



# Evolution of Circumstellar Disks and Planet Formation



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(NASA Ames), Miguel de Val-Borro (NASA Goddard), Nienke van der Marel  
(IfA Hawaii), Andras Zsom (Brown).



Rio de Janeiro, Aug 10<sup>th</sup>, 2017

# Protoplanetary Disks



## PP disk fact sheet

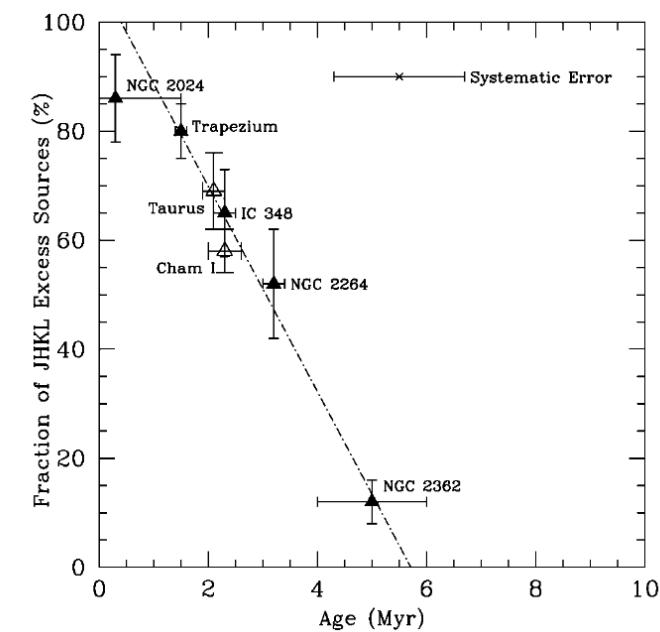
Density:  $10^{13} - 10^{15} \text{ cm}^{-3}$   
(Air:  $10^{21} \text{ cm}^{-3}$ )

Temperature: 10-1000 K

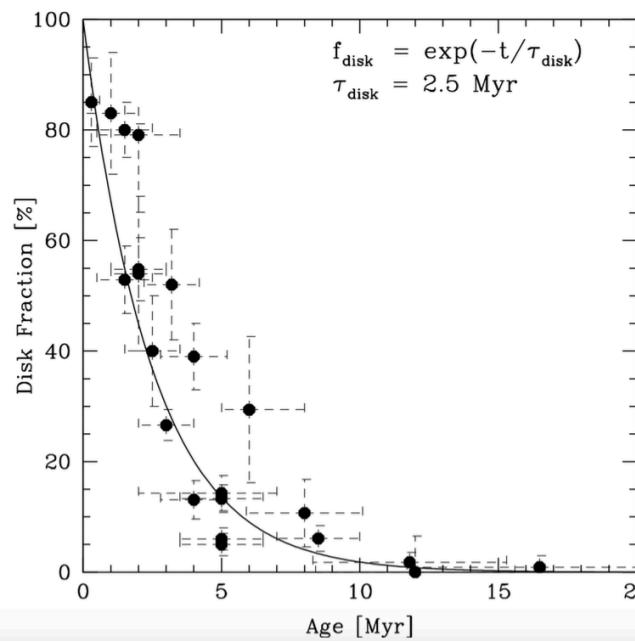
Scale: 0.1-100AU  
(1 AU =  $1.49 \times 10^{13} \text{ cm}$ )

Mass:  $10^{-3} - 10^{-1} M_{\text{sun}}$   
( $1 M_{\text{sun}} = 2 \times 10^{33} \text{ g}$ )

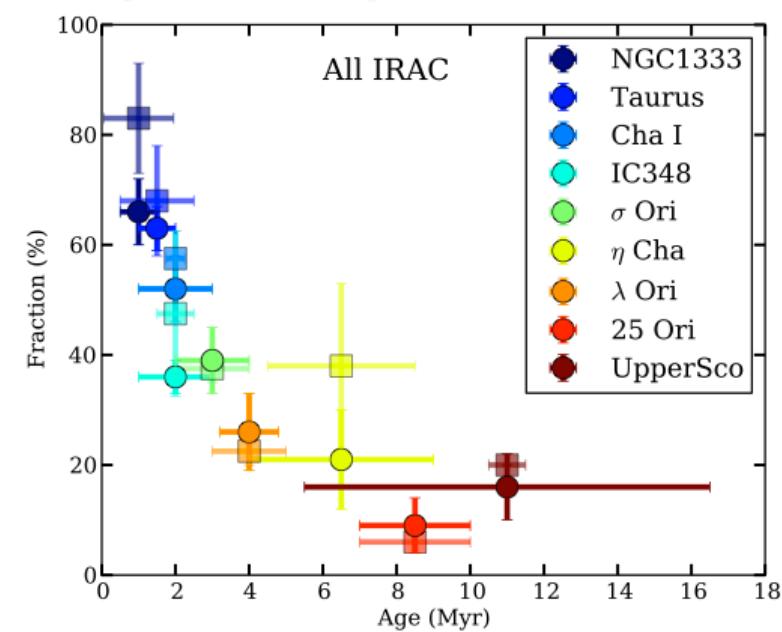
# Disk lifetime



(Haisch et al. 2001)

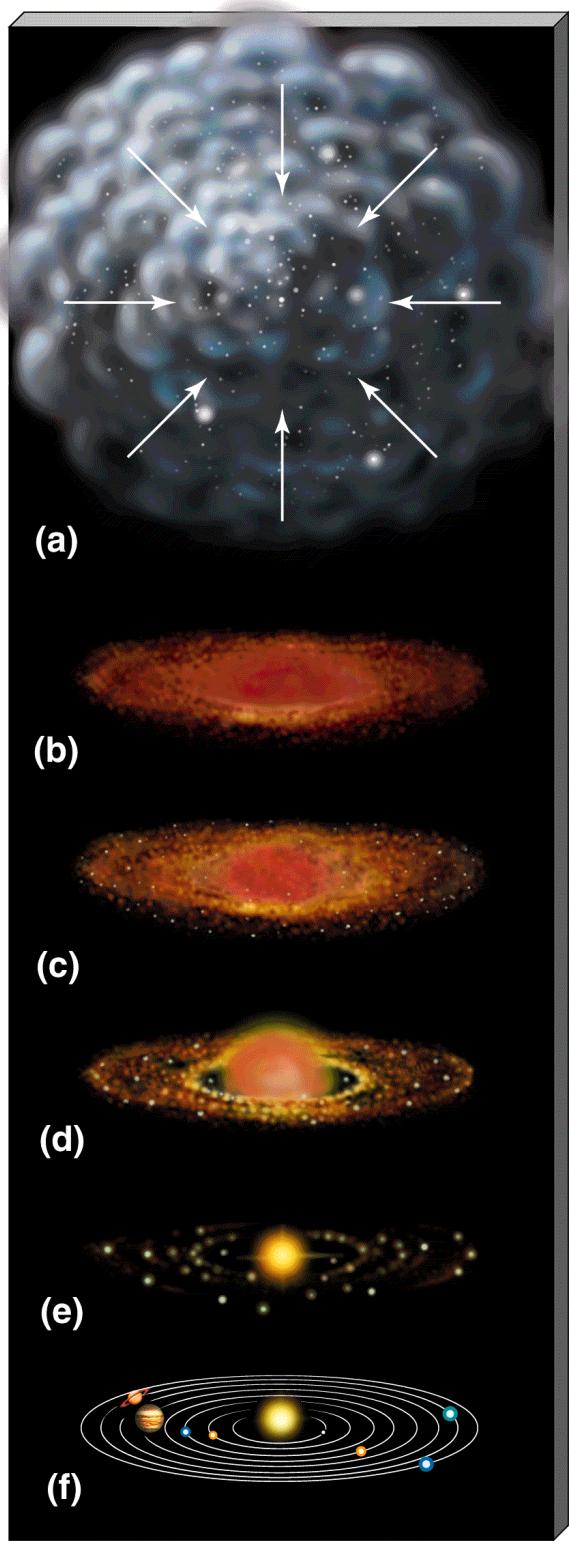


(Mamajek et al. 2009)



(Ribas et al. 2014)

Disks dissipate with an e-folding time of 2.5 Myr



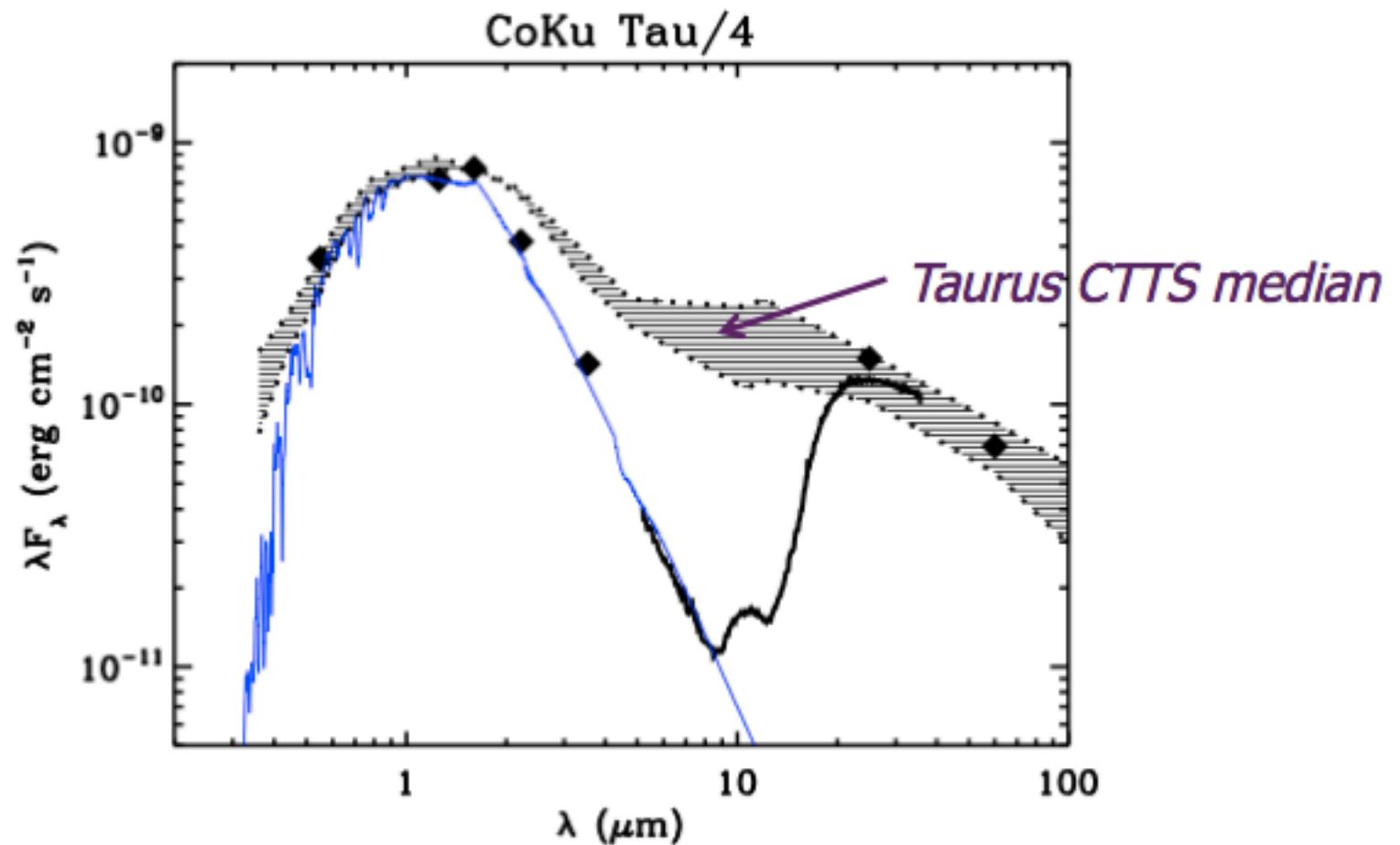
## Planet Formation

Gas-rich phase (< 10 Myr)  
*Primordial Disks*

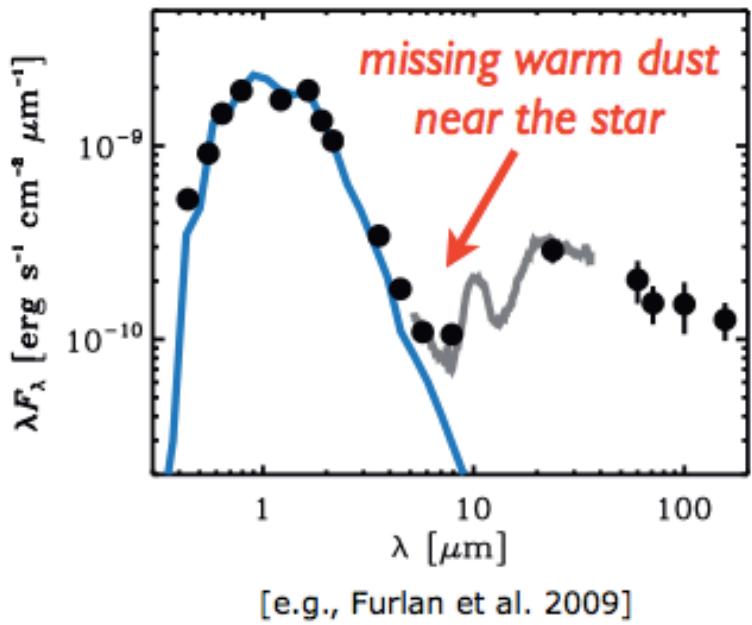
Thinning phase (~10 Myr)  
*Transition Disks (?)*

Gas-poor phase (>10 Myr)  
*Debris Disks*

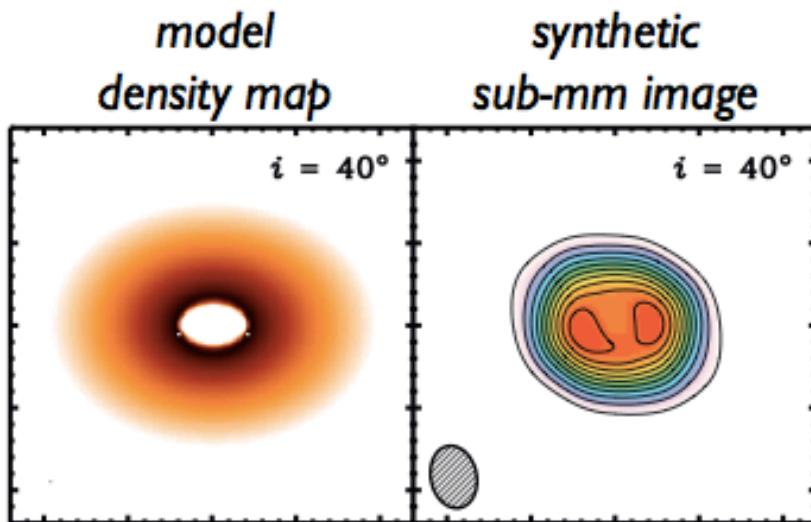
# Transition Disks: Disks with missing hot dust.



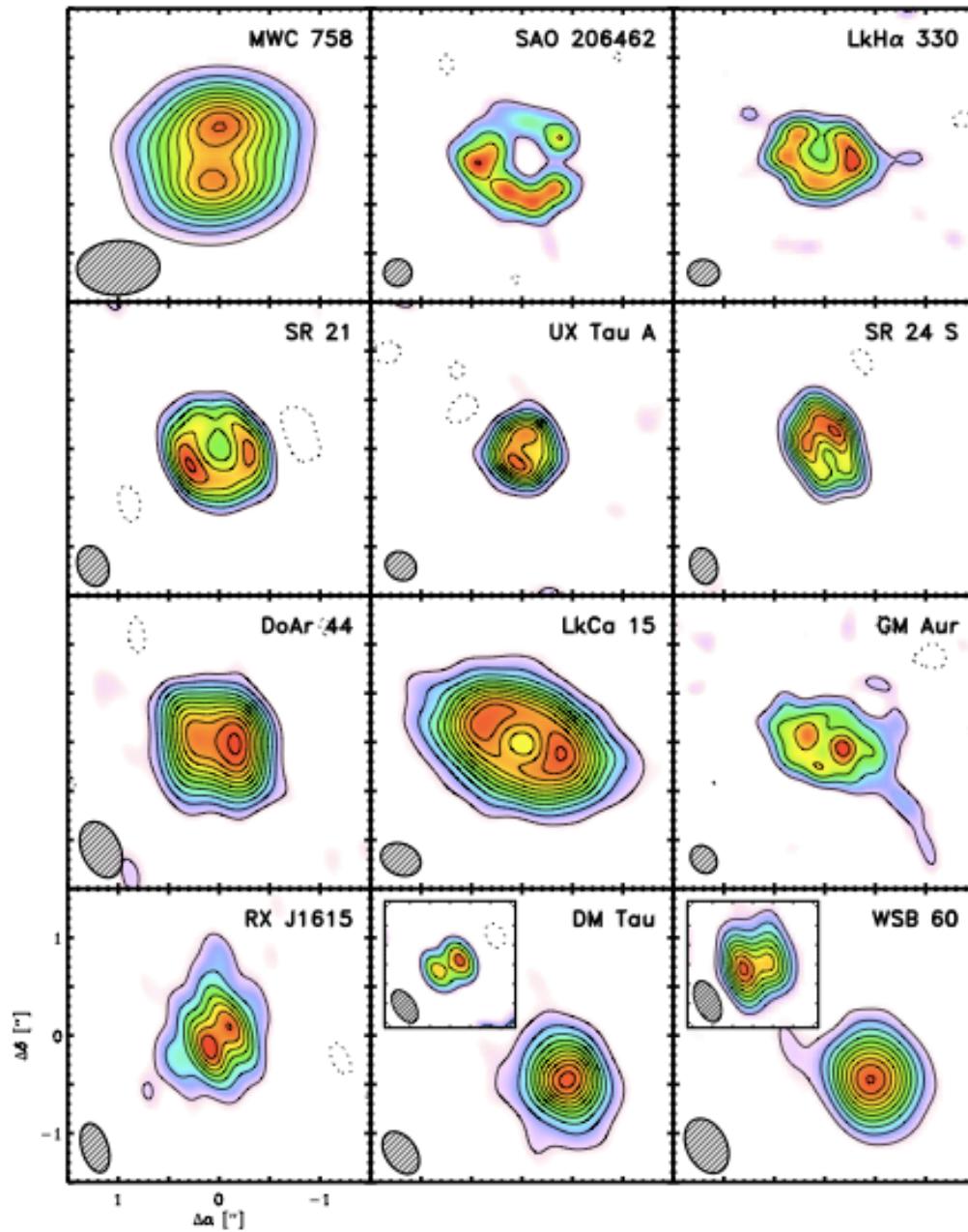
# Transition Disks: Disks with missing hot dust.



a disk with a large reduction  
in optical depth near the star  
(i.e., a “cavity” or “hole”)

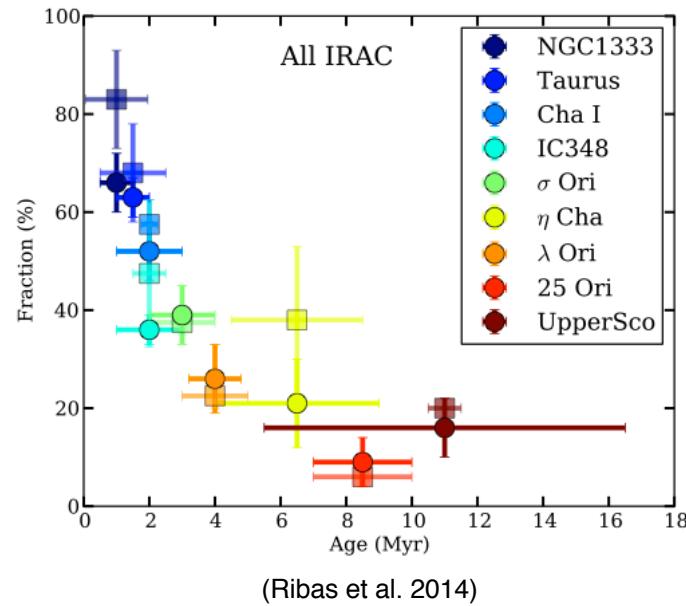


# Resolved transition disks with the Sub-millimeter Array (SMA)



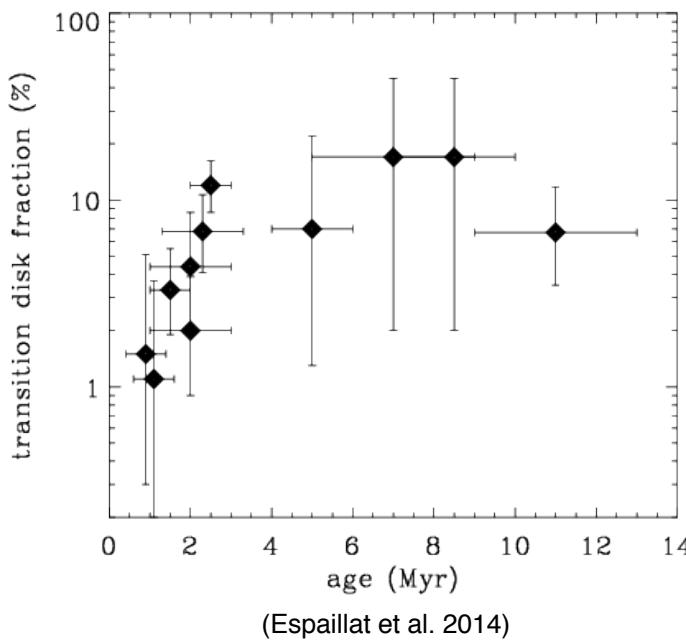
0.85mm  
0.3" ~ 20 AU resolution

# Are transition disks related to disk evolution?



(Ribas et al. 2014)

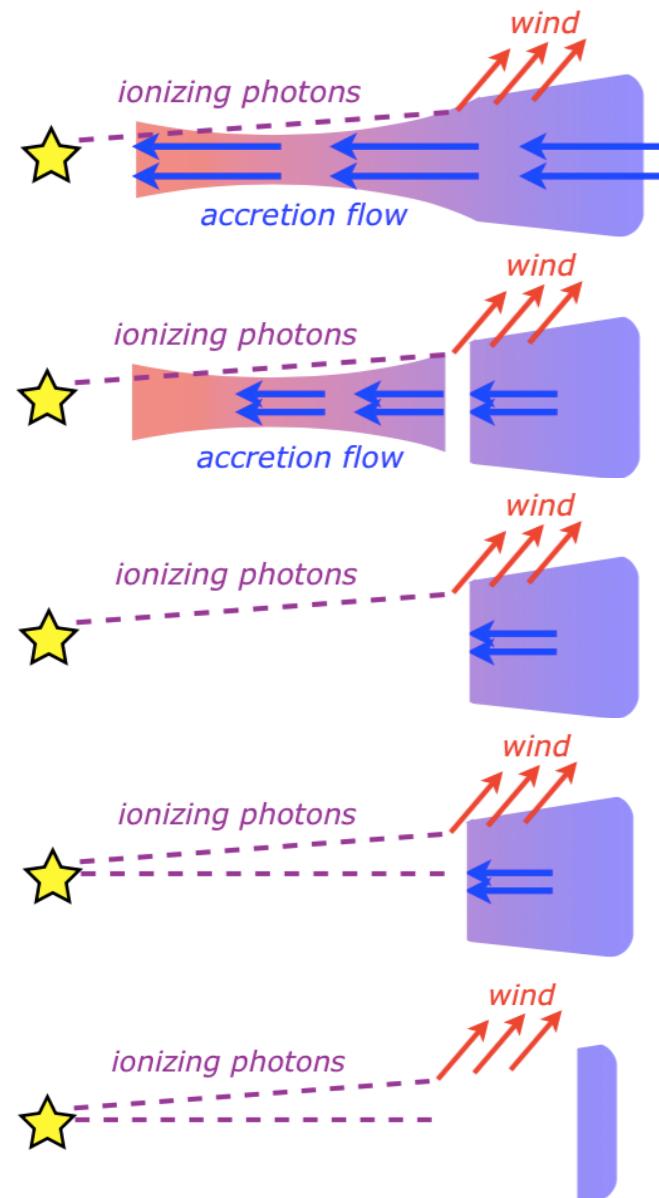
“Total” disk fraction



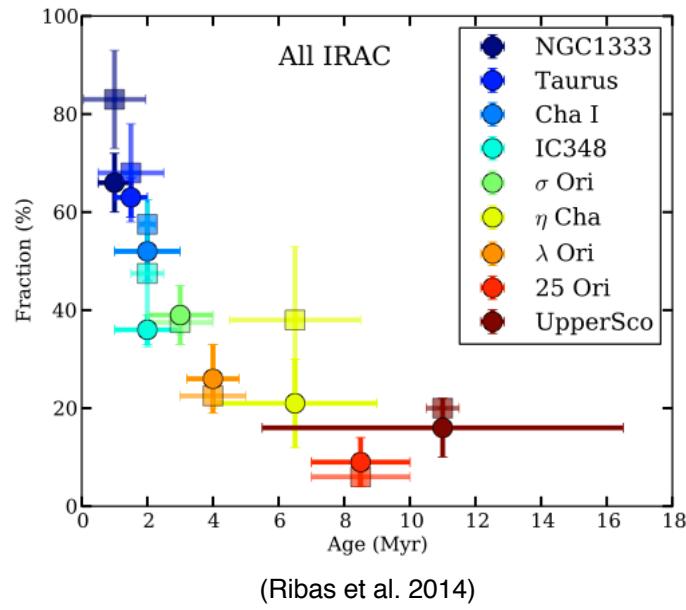
Transition disk fraction

(Espaillat et al. 2014)

