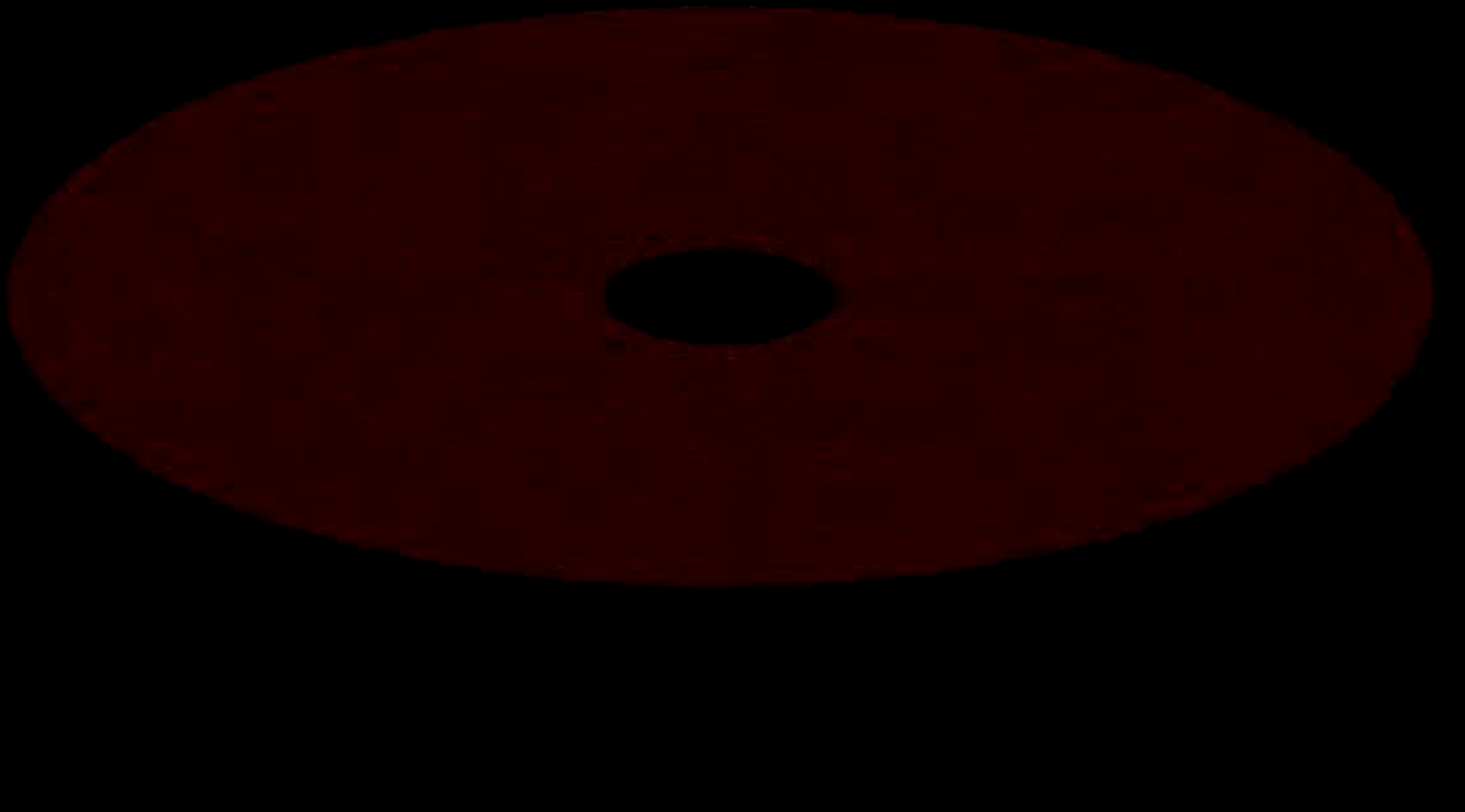
(b)

(c)

(d)

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