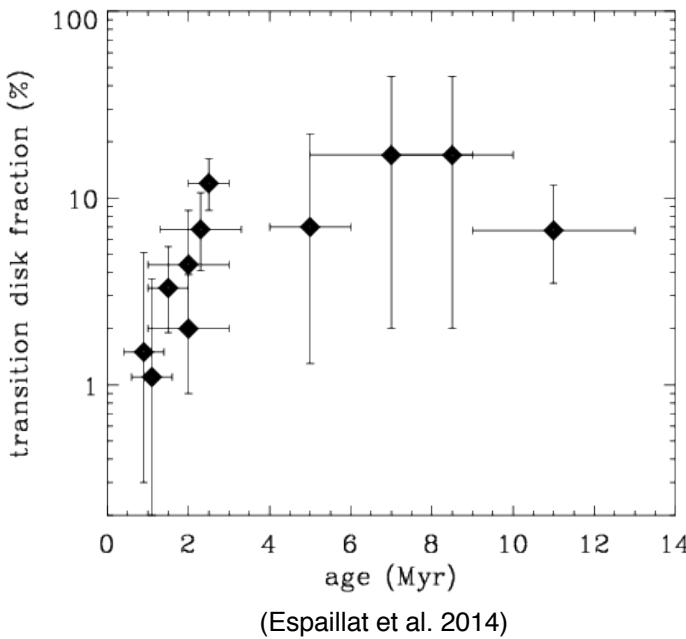
# Photoevaporation



# Look again...



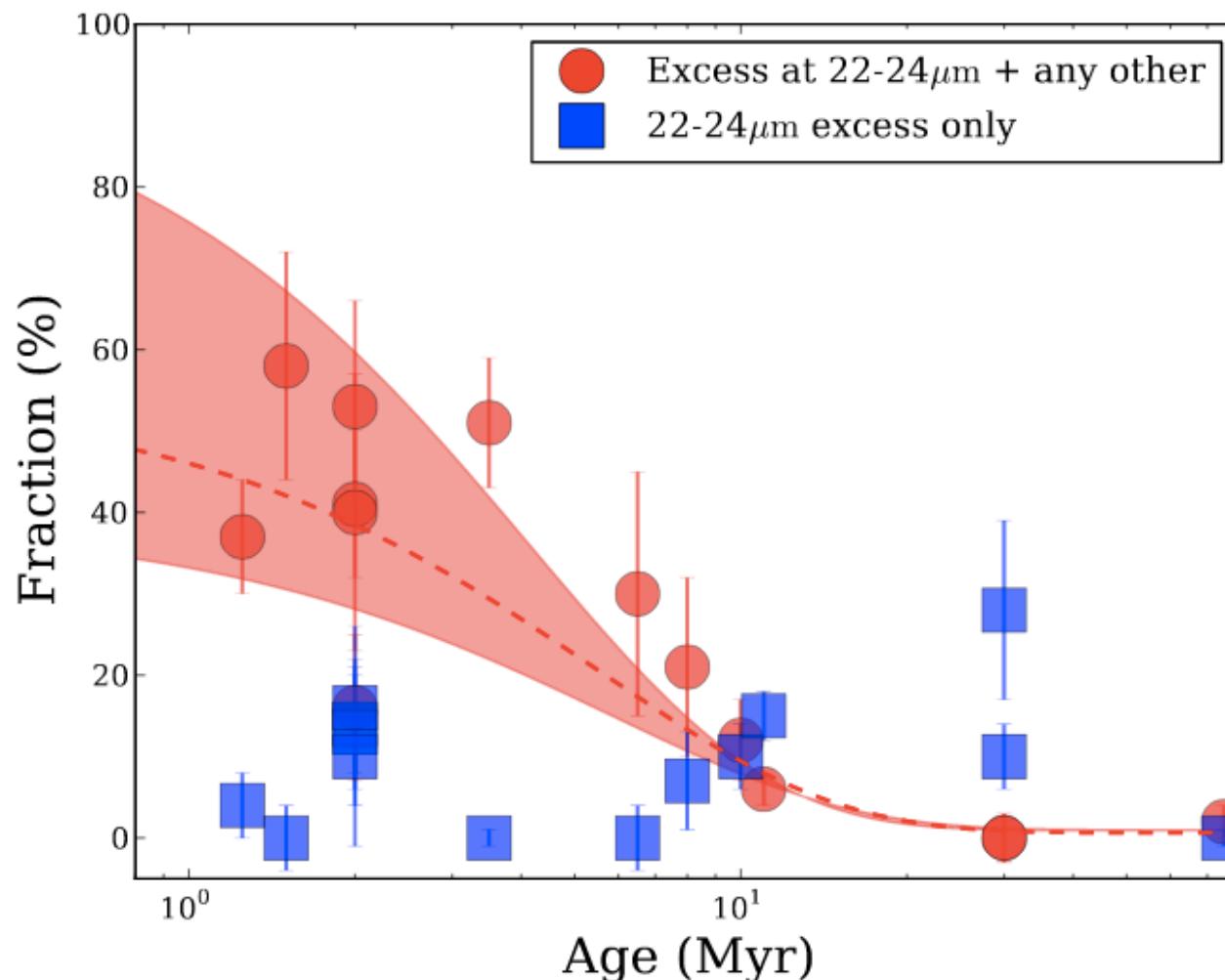
“Total” disk fraction



Transition disk fraction

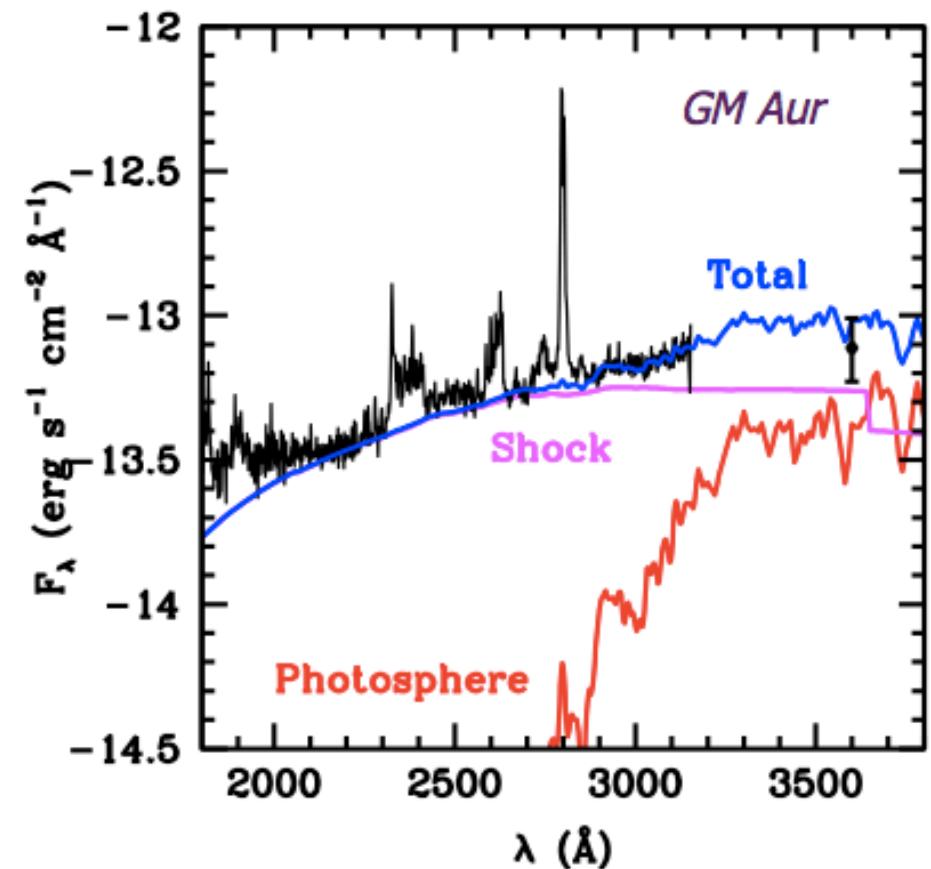
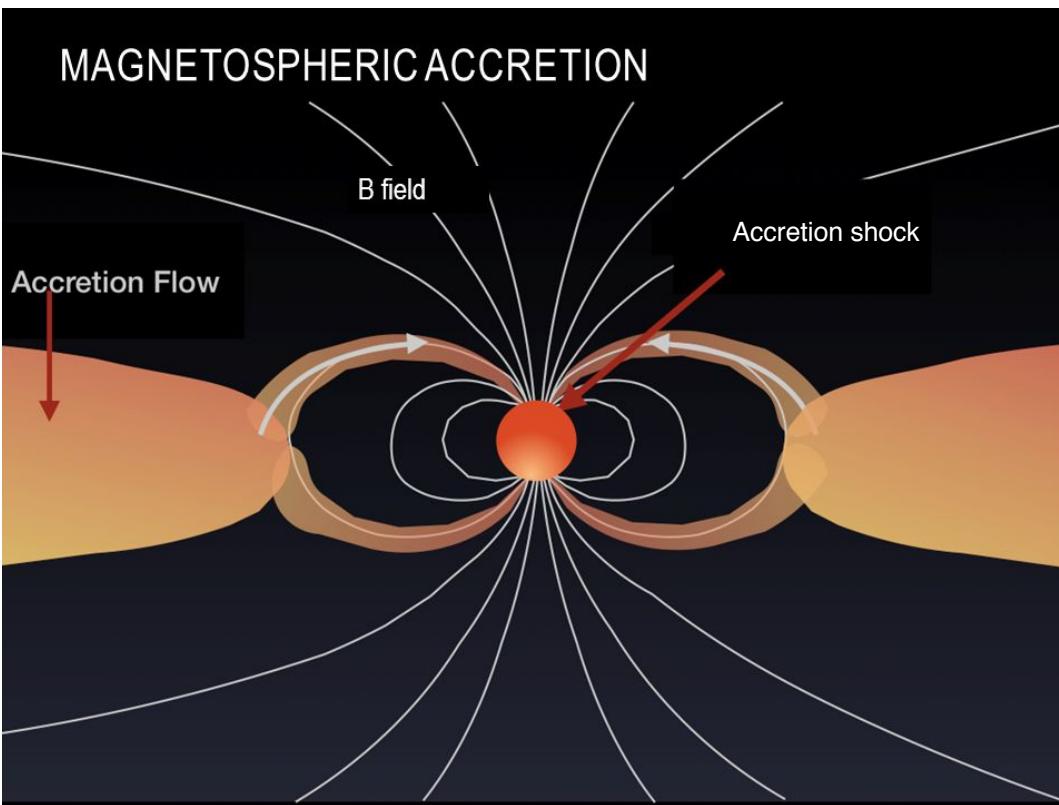
# Transition disks linked to disk evolution?

The distribution in age is consistent with a uniform distribution.

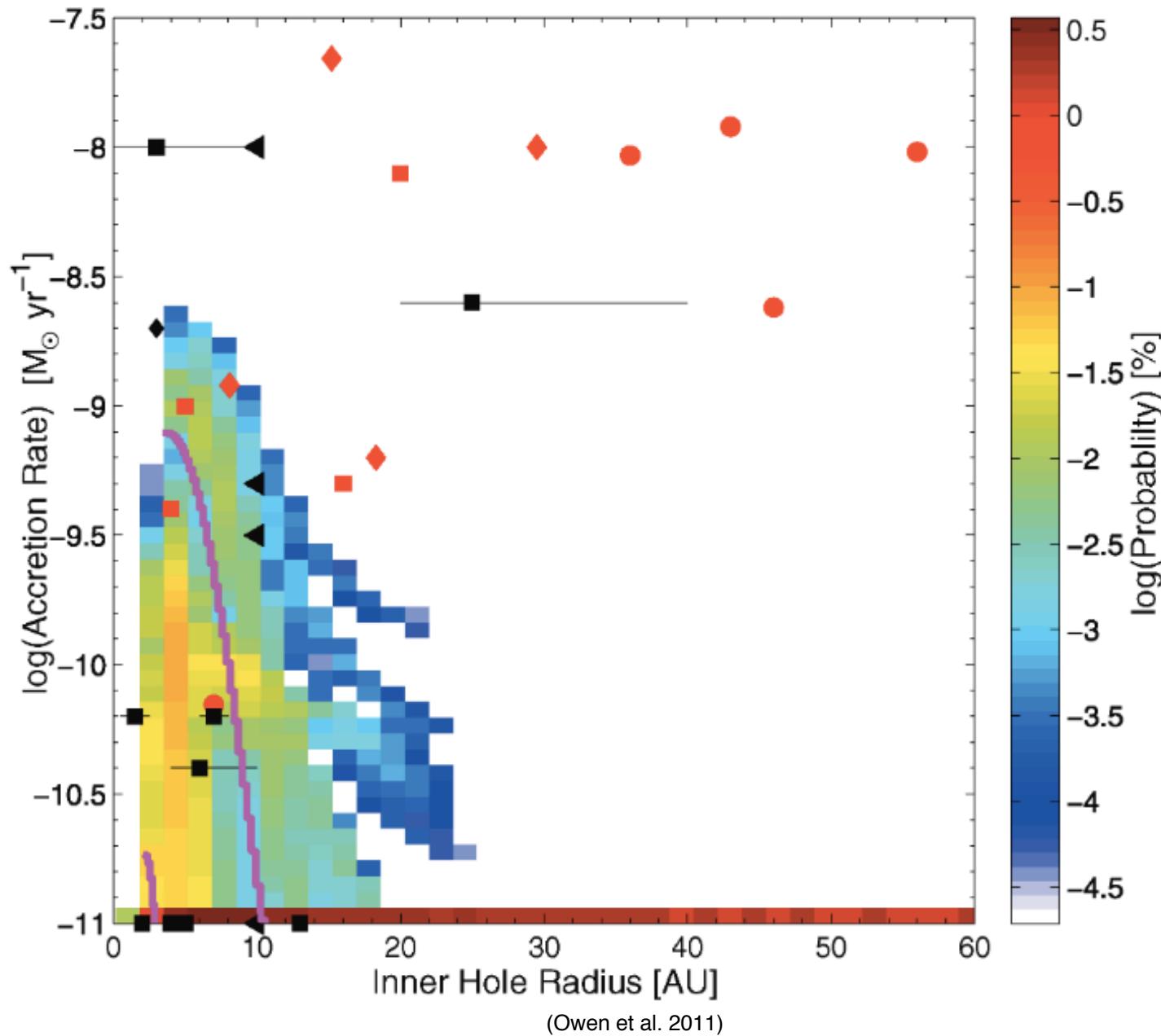


# UV excess

Many transitional disks show signs of accretion, at the level of primordial (classical T-Tauri) disks.



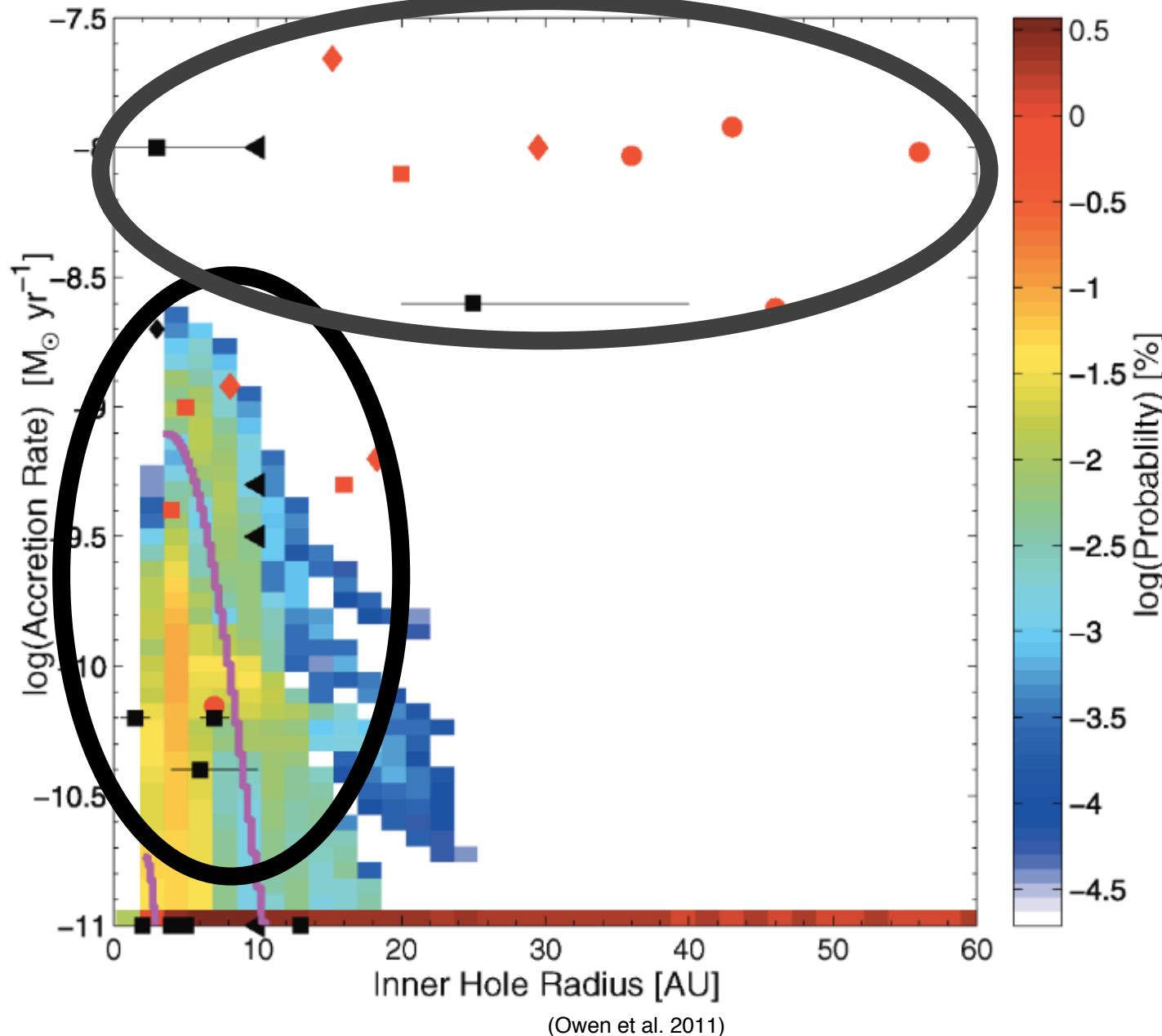
# Bimodal distribution of transition disks



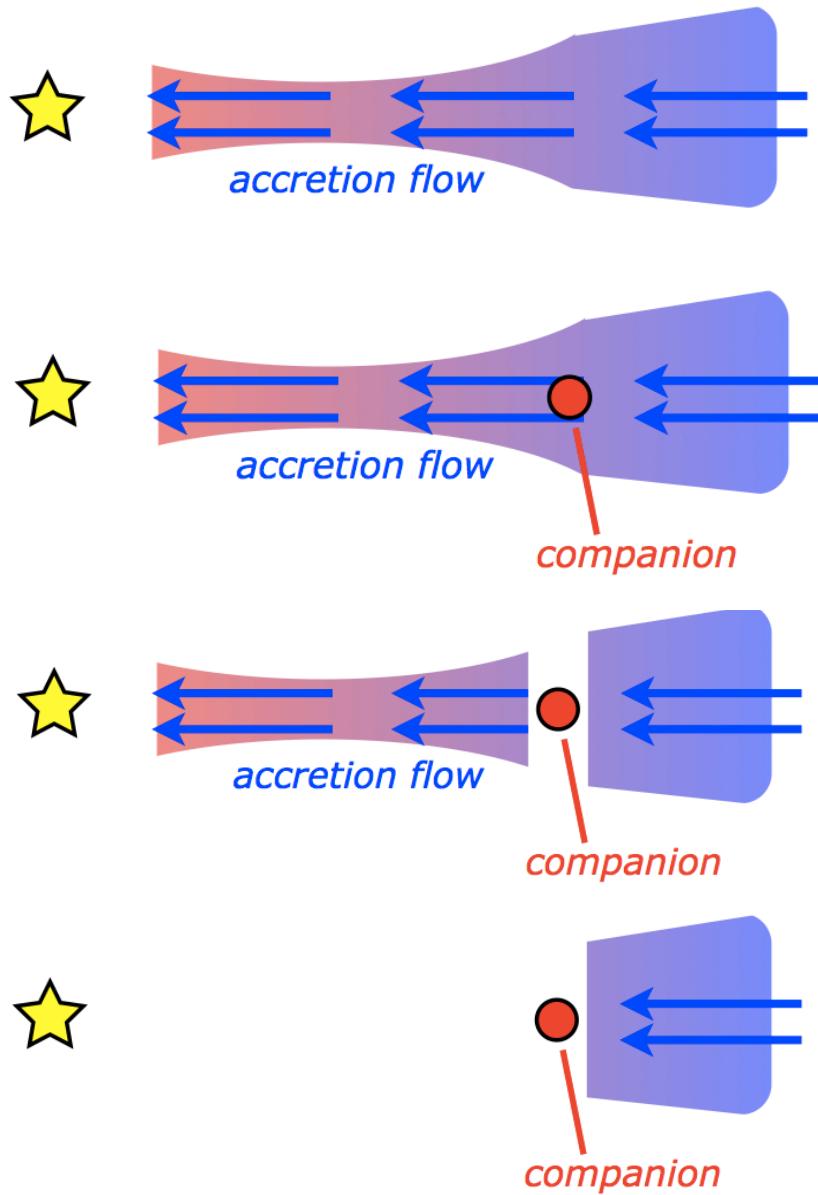
# Bimodal distribution of transition disks

*Not explained by photo-evaporation*

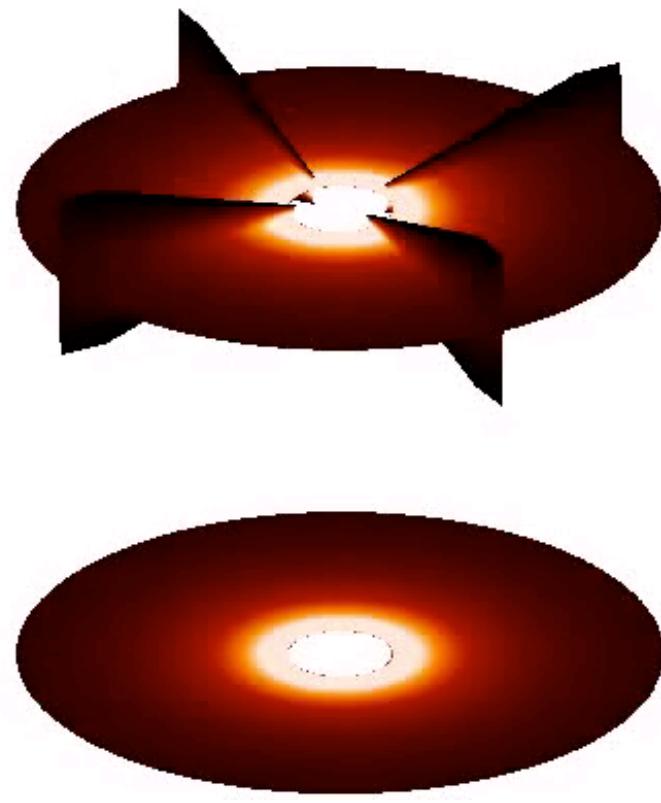
*Explained by  
Photo-  
evaporation*



# Planetary companion

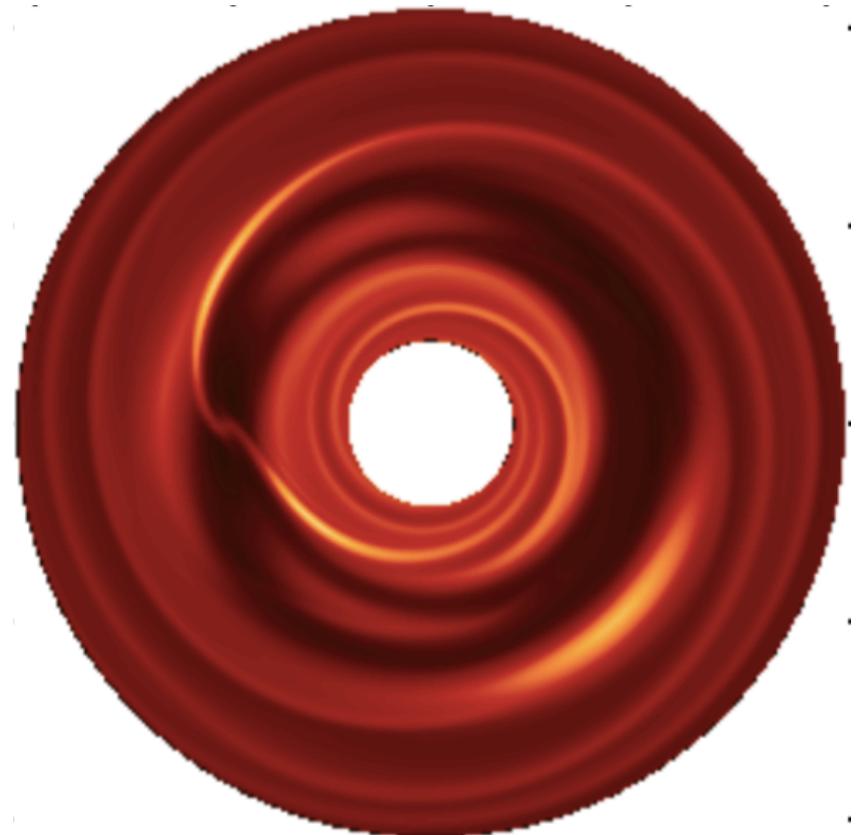
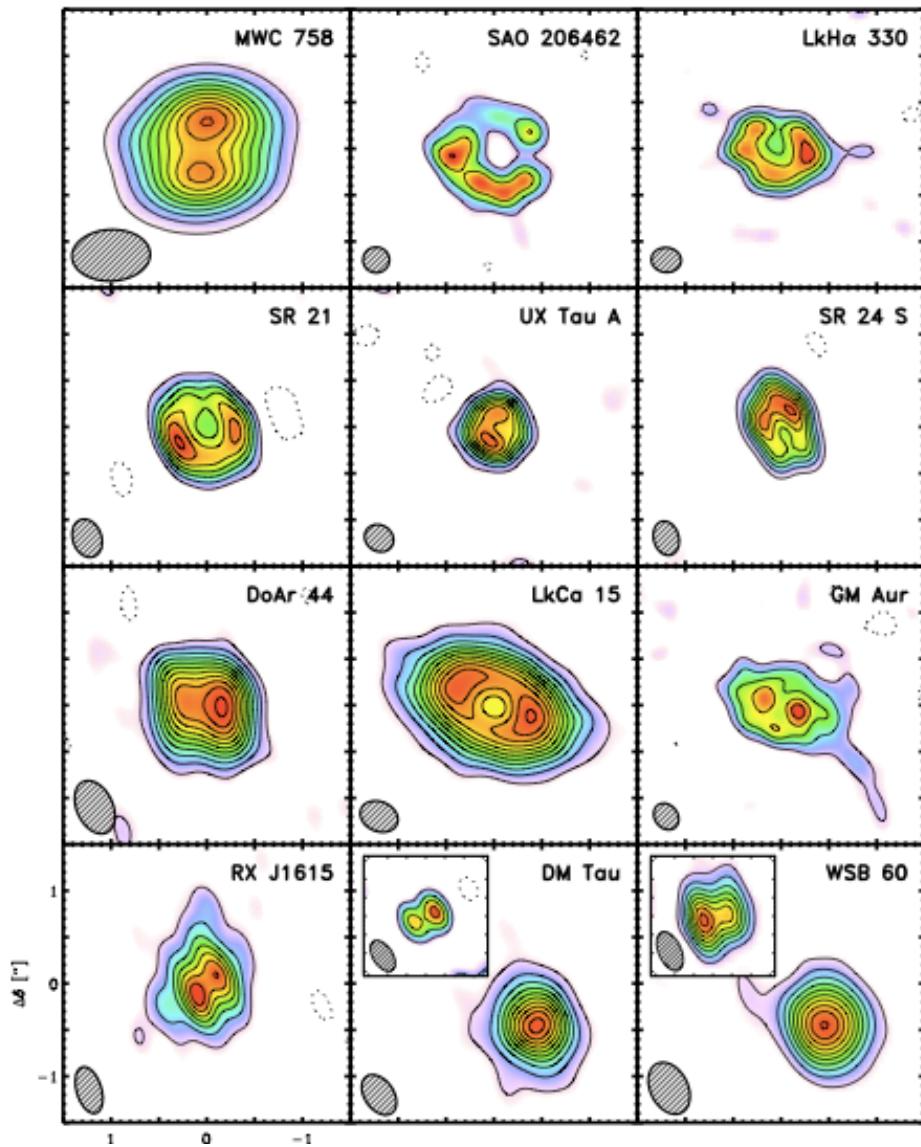


$t = 0.1$



(Lyra 2009)

# These cavities may be the telltale signature of forming planets



(Lyra et al. 2009b)

A way to directly study planet-disk interaction

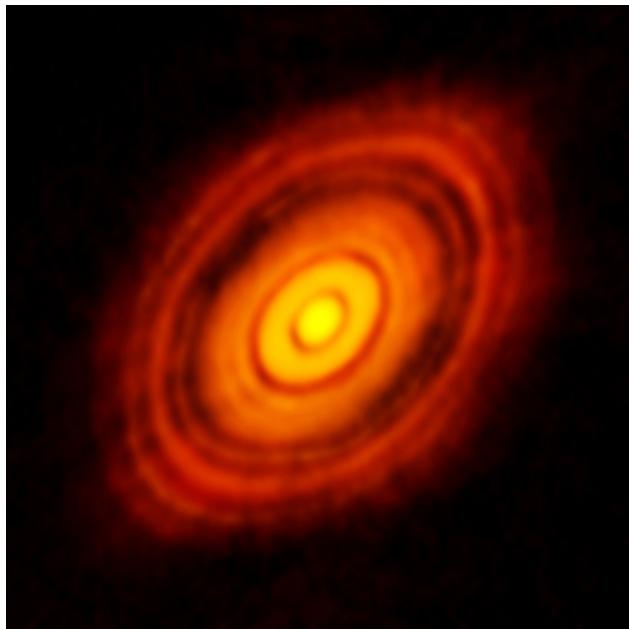
# Planet-disk interaction: gaps, spirals, and vortices.



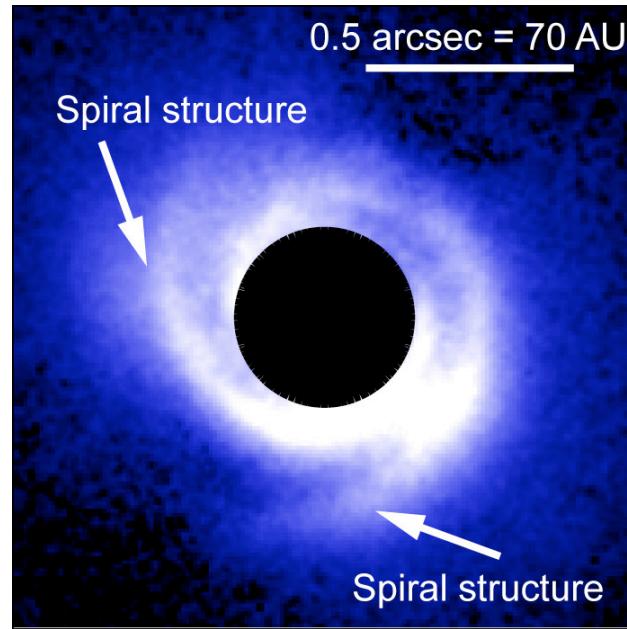
(Lyra et al. 2009b)

# Observational evidence: gaps, spirals, and vortices

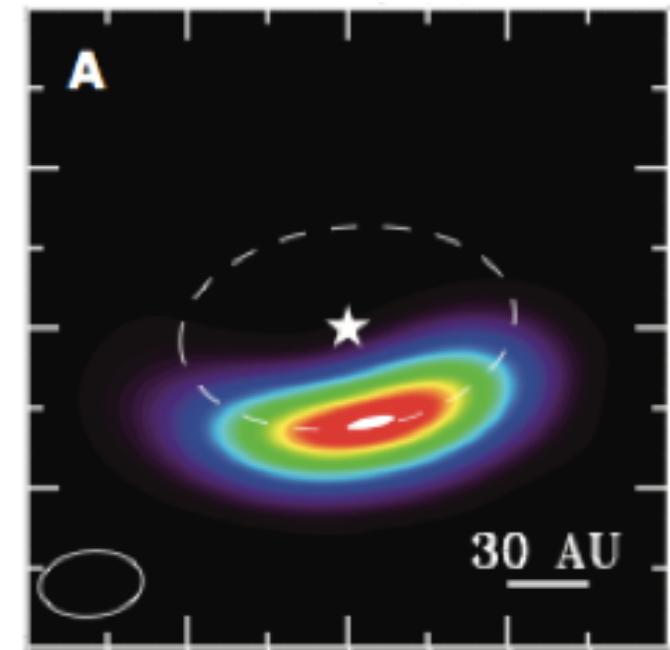
HL Tau



SAO 206462



Oph IRS 48



The ALMA Partnership et al. (2015)

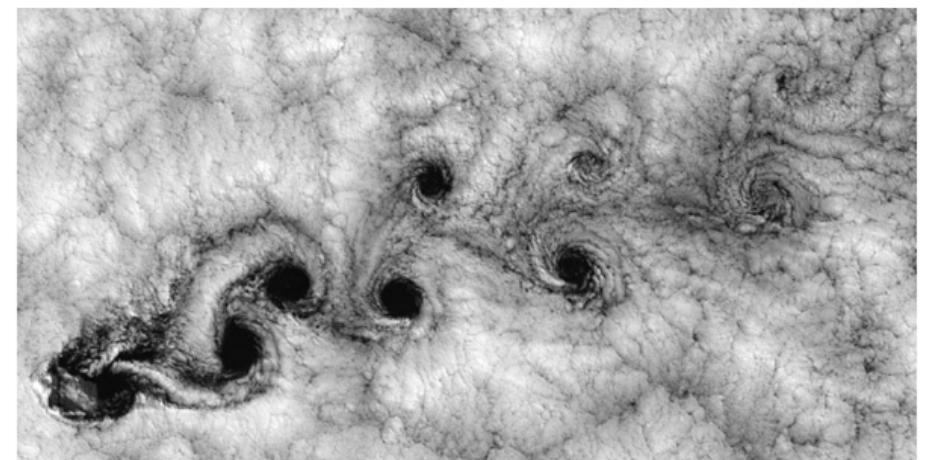
Muto et al. (2012)

van der Marel et al. (2013)

# Vortices – an ubiquitous fluid mechanics phenomenon

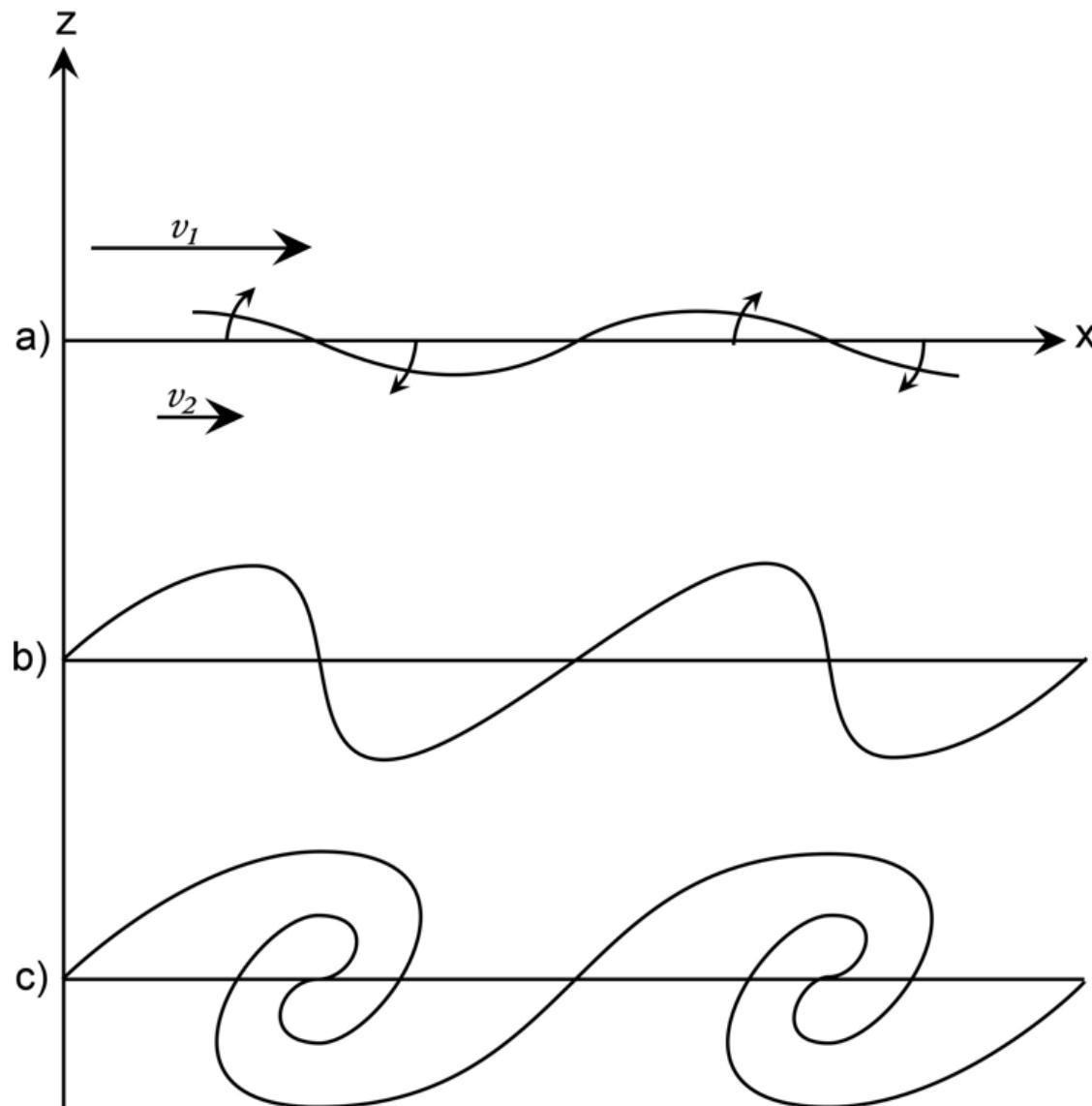


## Von Kármán *vortex street*



# Rossby wave instability

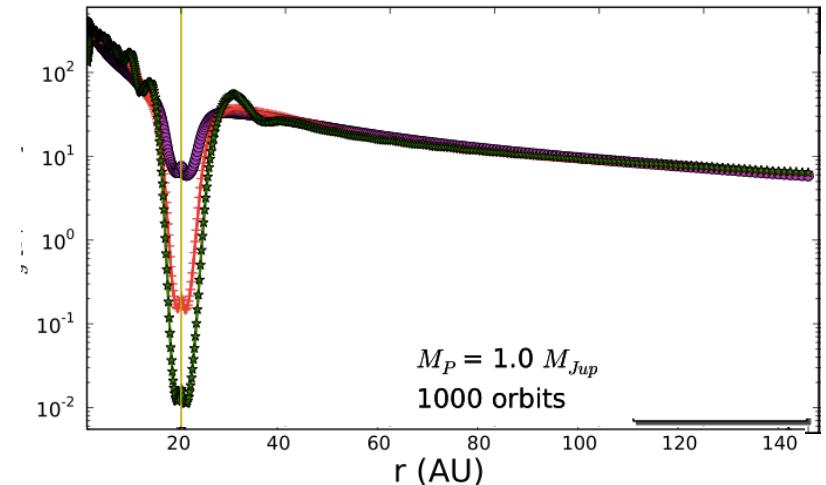
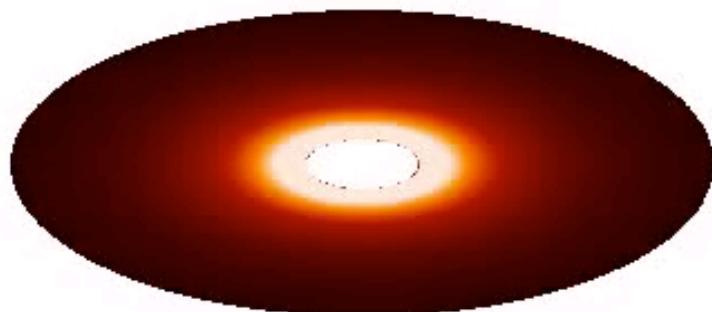
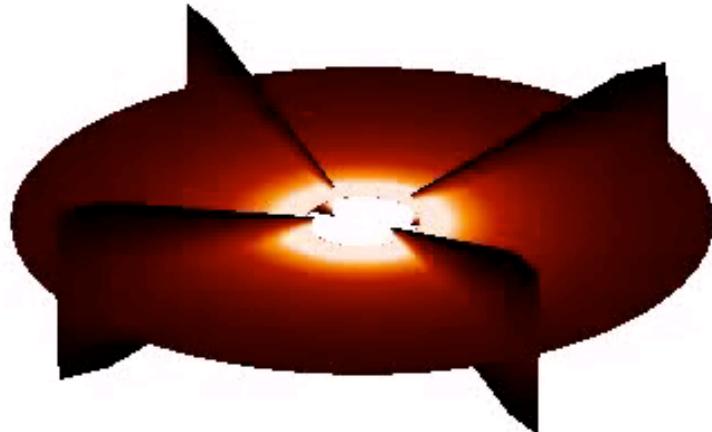
(or Kelvin-Helmholtz instability in differentially rotating gas)



© Brooks Martner

# Planet-disk interaction: gaps, spirals, and vortices.

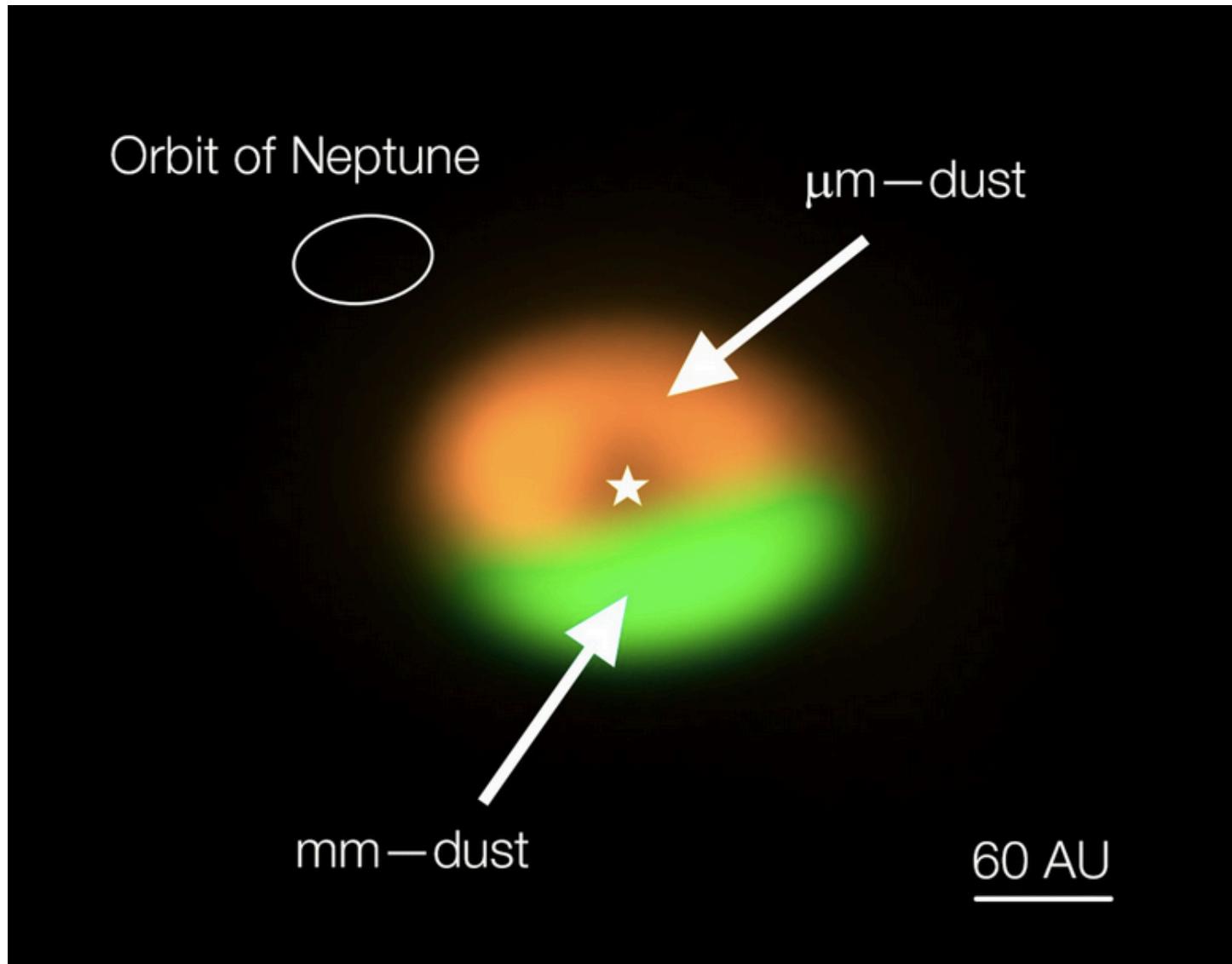
$t = 0.1$



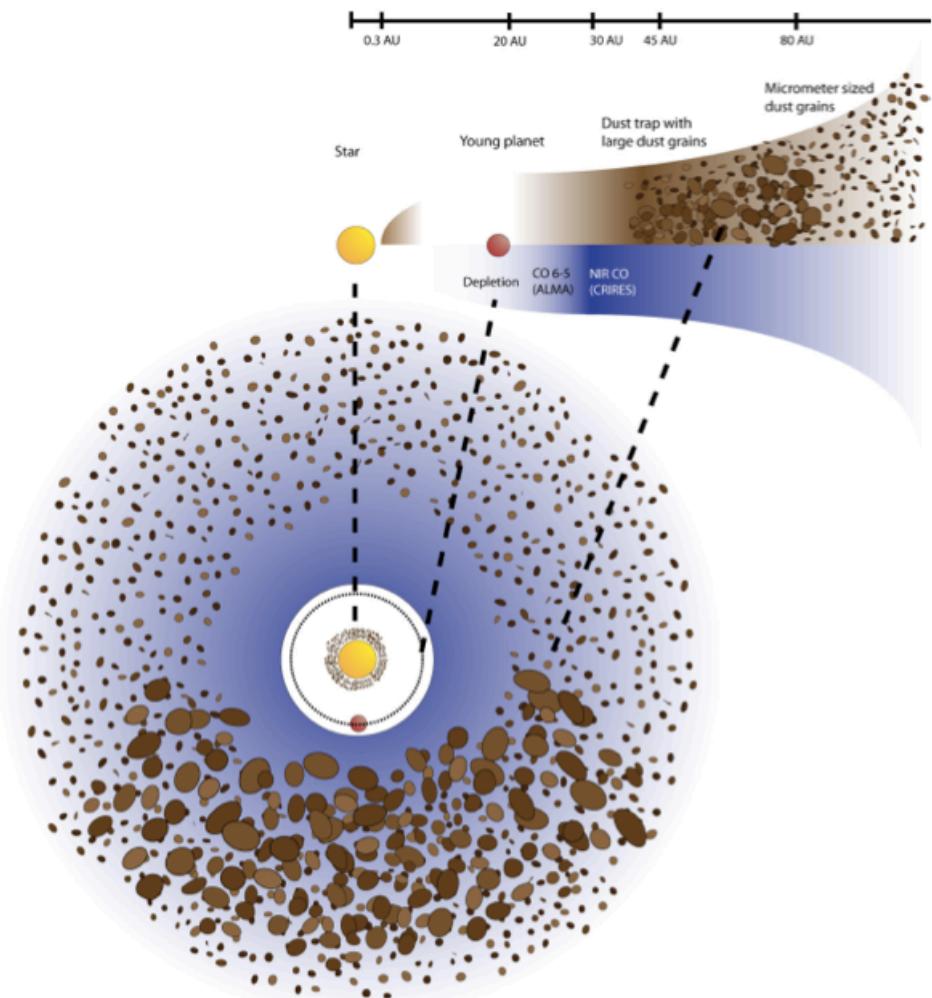
Planet tides carve gap

Gap walls are unstable to  
Rossby wave Instability

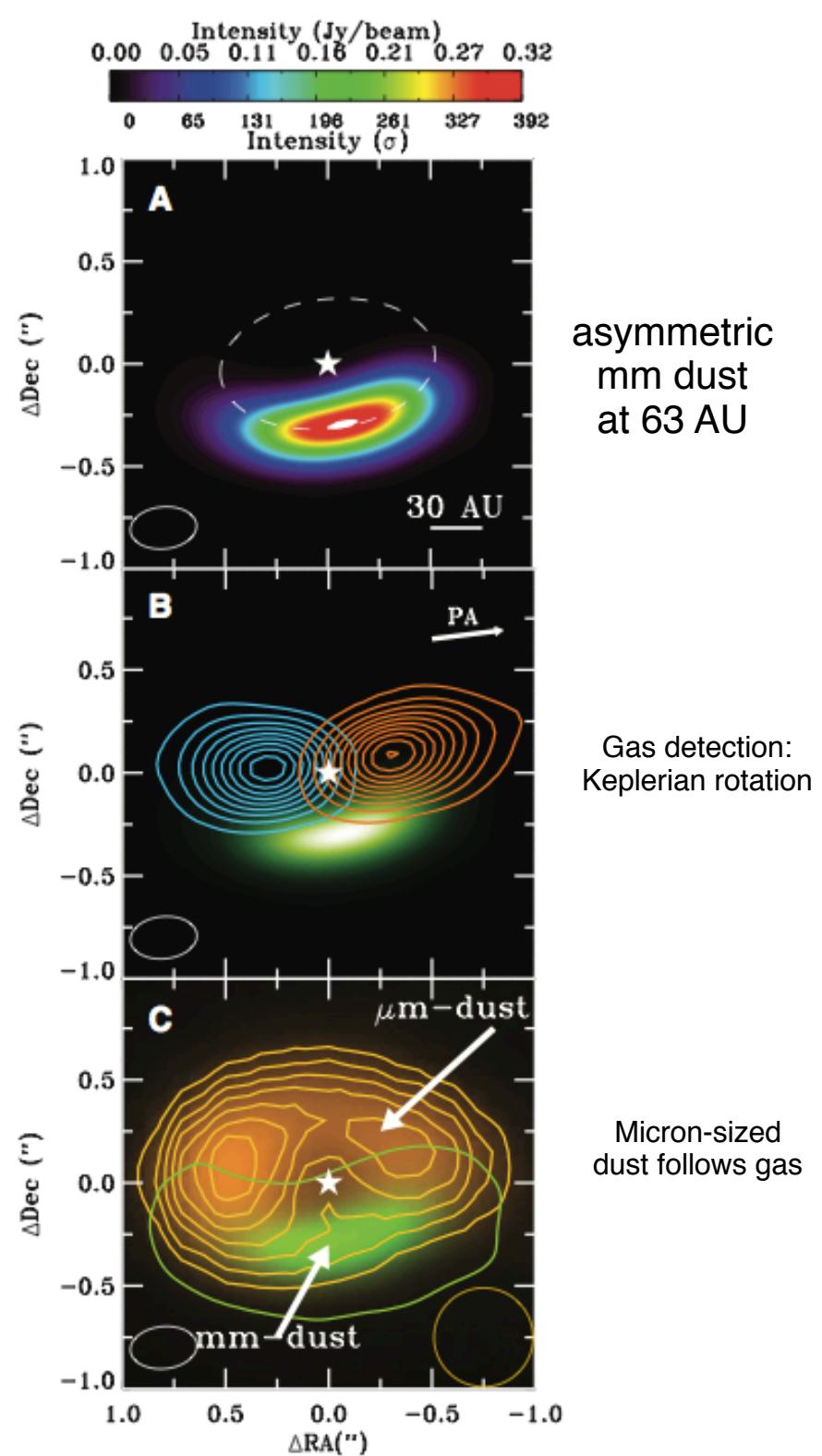
# Oph IRS 48



# The Oph IRS 48 “dust trap”

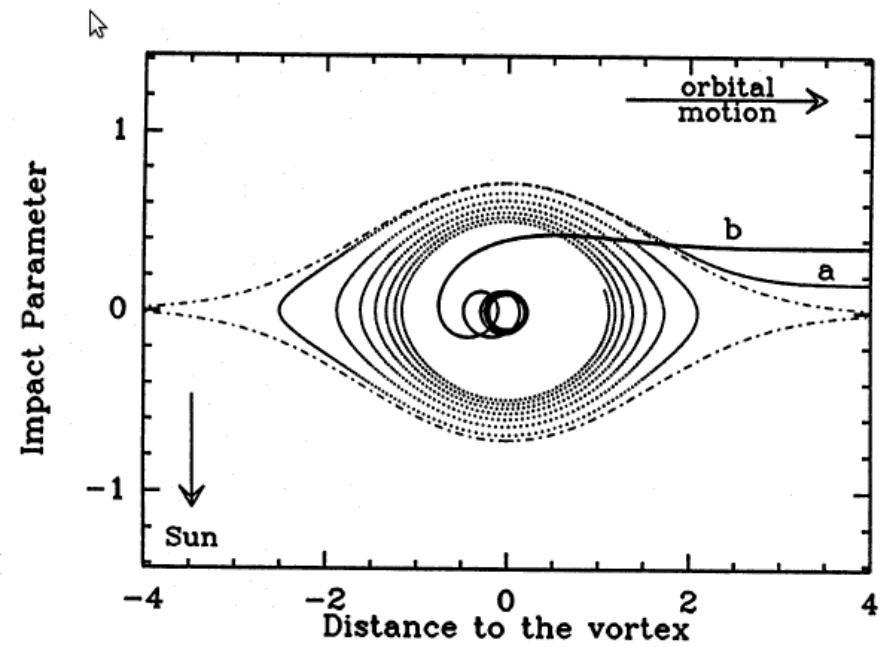
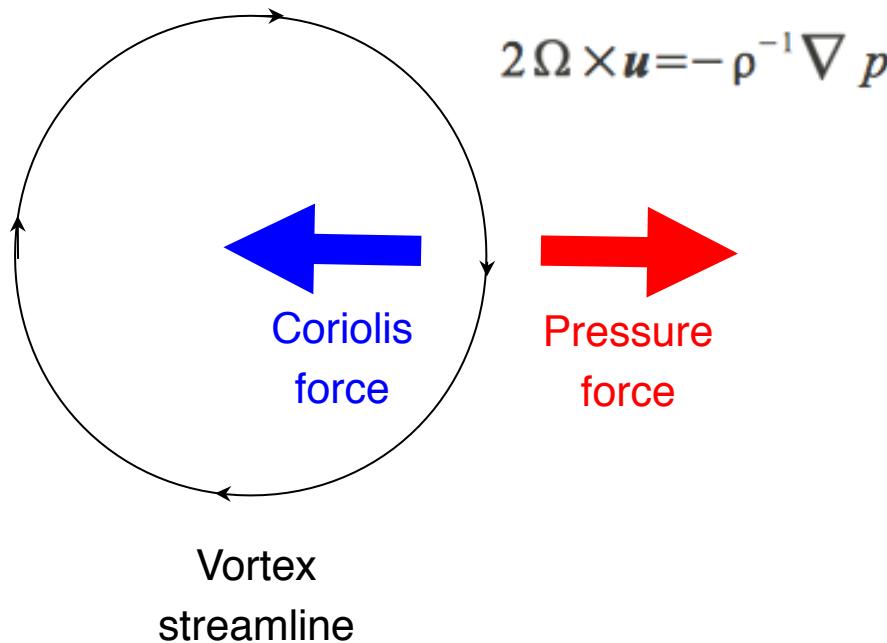


van der Marel et al. (2013)



# The Tea-Leaf effect

Geostrophic balance:



Grains do not feel the pressure gradient.  
They sink towards the center, where they accumulate.

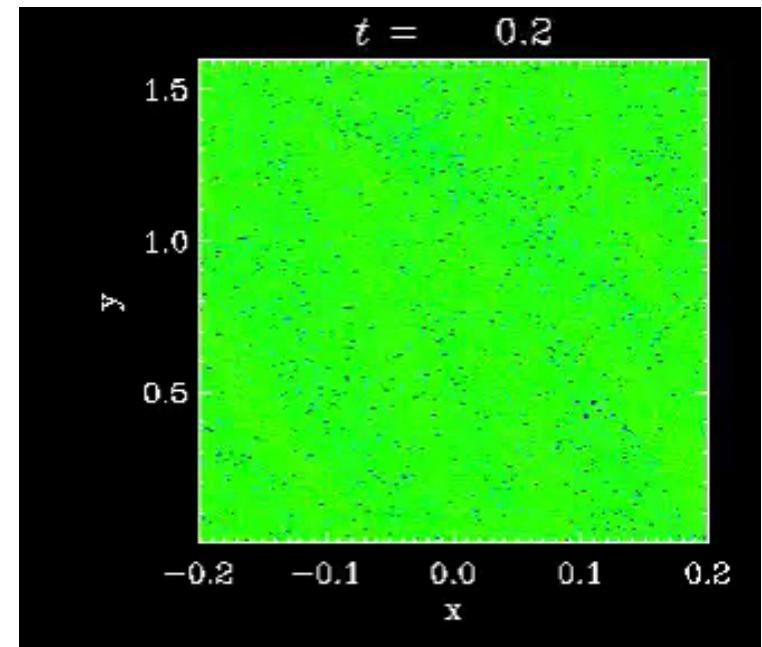
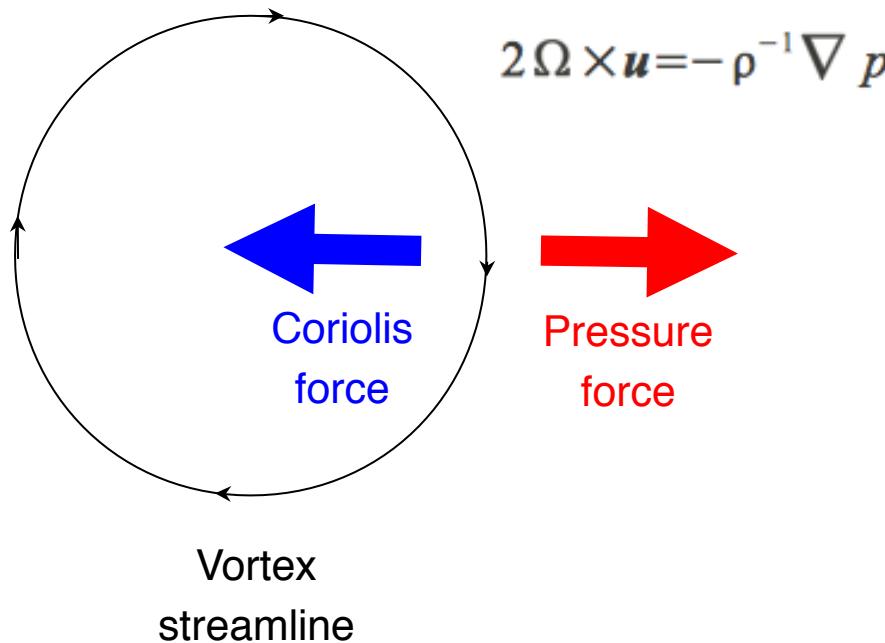
Aid to planet formation

(Barge & Sommeria 1995, Tanga et al. 1996, Barranco & Marcus 2005)

Speed up planet formation enormously  
(Lyra et al. 2008b, 2009ab, Raettig et al. 2012)

# The Tea-Leaf effect

Geostrophic balance:



Raettig, Lyra, & Klahr (2013)

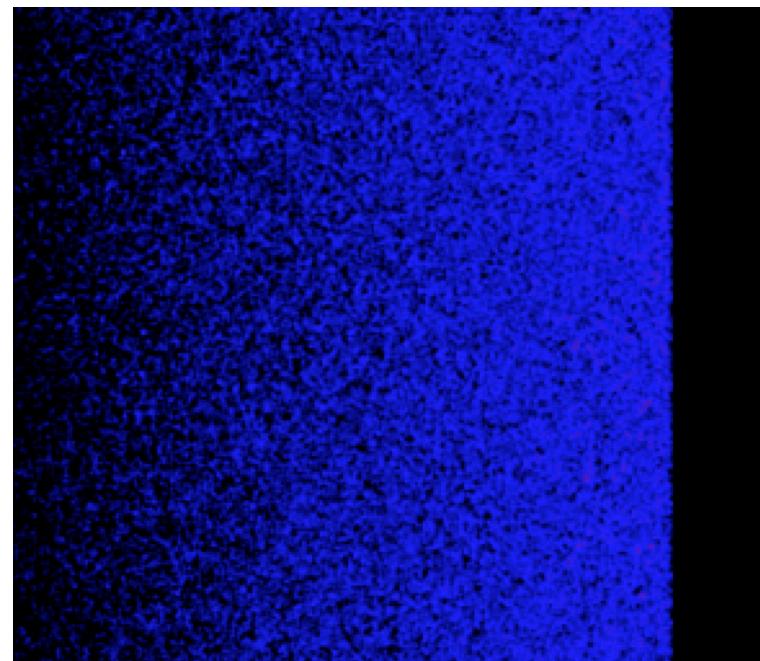
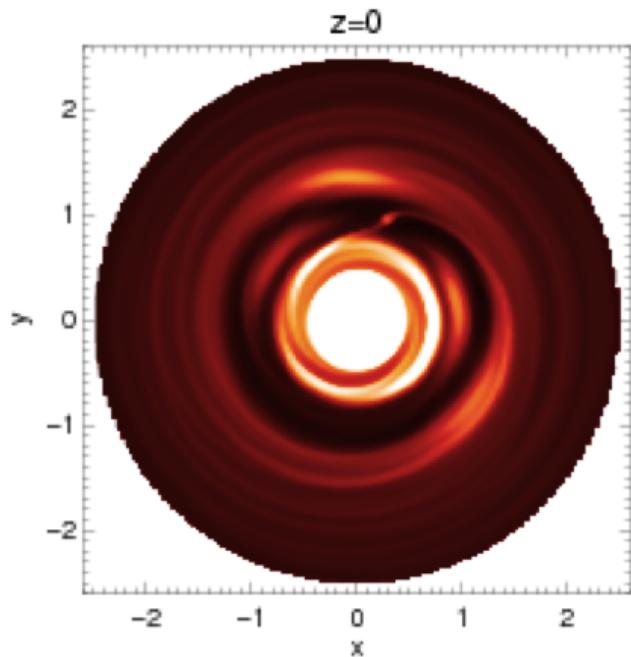
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# Planet Formation in gap edge vortices

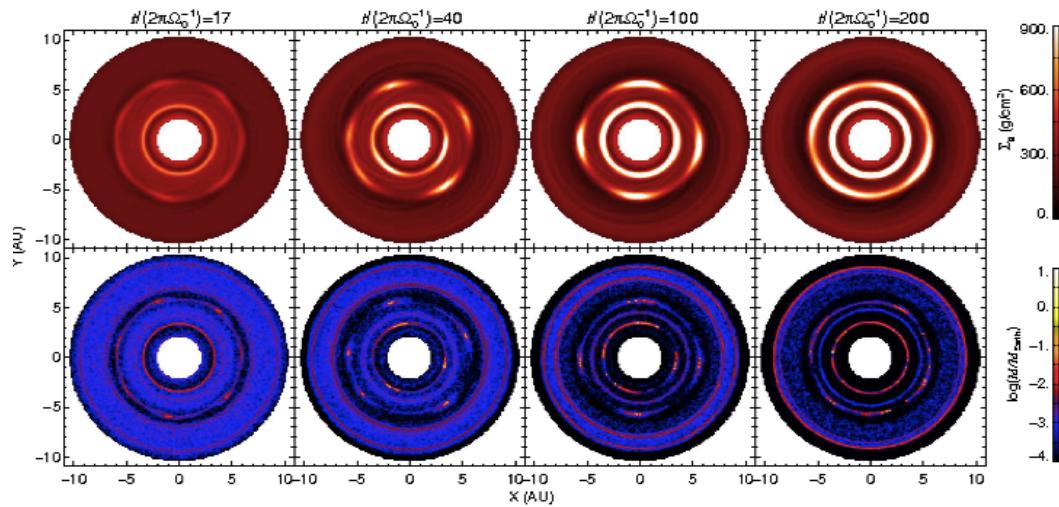
Lyra et al. (2009b),  
see also de Val-Borro et al. (2007)



Burst of formation in gap vortices

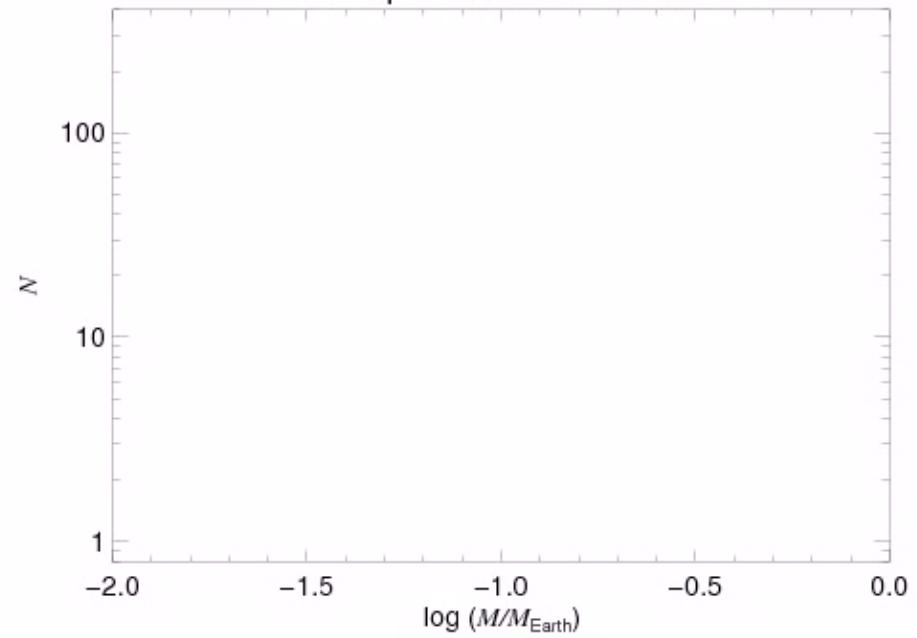
# Vortices and Planet Formation

Gas



Grains

Mass Spectrum  $t=0/204$  orbits

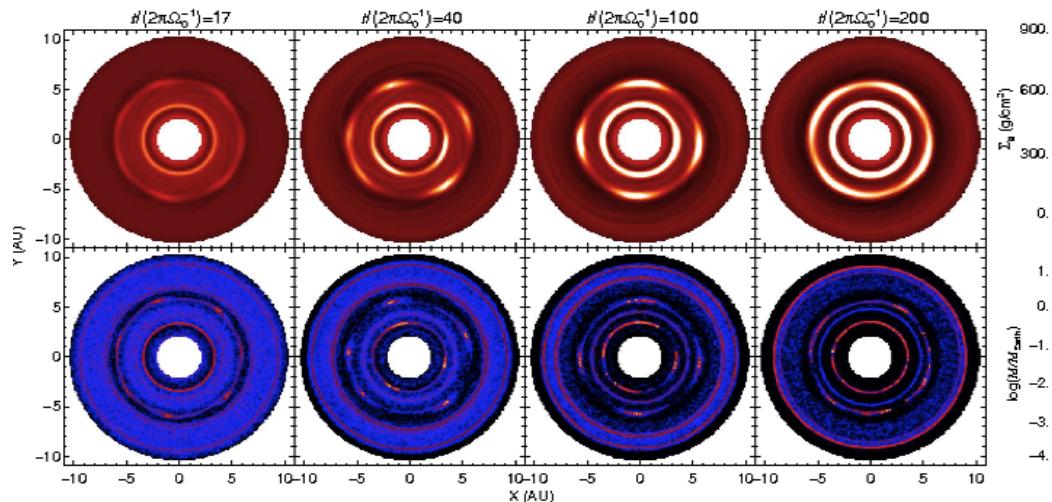


Collapse into Mars mass objects

(Lyra et al. 2008b, 2009a,  
Lambrechts & Johansen 2012)

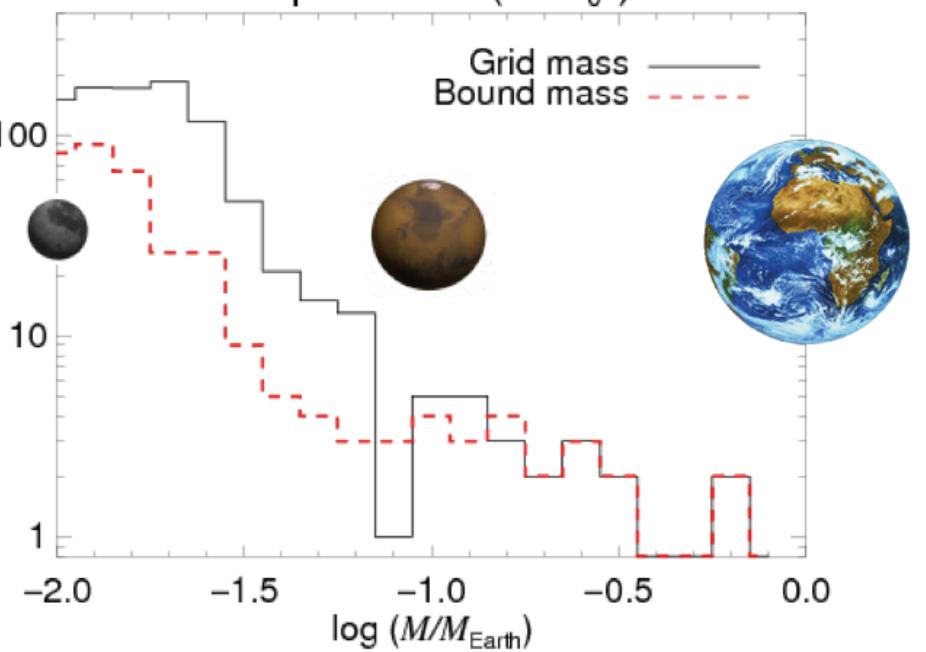
# Vortices and Planet Formation

Gas



Grains

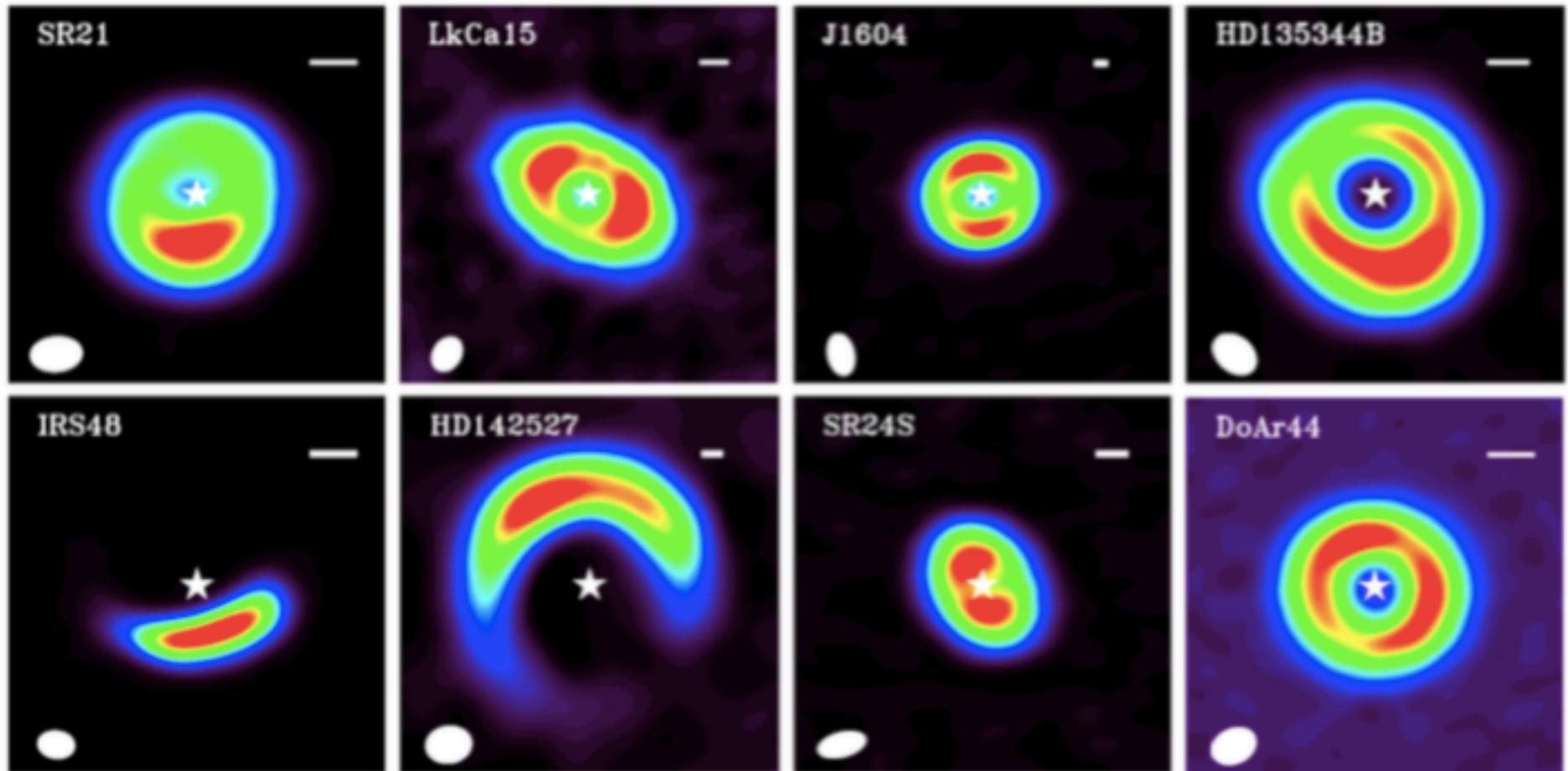
Mass Spectrum  $t/(2\pi\Omega_0^{-1})=200$



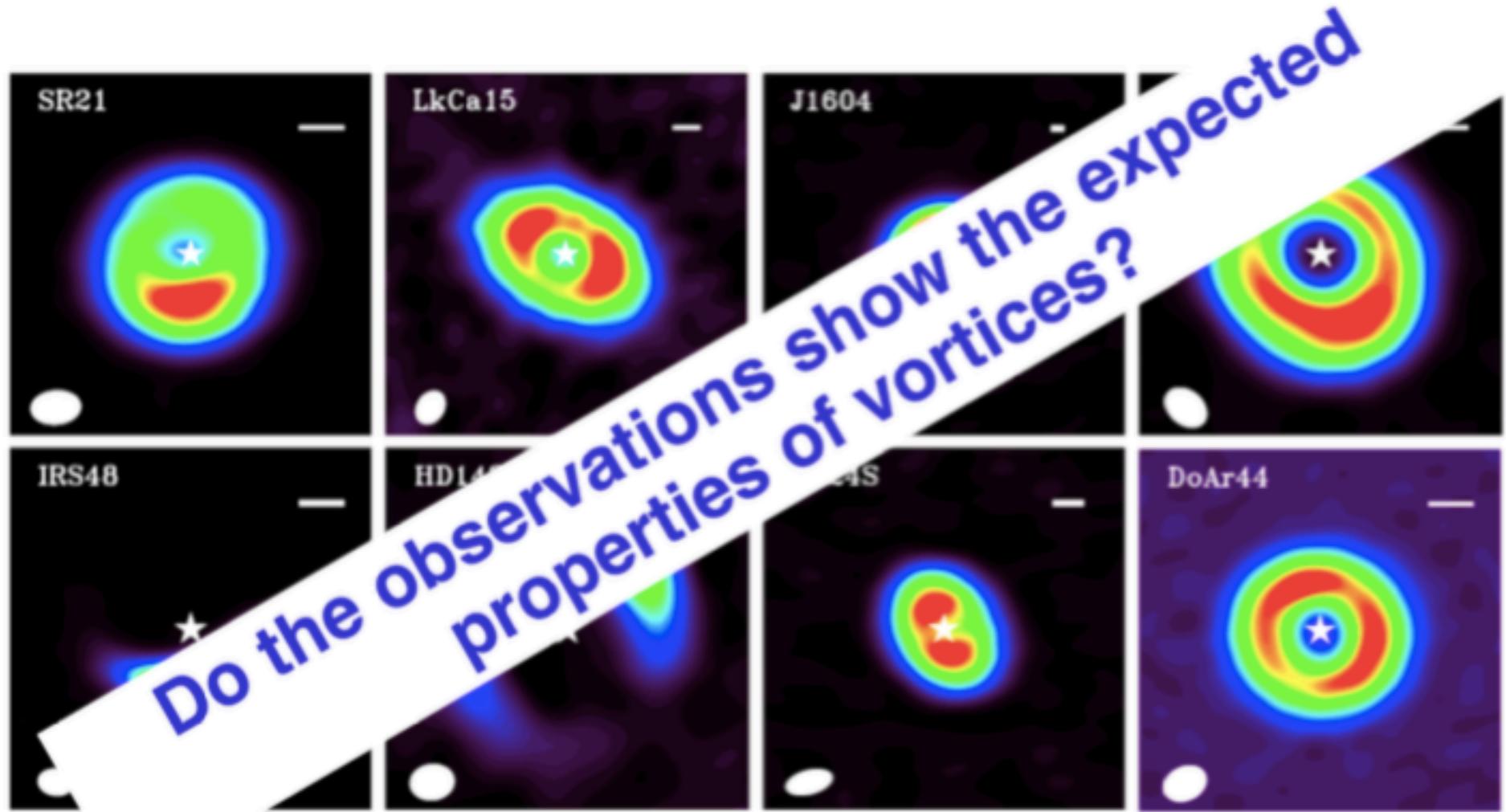
Collapse into Mars mass objects

(Lyra et al. 2008b, 2009a,  
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# “Asymmetries” everywhere!

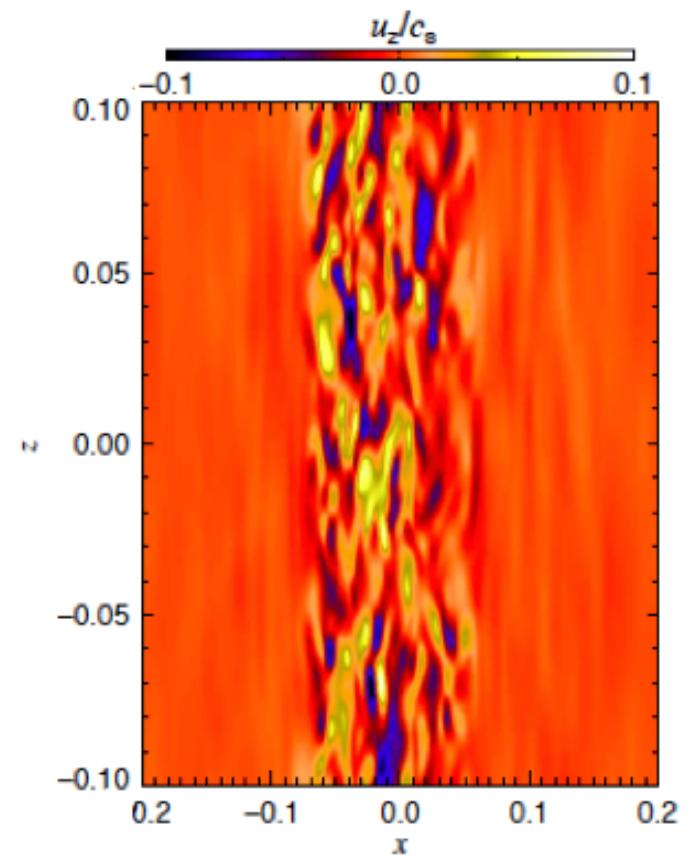
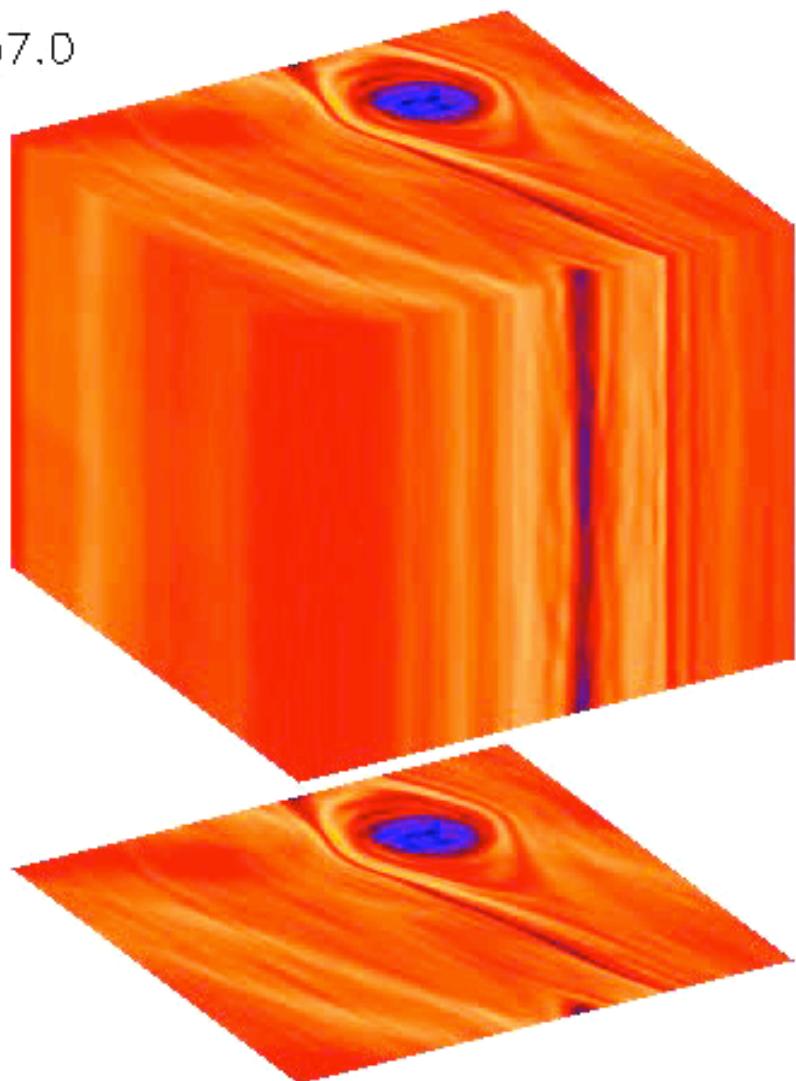


**“Asymmetries” everywhere!**



## Turbulence in vortex cores

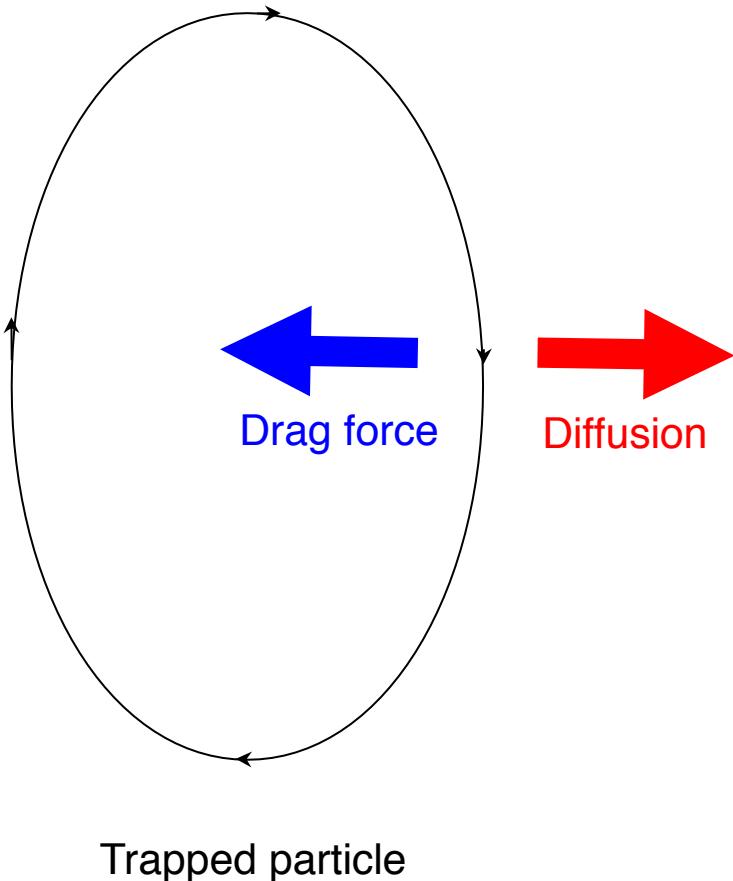
$t=1257.0$



Lyra & Klahr (2011)

$V_{\max}$ : ~10% of sound speed  
 $V_{rms}$ : ~3% of sound speed

# Drag-Diffusion Equilibrium



Dust continuity equation

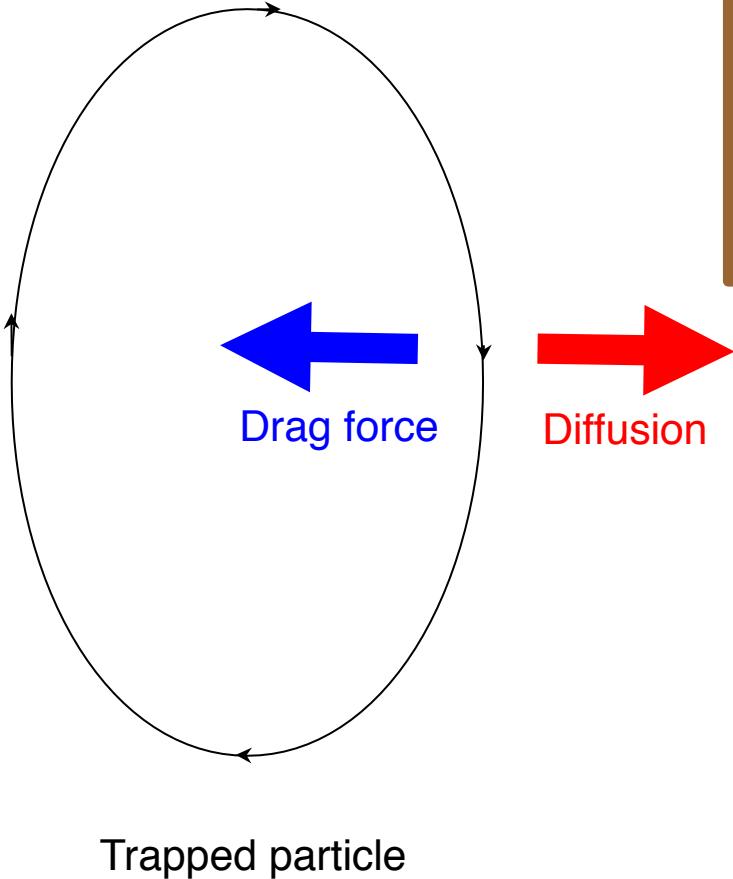
$$\frac{\partial \rho_d}{\partial t} = -(\mathbf{v} \cdot \nabla) \rho_d - \rho_d \nabla \cdot \mathbf{v} + D \nabla^2 \rho_d,$$

advection

compression

diffusion

# Lyra-Lin solution



Steady-state solution

$$\rho_d(a,z) = \epsilon \rho_0 (S + 1)^{3/2} \exp \left\{ - \frac{[a^2 f^2(\chi) + z^2]}{2H^2} (S + 1) \right\}$$

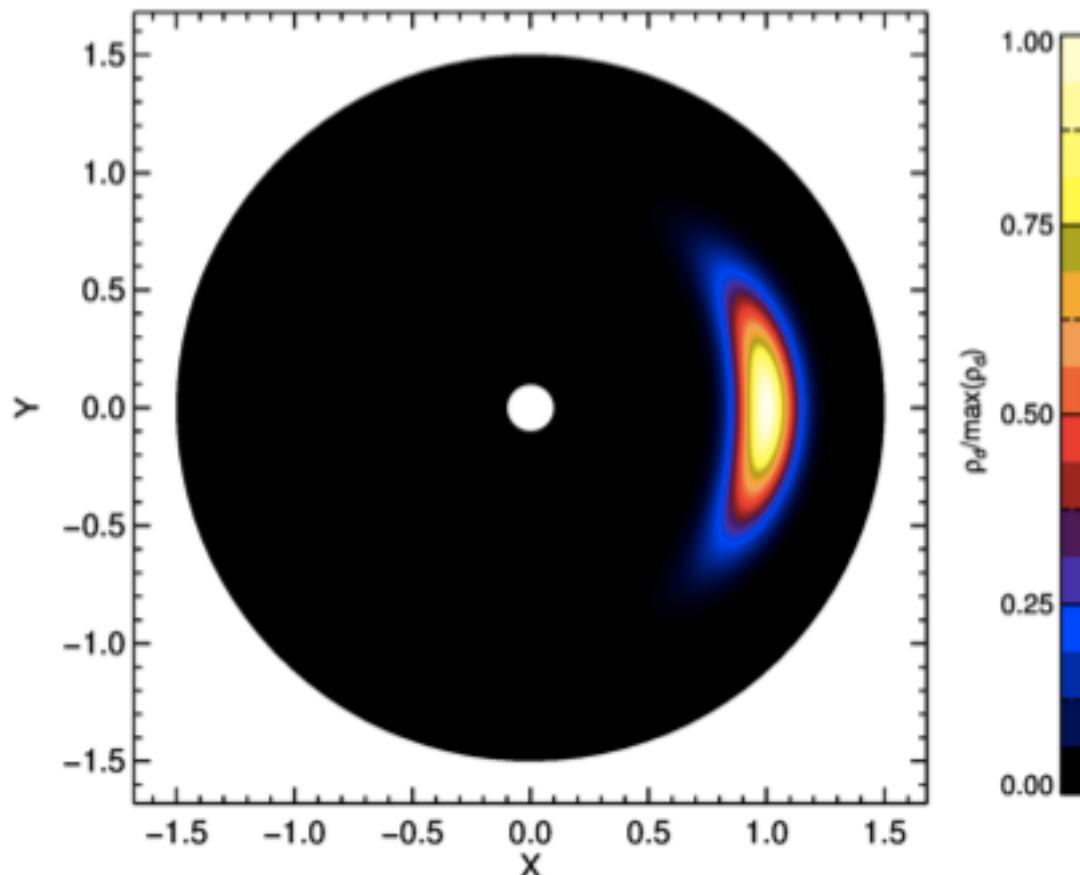
Lyra & Lin (2013)

$$S = \frac{St}{\delta}$$

$$\delta = v_{\text{rms}}^2 / c_s^2,$$

- |           |                                   |
|-----------|-----------------------------------|
| $a$       | = vortex semi-minor axis          |
| $H$       | = disk scale height (temperature) |
| $\chi$    | = vortex aspect ratio             |
| $\delta$  | = diffusion parameter             |
| $St$      | = Stokes number (particle size)   |
| $f(\chi)$ | = model-dependent scale function  |

# Analytical solution for dust in drag-diffusion equilibrium



Solution for

$$H/r=0.1 \quad \chi=4 \quad S=1$$

## Solution

$$\rho_d(a) = \rho_{d\max} \exp\left(-\frac{a^2}{2H_V^2}\right),$$

$$H_V = \frac{H}{f(\chi)} \sqrt{\frac{1}{S+1}}$$

$$S = \frac{St}{\delta}$$

$$\delta = v_{\text{rms}}^2 / c_s^2,$$

$a$  = distance to vortex center  
 $H$  = disk scale height (temperature)  
 $\chi$  = vortex aspect ratio  
 $\delta$  = diffusion parameter  
St = Stokes number (grain size)  
 $f(\chi)$  = model-dependent scale function

## Derived quantities

$$\rho_d(a,z) = \epsilon \rho_0 (S+1)^{3/2} \exp \left\{ - \frac{[a^2 f^2(\chi) + z^2]}{2H^2} (S+1) \right\}$$

Lyra & Lin (2013)

Gas distribution

$$\rho_g(a) = \rho_{g\max} \exp \left( - \frac{a^2}{2H_g^2} \right),$$

Maximum dust density

$$\rho_{d\max} = \epsilon \rho_0 (S+1)^{3/2}$$

Gas contrast

$$\frac{\rho_{g\max}}{\rho_{g\min}} = \exp \left[ \frac{f^2(\chi)}{2\chi^2 \omega_V^2} \right],$$

Dust contrast

$$\frac{\rho_{d\max}}{\rho_{d\min}} = \frac{\rho_{g\max}}{\rho_{g\min}} \exp(S),$$

Total trapped mass

$$\int \rho_d(a,z) dV = (2\pi)^{3/2} \epsilon \rho_0 \chi H H_g^2$$

Vortex size

$$a_s = H(\chi \omega_V)^{-1}$$

$H$  = disk scale height (temperature)

$\chi$  = vortex aspect ratio

$\delta$  = diffusion parameter

St = Stokes number (particle size)

$f(\chi)$  = model-dependent scale function

$\epsilon$  = dust-to-gas ratio

## Applying the model to Oph IRS 48

### Observed parameters

Aspect ratio: 3.1

Dust contrast: 130

Temperature: 60K

**Trapped mass:  $9 M_{Earth}$**

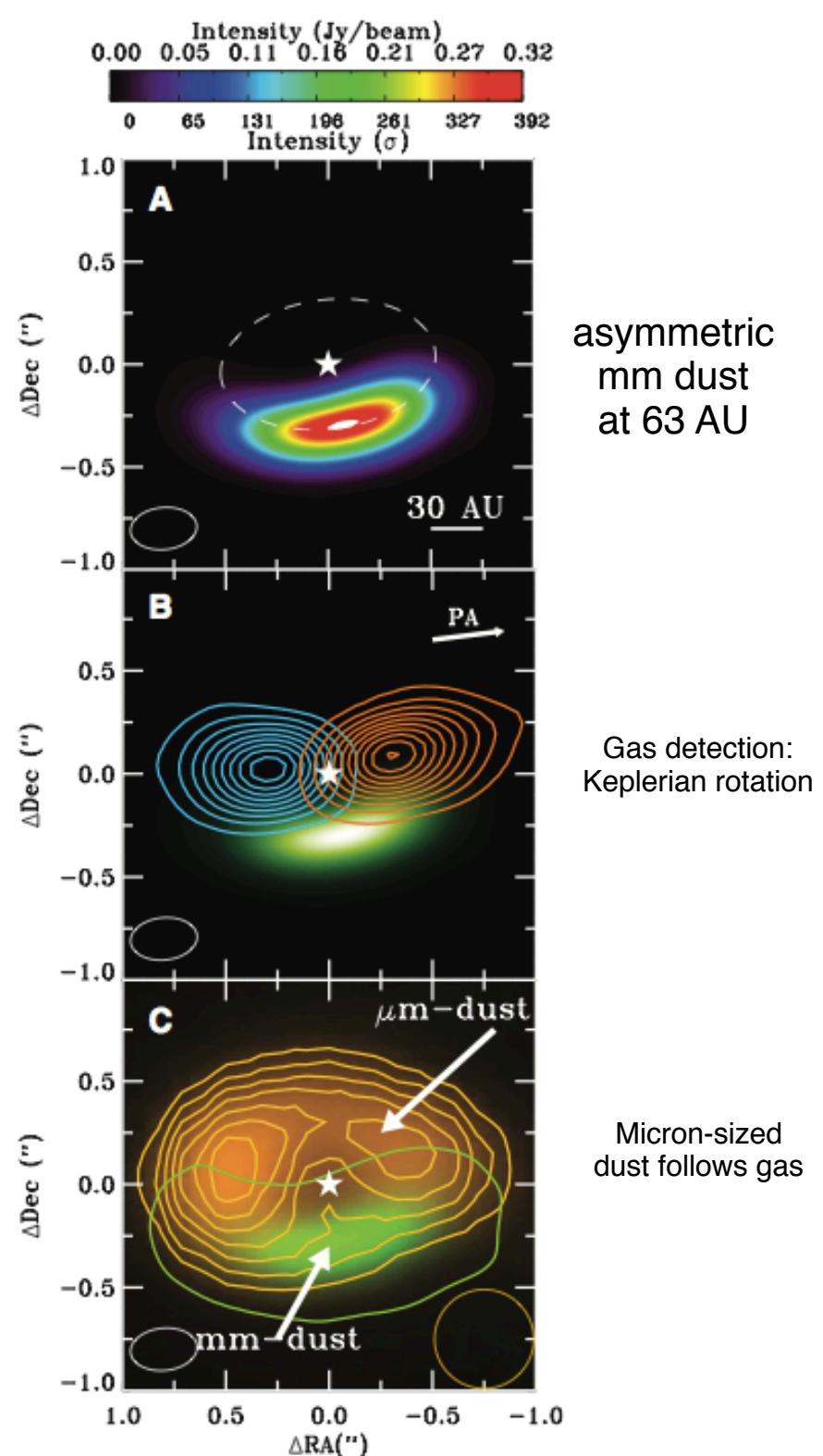
### Derived parameters

$S=4.8$

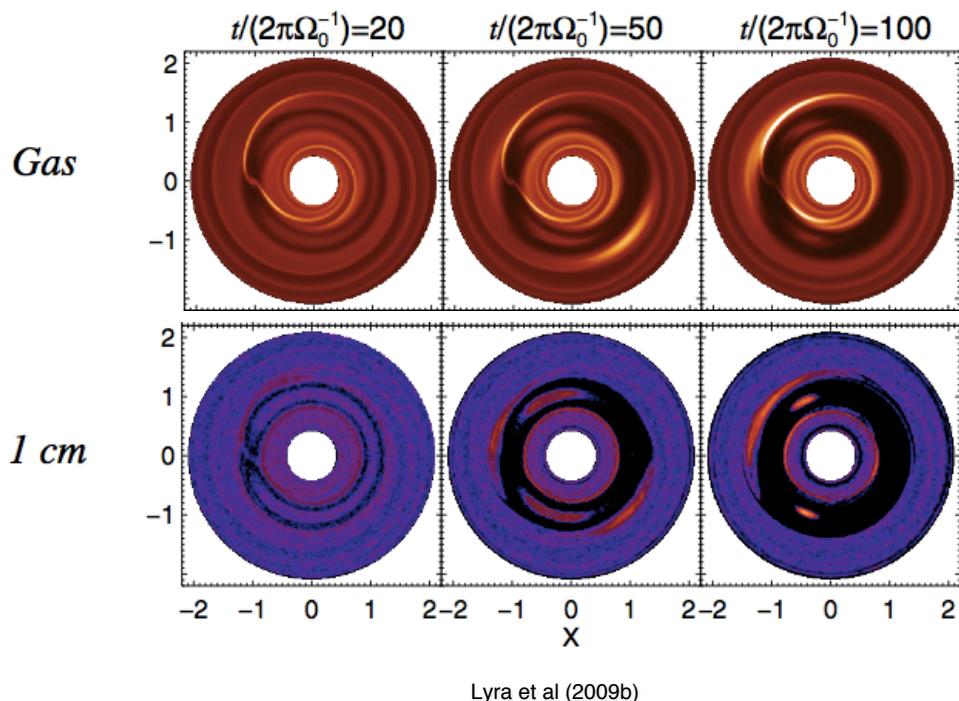
$St=0.008$        $\delta = 0.005$

**$V_{rms} = 4\% c_s$**

**Trapped mass:  $11 M_{Earth}$**



# Planetary gap vortices: Chicken or the Egg problem

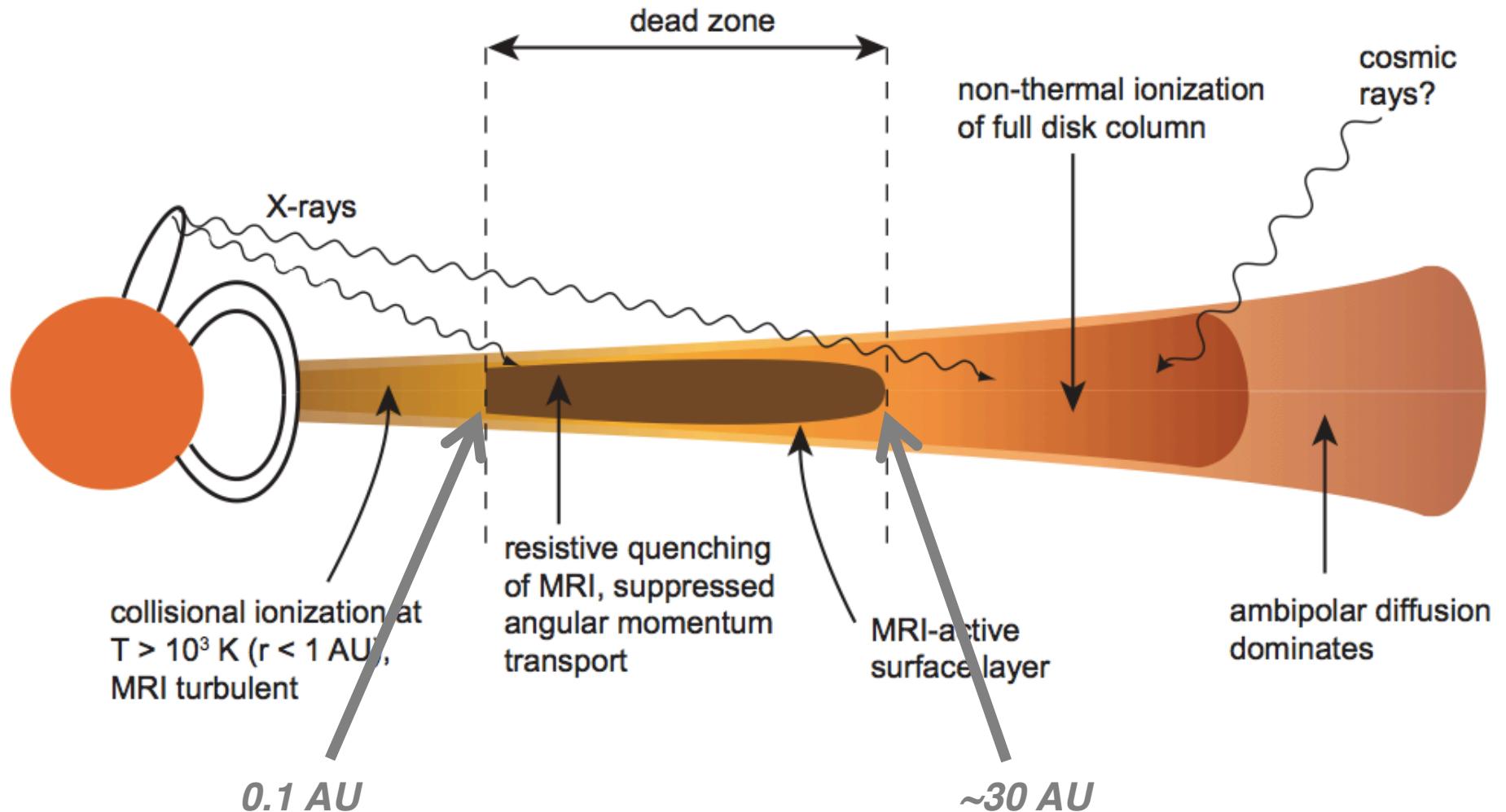


Does the observational detection of a particle trap in IRS 48 imply that traps are the answer to surmounting the radial drift barrier and allowing planet formation? Not immediately. Particle traps solve theoretical problems in planet formation that exist at millimeter to meter scales, and they are no solution at all if the only way to form them requires that gas giant planets already exist. The trap observed in the IRS 48 disk might instead catalyze the formation of additional

Armitage (2013)

# **Turbulence and Accretion in 3D Global MHD Simulations of Stratified Protoplanetary Disk**

# Dead zones

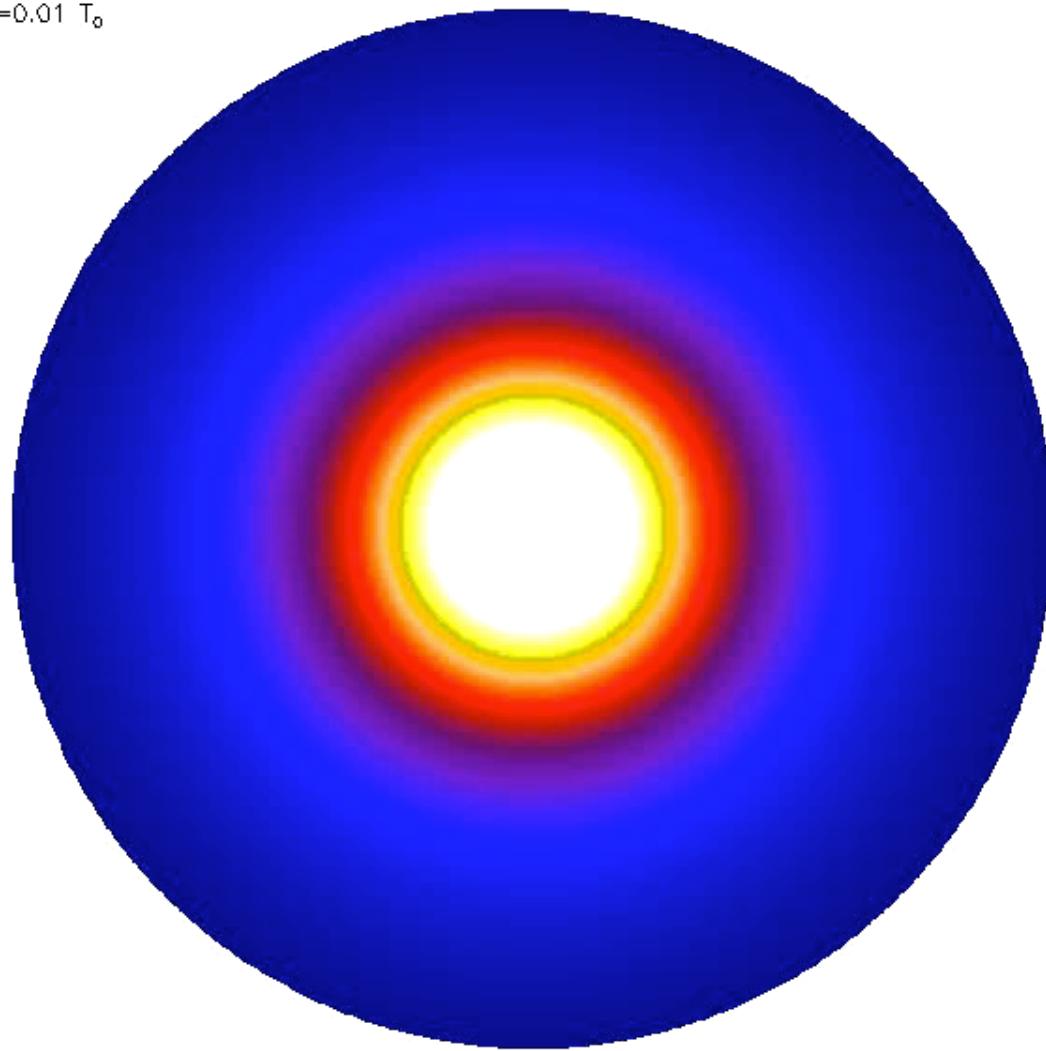


There should be a **magnetized, active zone**  
and a **non-magnetic, dead zone**



## Inner (0.1 AU) active/dead zone boundary

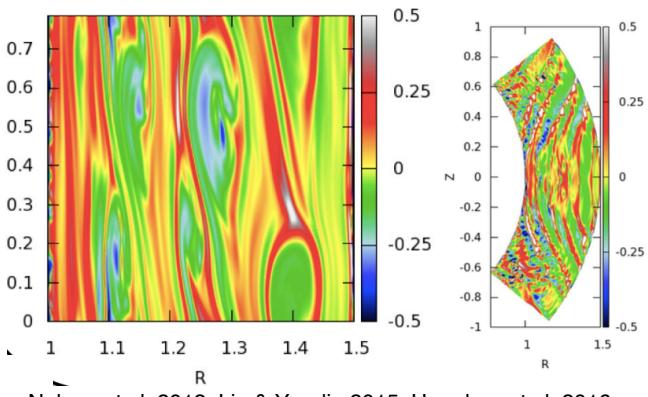
$t=0.01 T_0$



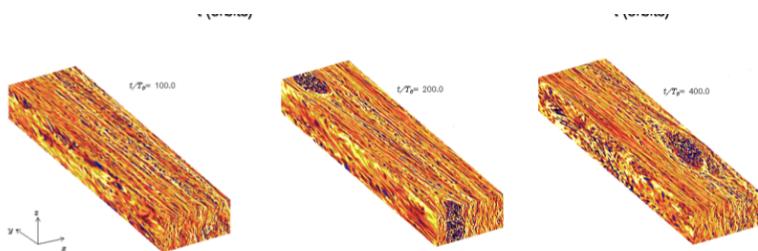
Magnetized inner disk + resistive outer disk  
Lyra & Mac Low (2012)

# Thermal Instabilities

## Vertical Shear Instability

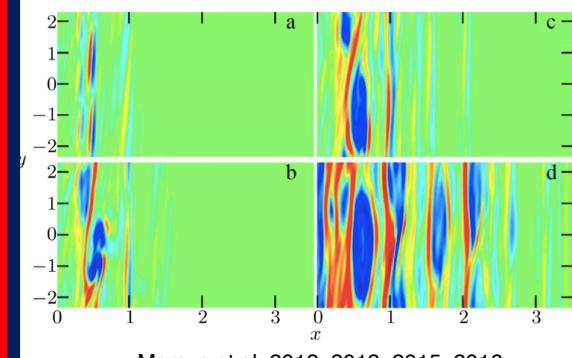


## Convective Overstability



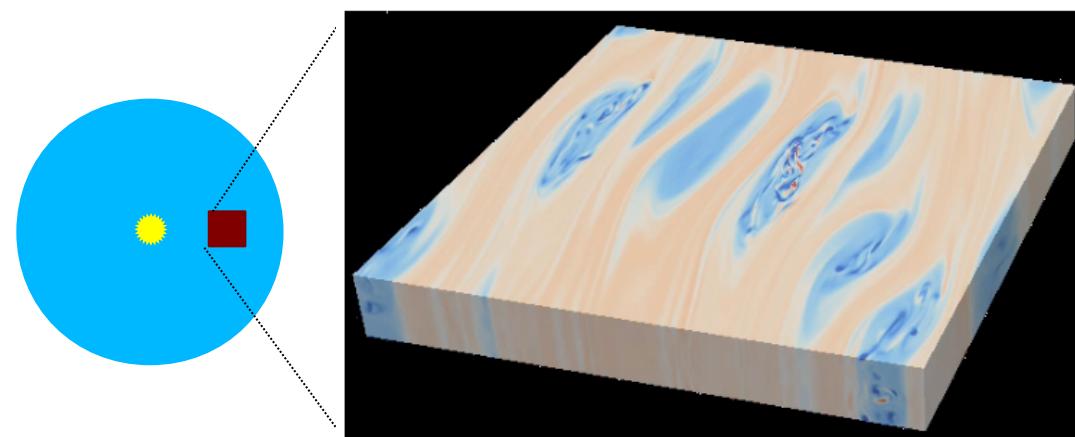
Klahr 2003, Lesur & Papaloizou 2010, Lyra & Klahr 2011, Lyra 2014

## Zombie Vortex Instability



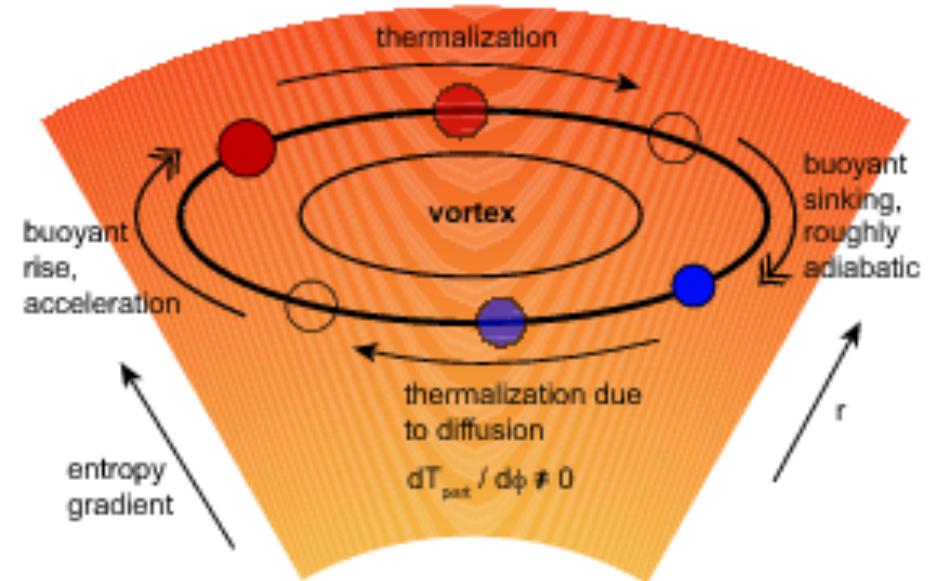
Marcus et al. 2012, 2013, 2015, 2016  
Umurhan et al. 2016, Lesur & Latter 2016

# Convective Overstability



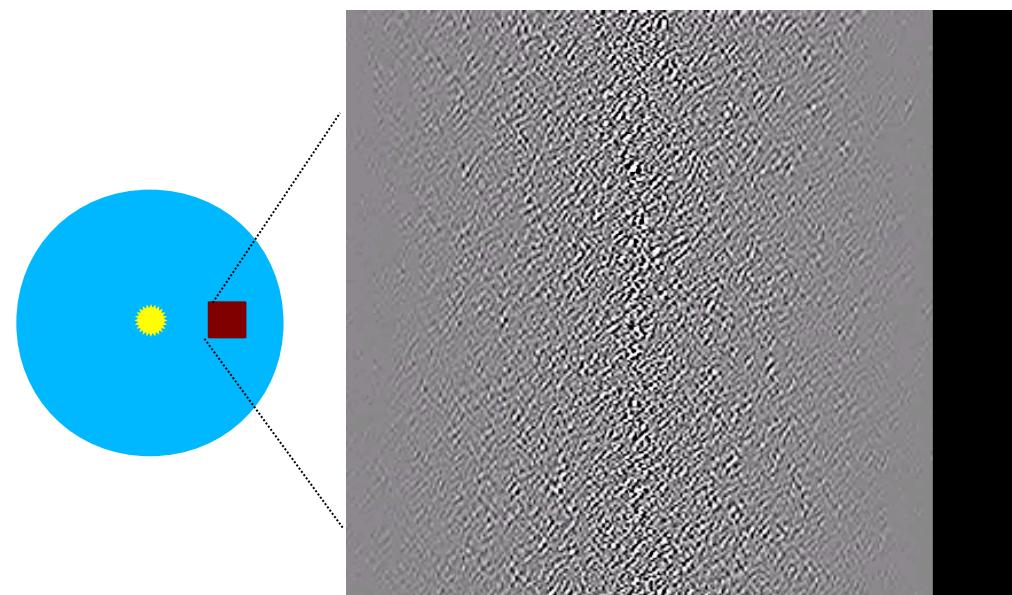
Lesur & Papaloizou (2010)

Sketch of the  
Convective Overstability



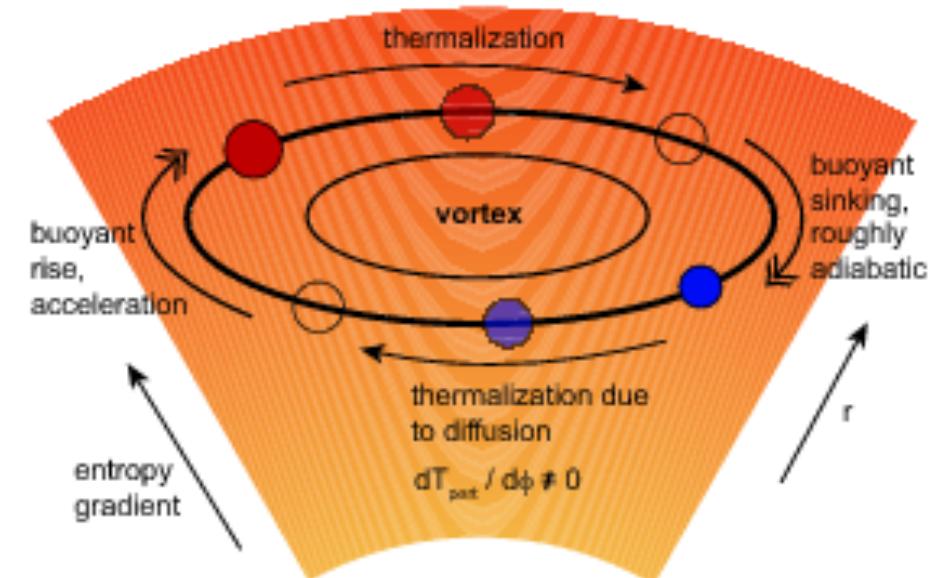
Armitage (2010)

# Convective Overstability



Lyra & Klahr (2011)

Sketch of the  
Convective Overstability



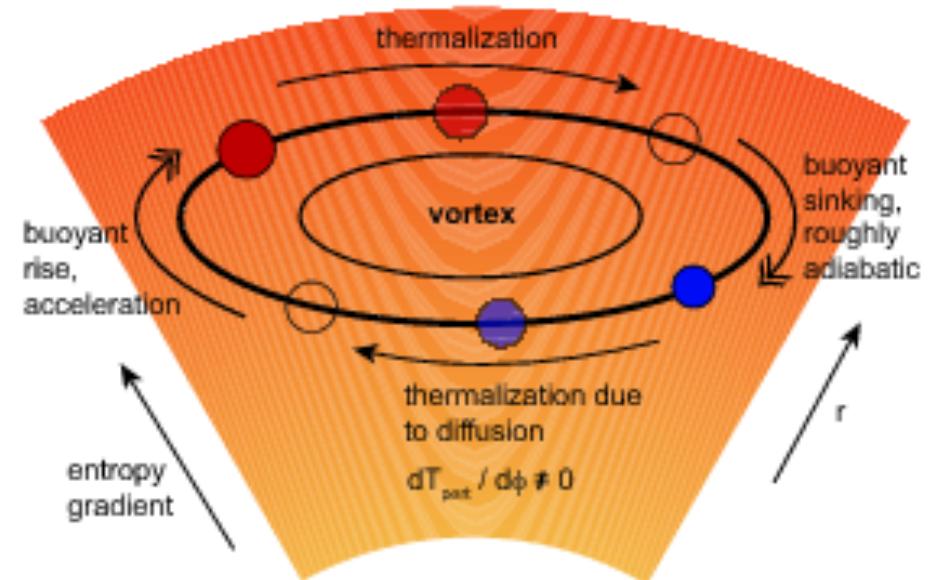
Armitage (2010)

# Convective Overstability

Sketch of the  
Convective Overstability

1. *Radial entropy gradient*
2. *Finite cooling time*

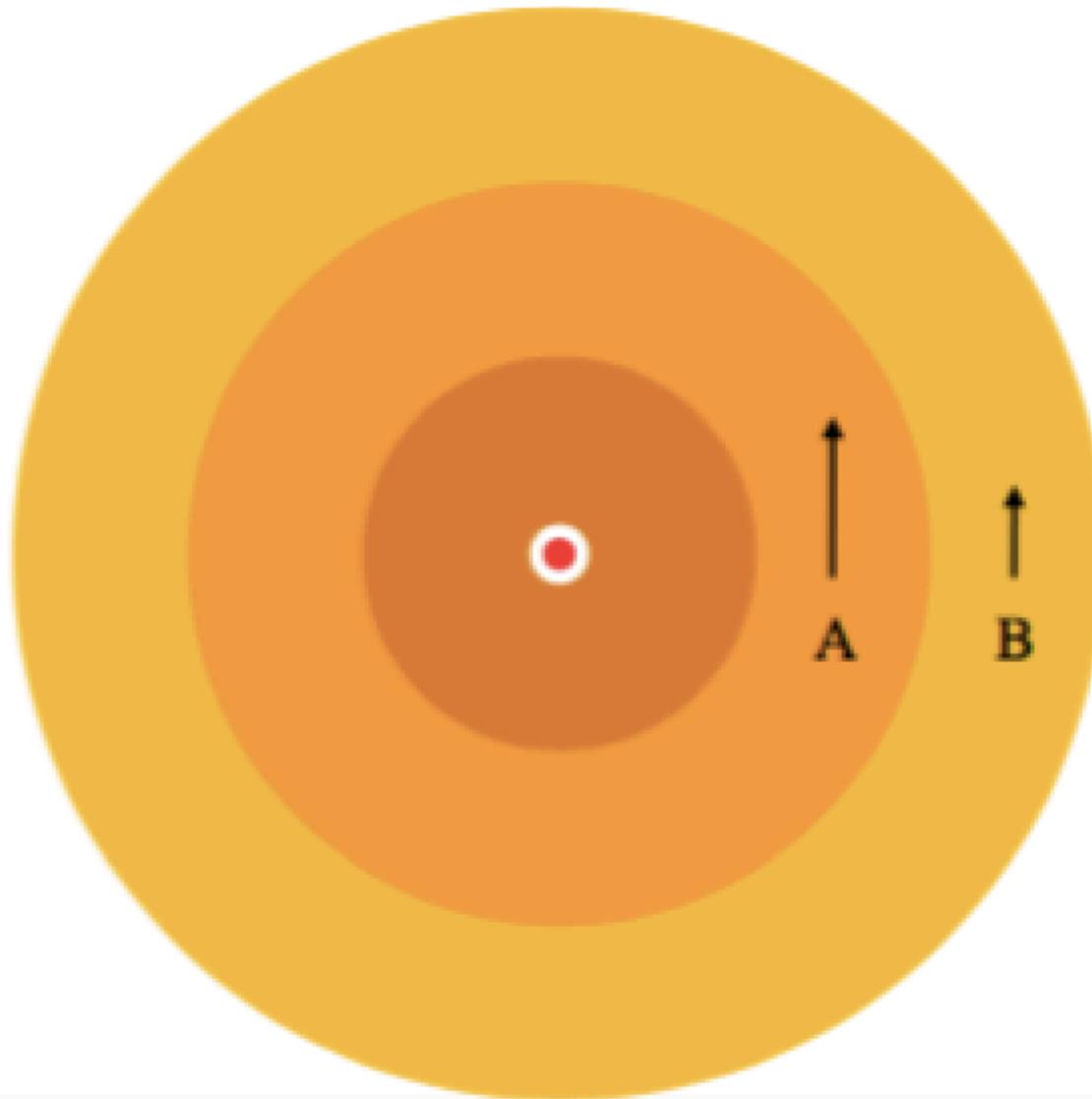
Lesur & Papaloizou (2010)



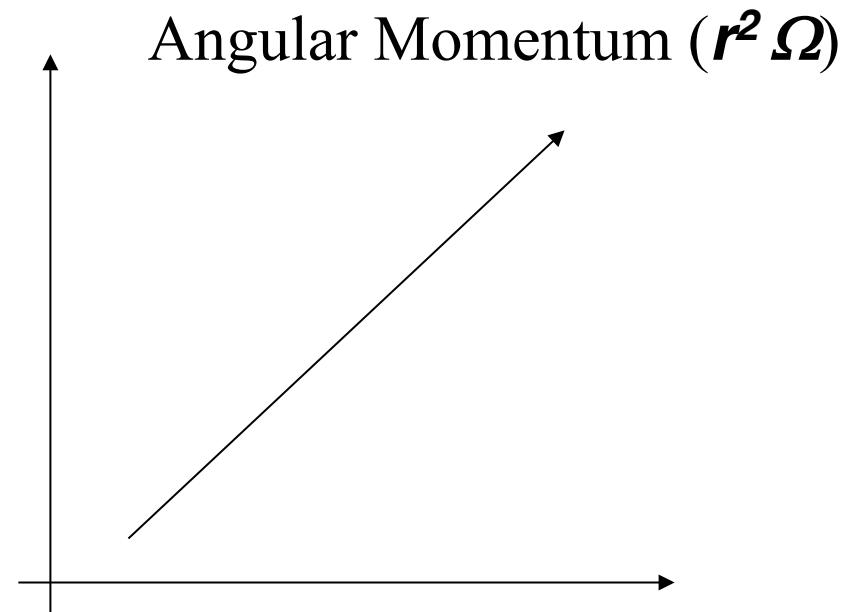
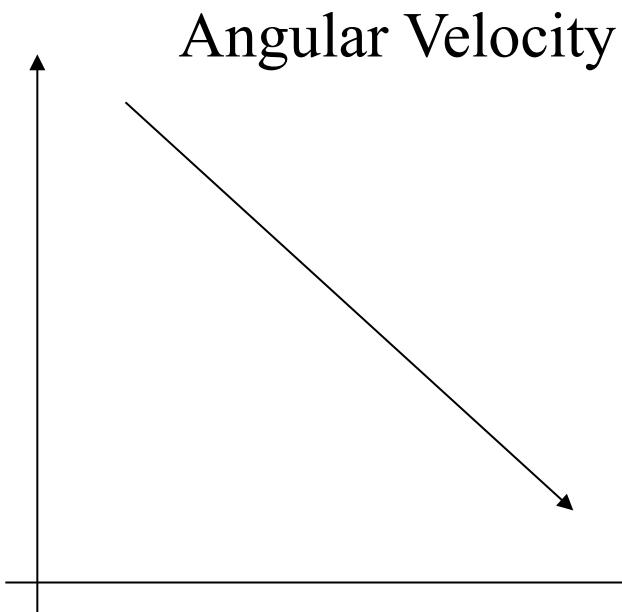
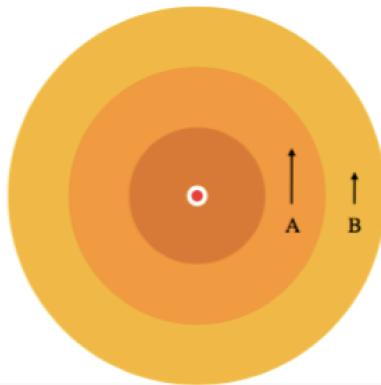
Armitage (2010)

# Vertical Shear Instability

Rayleigh criterion

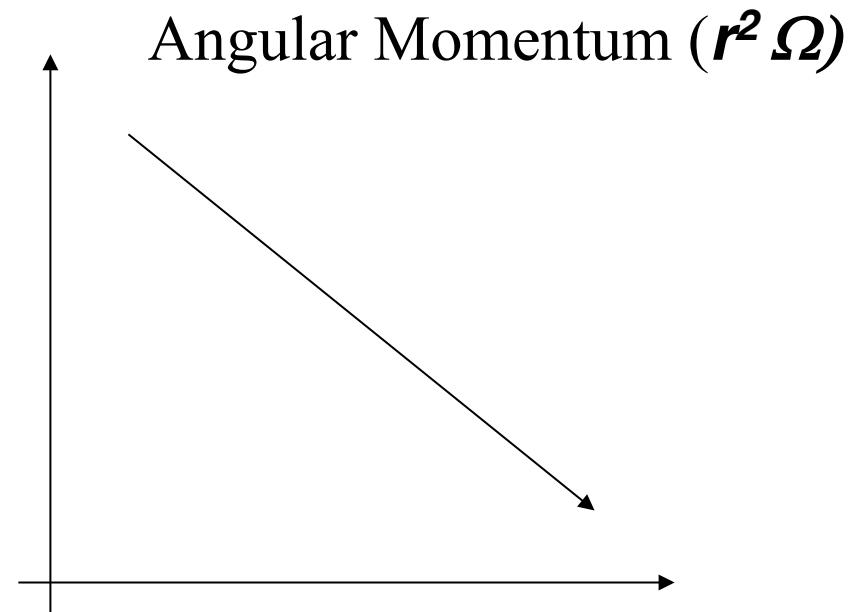
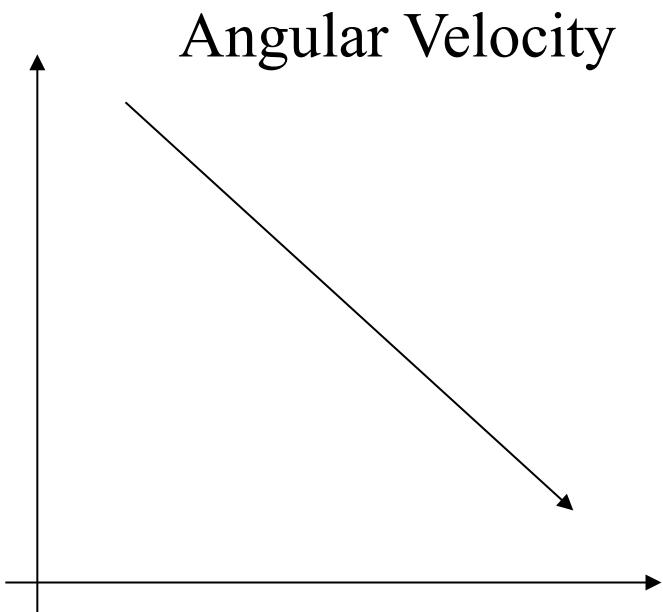
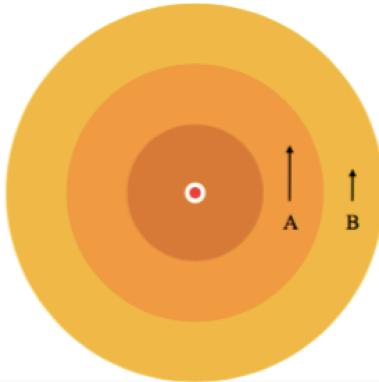


# Stability (Rayleigh criterion)



$$|d \ln \Omega / d \ln r| < 2$$

# Rayleigh instability

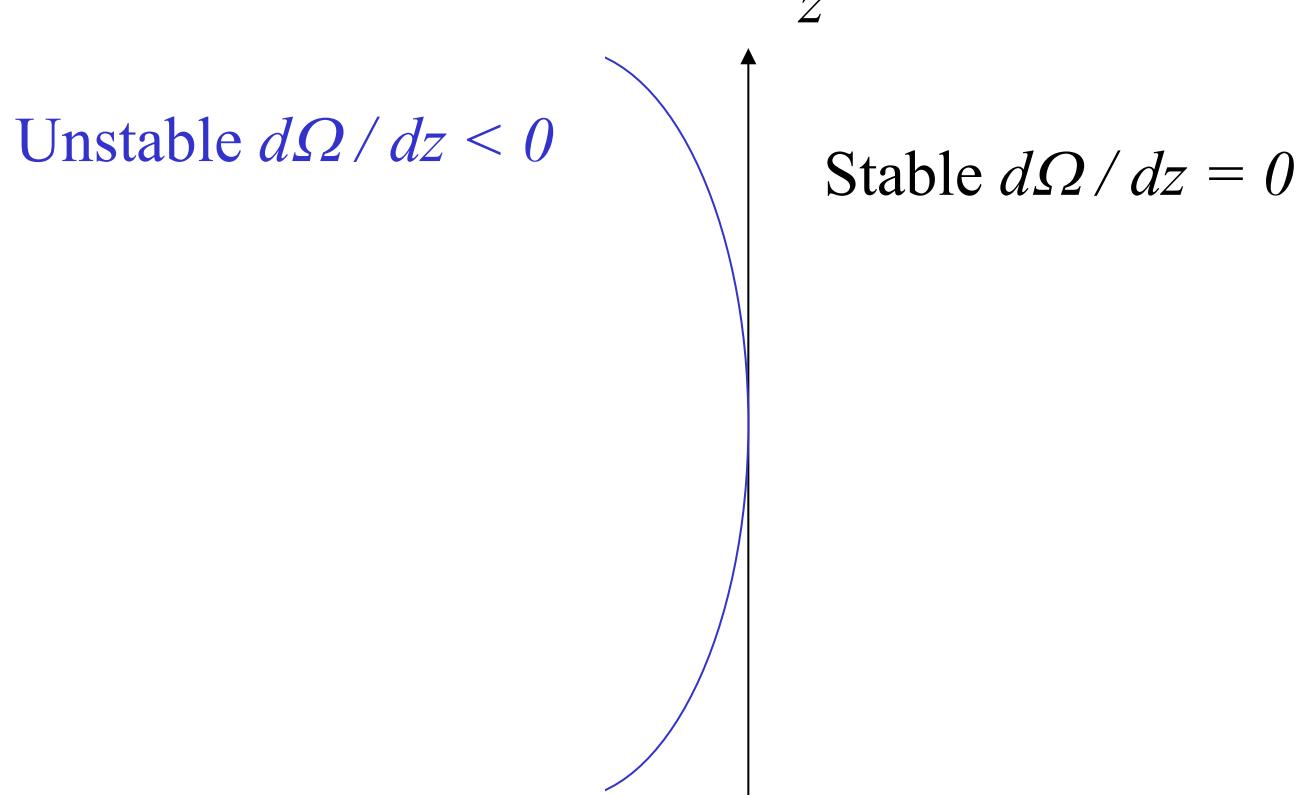


$$|d \ln \Omega / d \ln r | > 2$$

# Vertical shear instability

$$\rho_{\text{mid}} = \rho_0 \left( \frac{R}{R_0} \right)^p, \quad \Omega = \Omega_K \left[ 1 + \frac{1}{2} \left( \frac{H}{R} \right)^2 \left( p + q + \frac{q}{2} \frac{Z^2}{H^2} \right) \right]$$
$$c_s^2 = c_0^2 \left( \frac{R}{R_0} \right)^q,$$

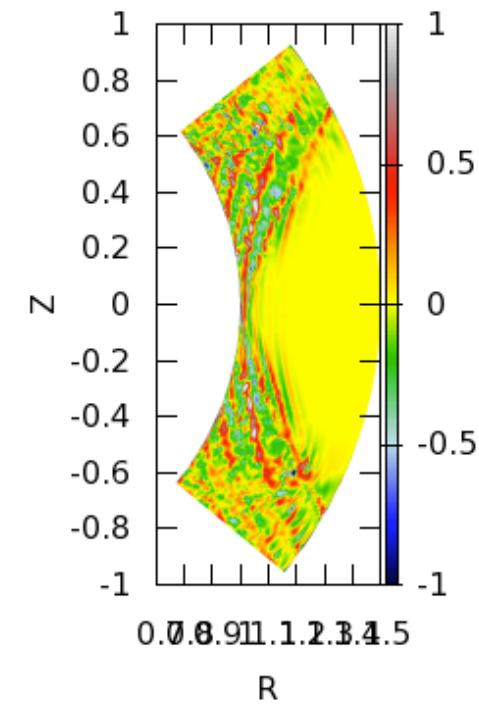
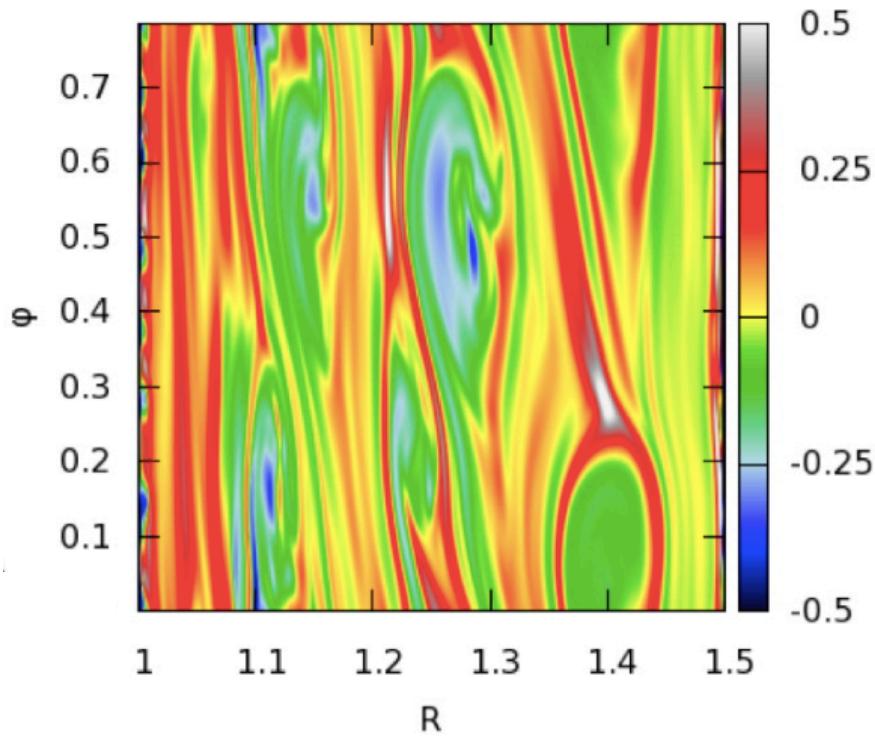
→



# Vertical shear instability

$$\rho_{\text{mid}} = \rho_0 \left( \frac{R}{R_0} \right)^p, \quad \Omega = \Omega_K \left[ 1 + \frac{1}{2} \left( \frac{H}{R} \right)^2 \left( p + q + \frac{q}{2} \frac{Z^2}{H^2} \right) \right]$$
$$c_s^2 = c_0^2 \left( \frac{R}{R_0} \right)^q,$$

→



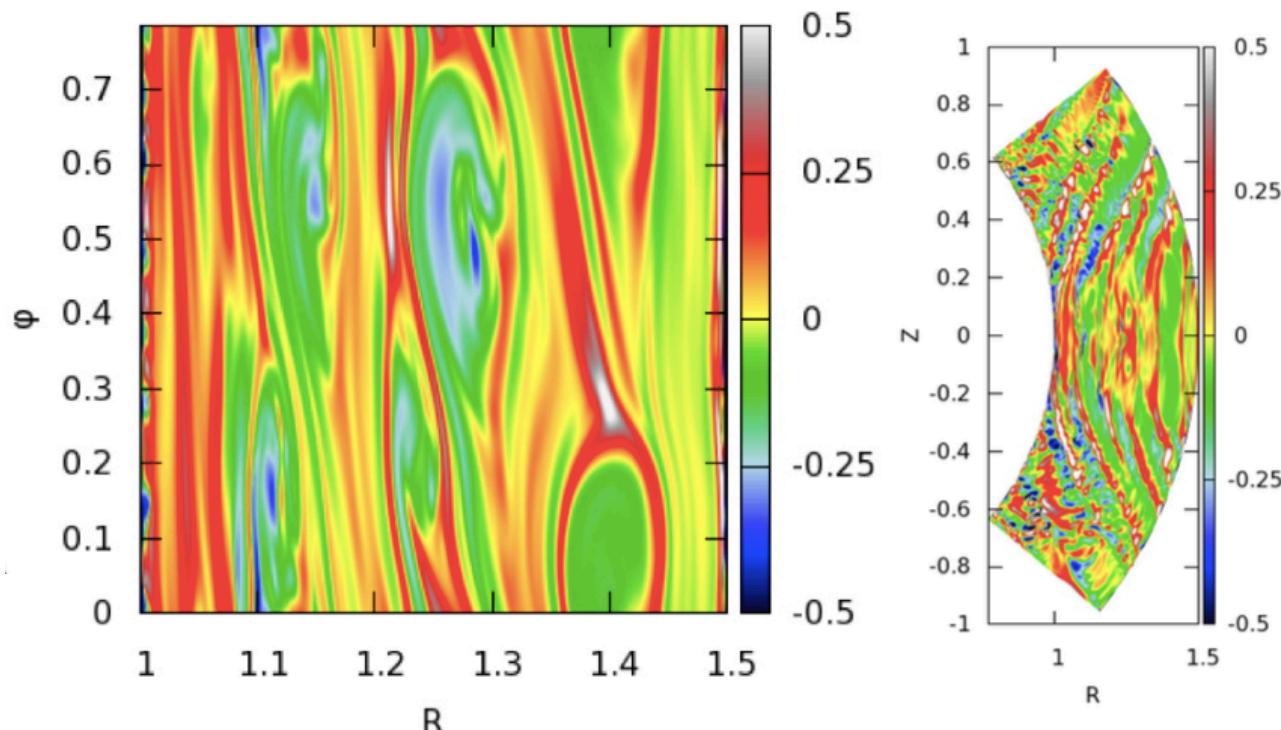
# Vertical shear instability

$$\rho_{\text{mid}} = \rho_0 \left( \frac{R}{R_0} \right)^p, \quad \Omega = \Omega_K \left[ 1 + \frac{1}{2} \left( \frac{H}{R} \right)^2 \left( p + q + \frac{q}{2} \frac{Z^2}{H^2} \right) \right]$$

$$c_s^2 = c_0^2 \left( \frac{R}{R_0} \right)^q,$$

$d\Omega / dz \neq 0 ; \kappa_z^2 < 0 \Rightarrow$  Rayleigh unstable

Solberg-Hoiland stability criterion  
 $\kappa^2 + N^2 > 0$



# Vertical shear instability

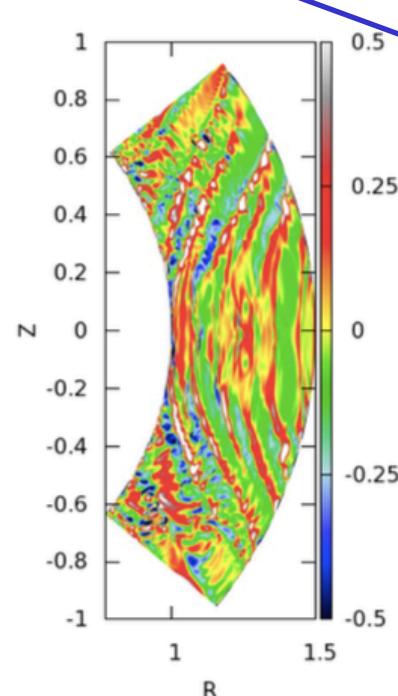
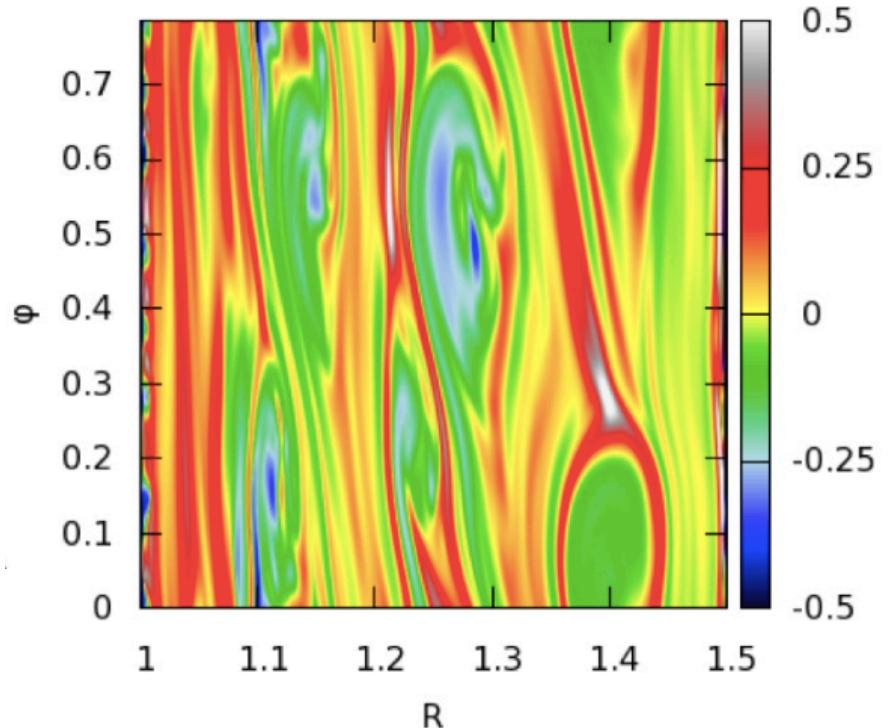
$$\rho_{\text{mid}} = \rho_0 \left( \frac{R}{R_0} \right)^p, \quad \Omega = \Omega_K \left[ 1 + \frac{1}{2} \left( \frac{H}{R} \right)^2 \left( p + q + \frac{q}{2} \frac{Z^2}{H^2} \right) \right]$$

$$c_s^2 = c_0^2 \left( \frac{R}{R_0} \right)^q,$$

$d\Omega / dz \neq 0 ; \kappa_z^2 < 0 \Rightarrow$  Rayleigh unstable

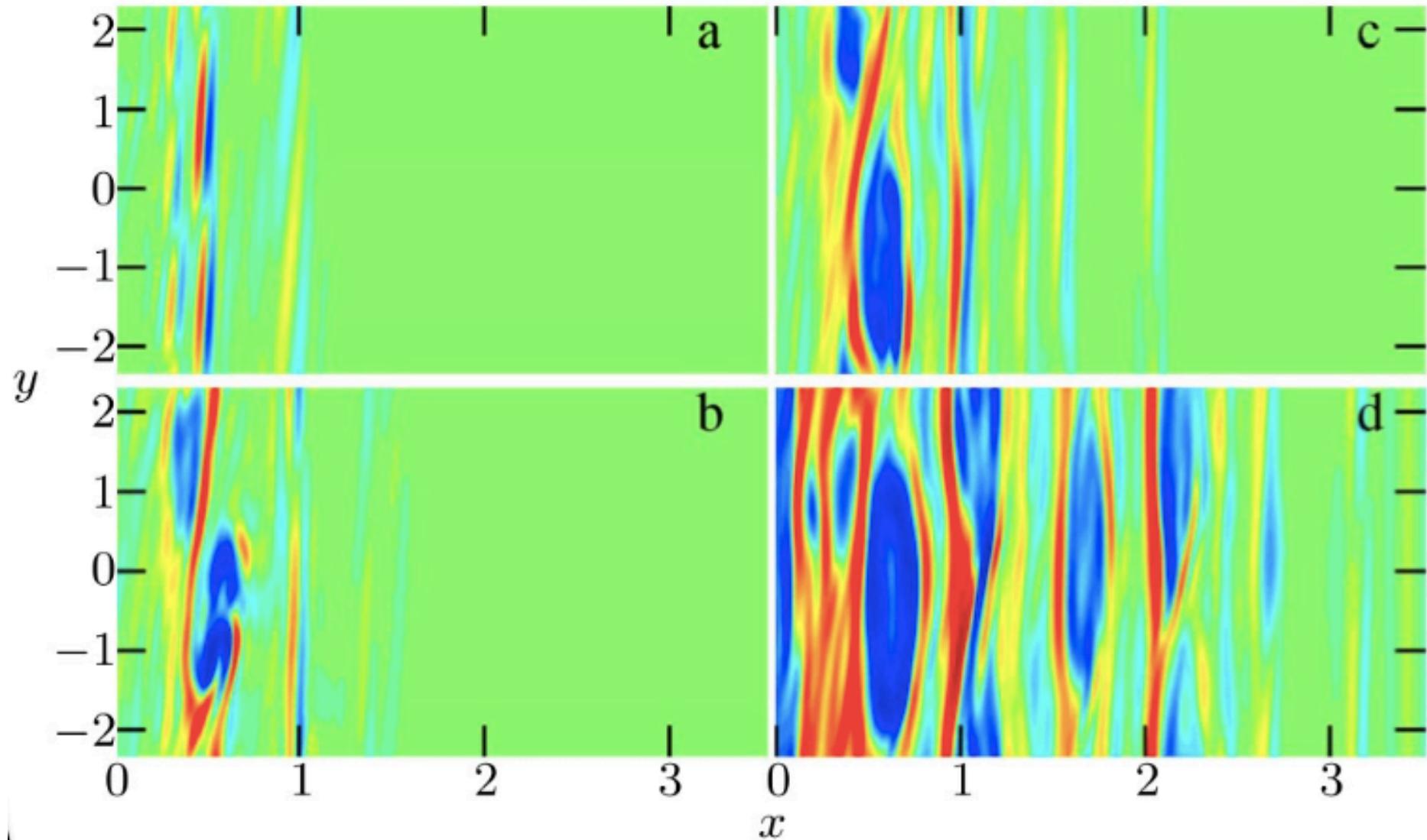
Solberg-Hoiland stability criterion

$$\kappa^2 + N^2 > 0$$



Buoyancy stabilizes!

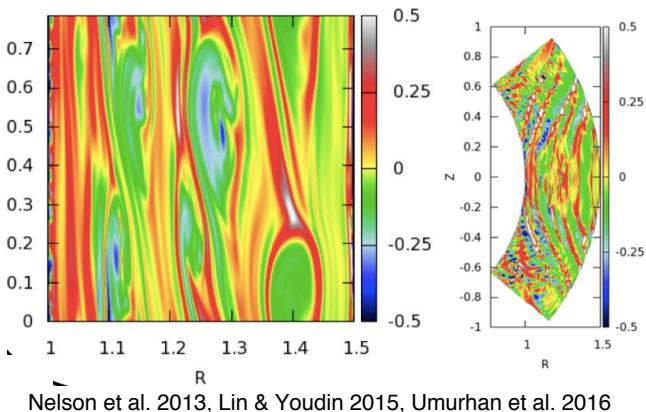
# Zombie Vortex Instability



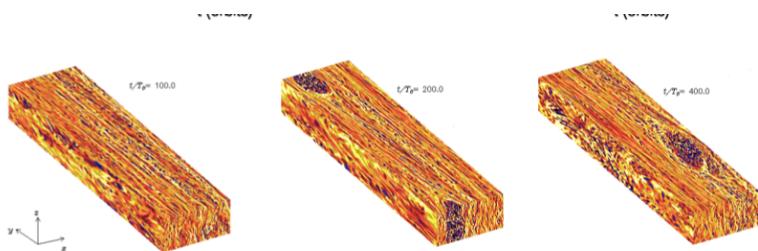
Cascade of baroclinic critical layers

# Thermal Instabilities

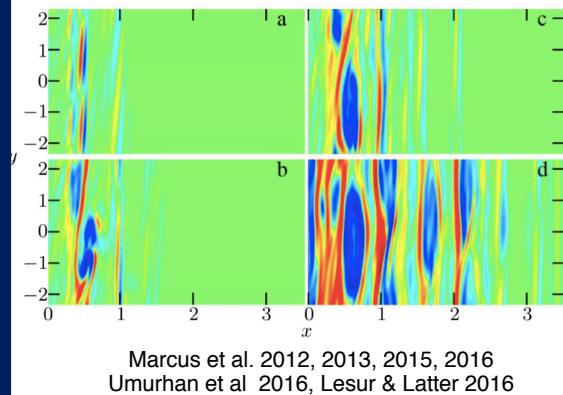
## Vertical Shear Instability



## Convective Overstability



## Zombie Vortex Instability



$\Omega\tau \ll 1$   
( $\kappa < 1 \text{ cm}^2/\text{g}$ )

$\Omega\tau \sim 1$   
( $\kappa \sim 1\text{--}50 \text{ cm}^2/\text{g}$ )

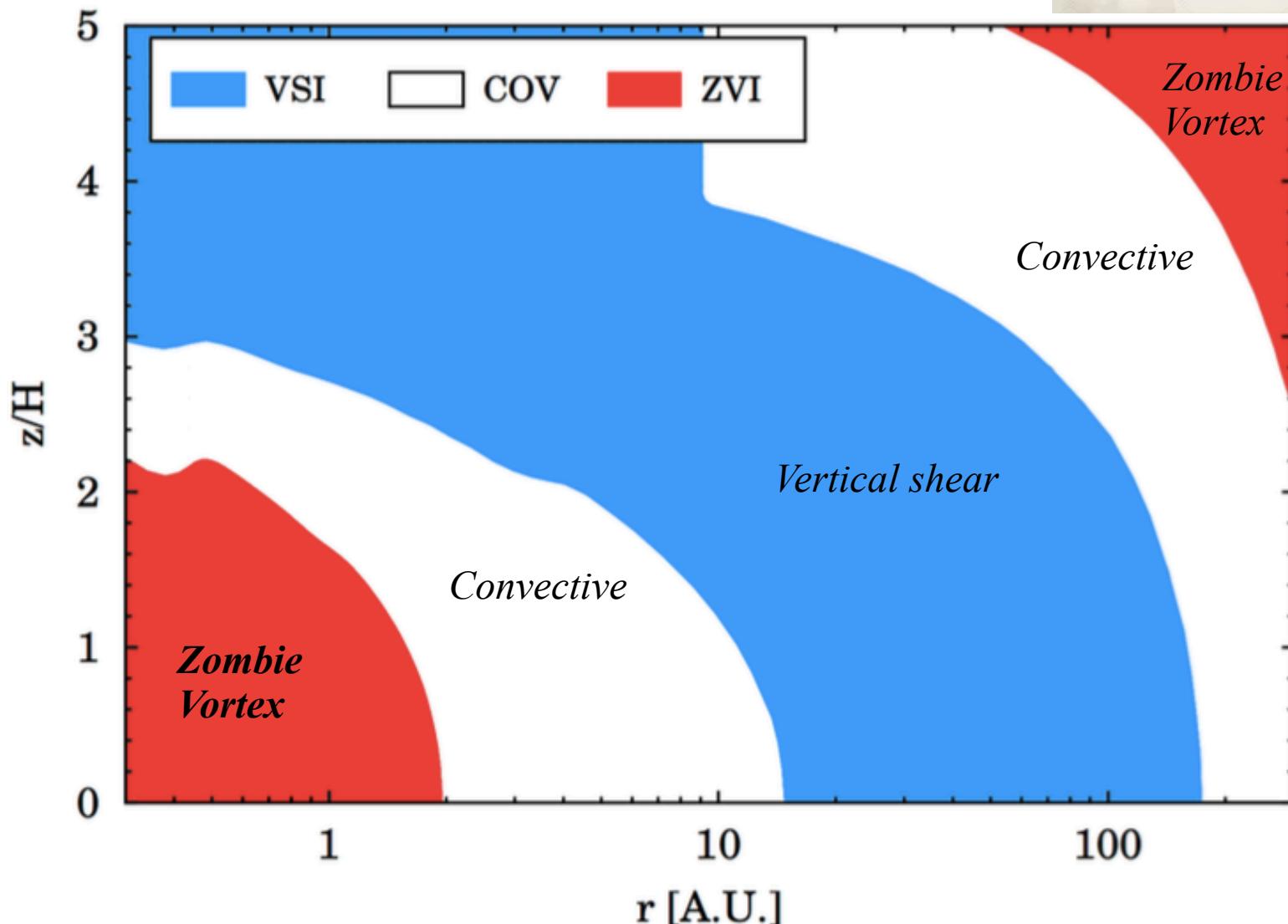
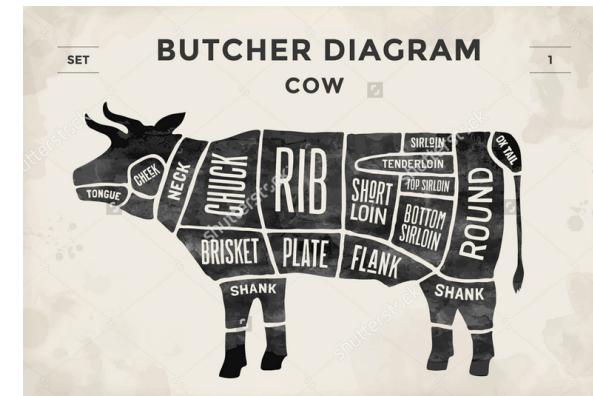
$\Omega\tau \gg 1$   
( $\kappa > 50 \text{ cm}^2/\text{g}$ )

*Opacity*



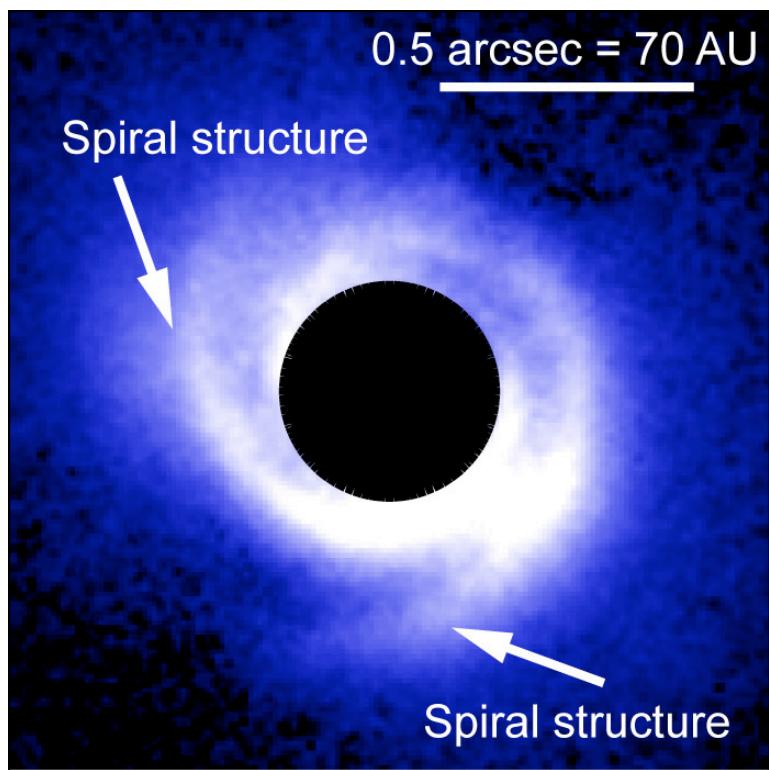
# Synthesis

A “butcher diagram” for hydro instabilities.



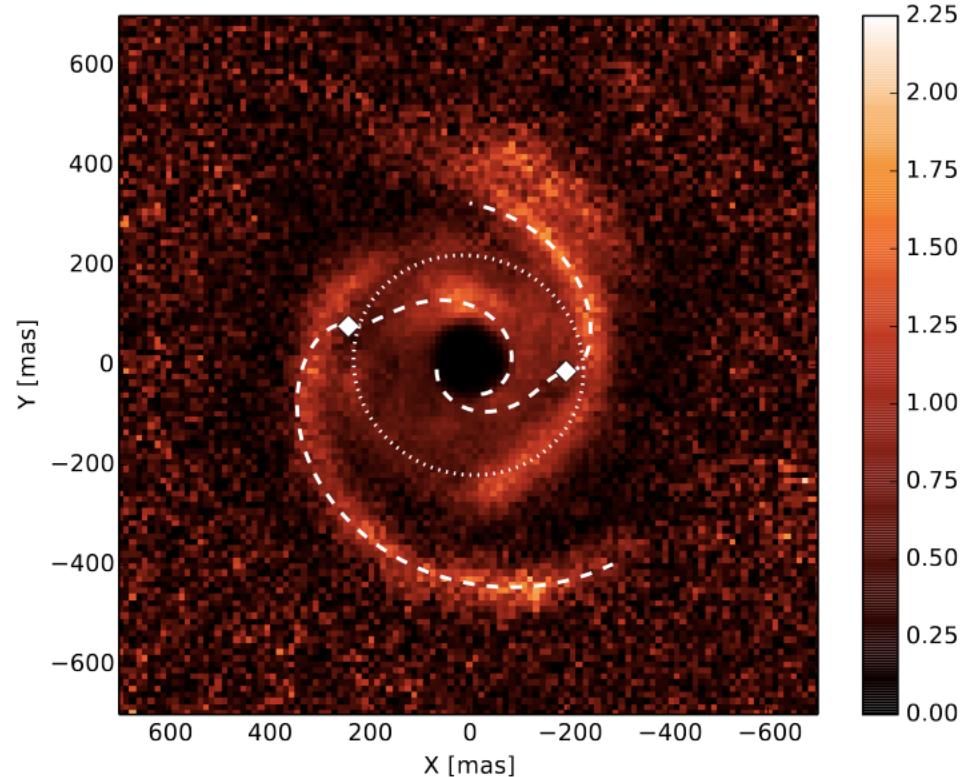
# Observational Evidence: Spirals

SAO 206462



Muto et al. (2012)

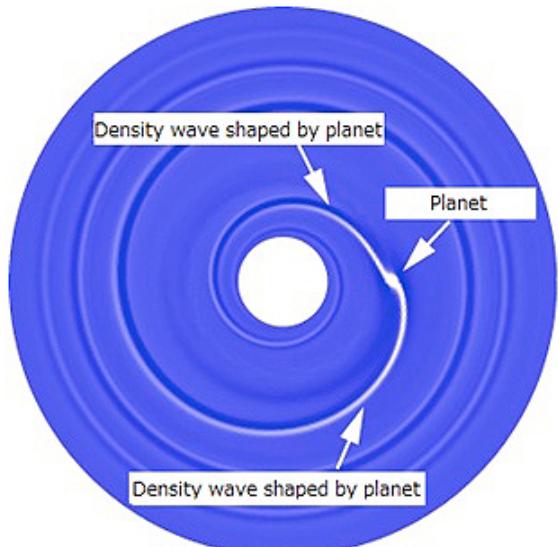
MWC 748



Benisty et al. (2015)

# Spiral arm fitting leads to problems

## Analytical spiral fit

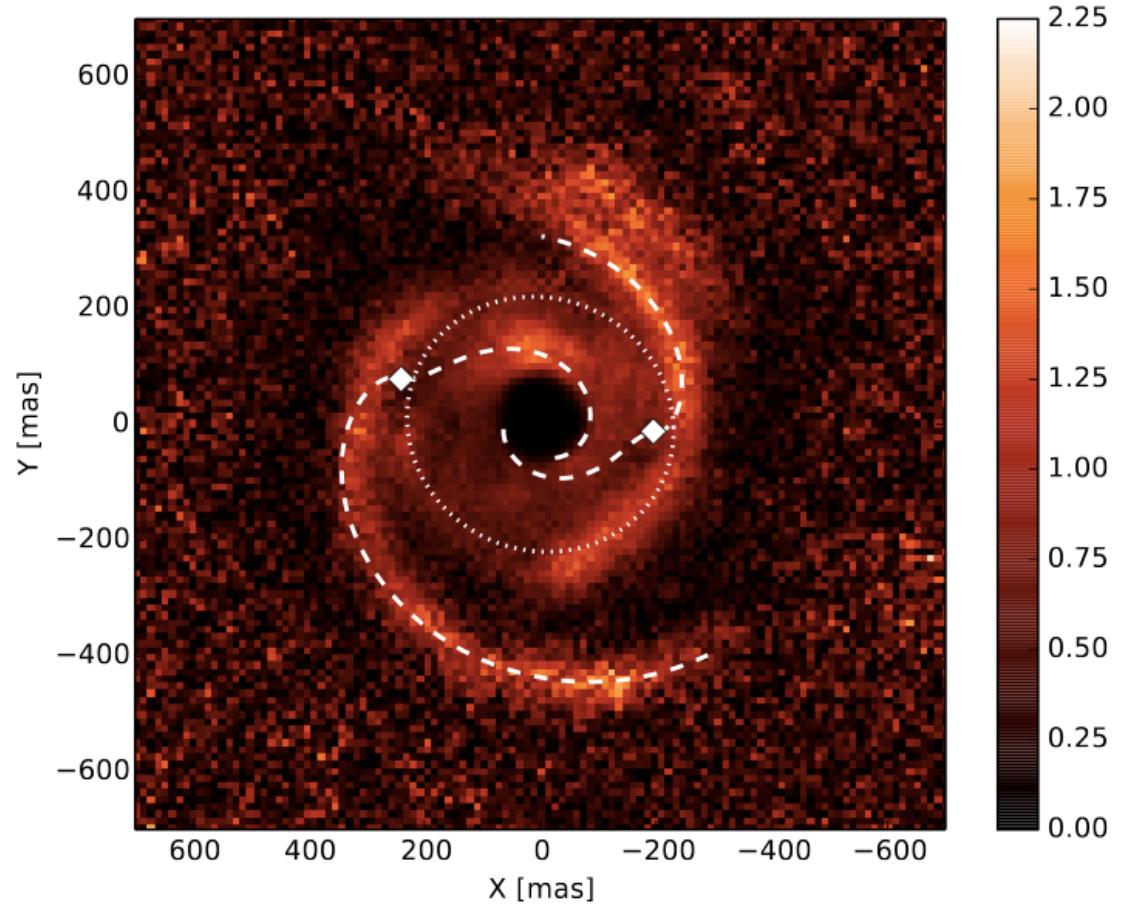


$$\theta(r) = \theta_c + \frac{\operatorname{sgn}(r - r_c)}{h_c} \times \left\{ \left( \frac{r}{r_c} \right)^{1+\beta} \left[ \frac{1}{1+\beta} - \frac{1}{1-\alpha+\beta} \left( \frac{r}{r_c} \right)^{-\alpha} \right] - \left( \frac{1}{1+\beta} - \frac{1}{1-\alpha+\beta} \right) \right\},$$

Rafikov (2002)

Muto et al. (2012)

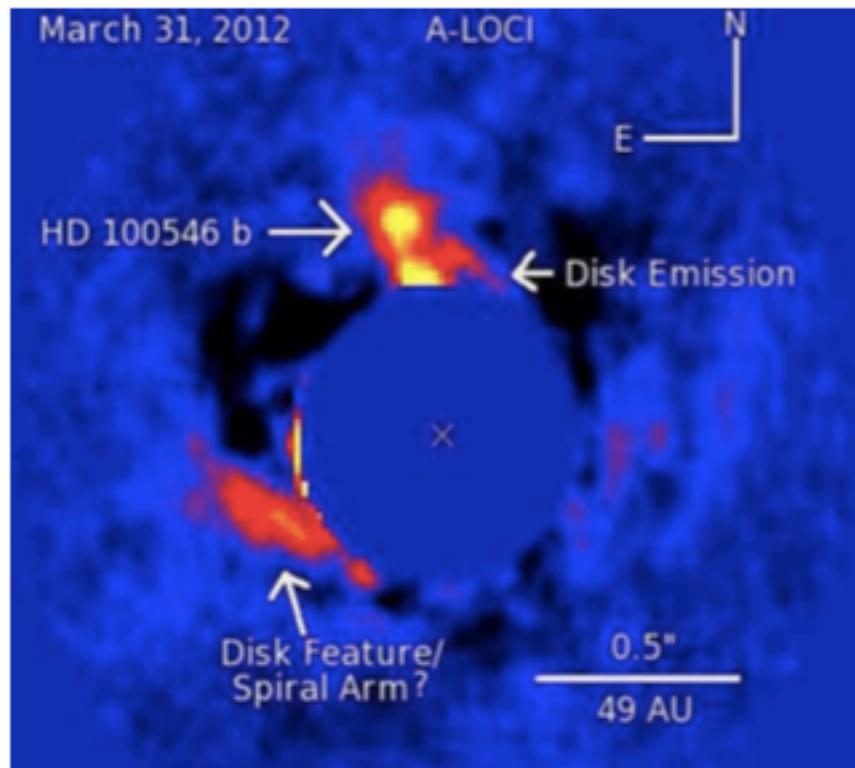
Spirals are **too wide**,  
**hotter** (300K) than ambient gas (50K).



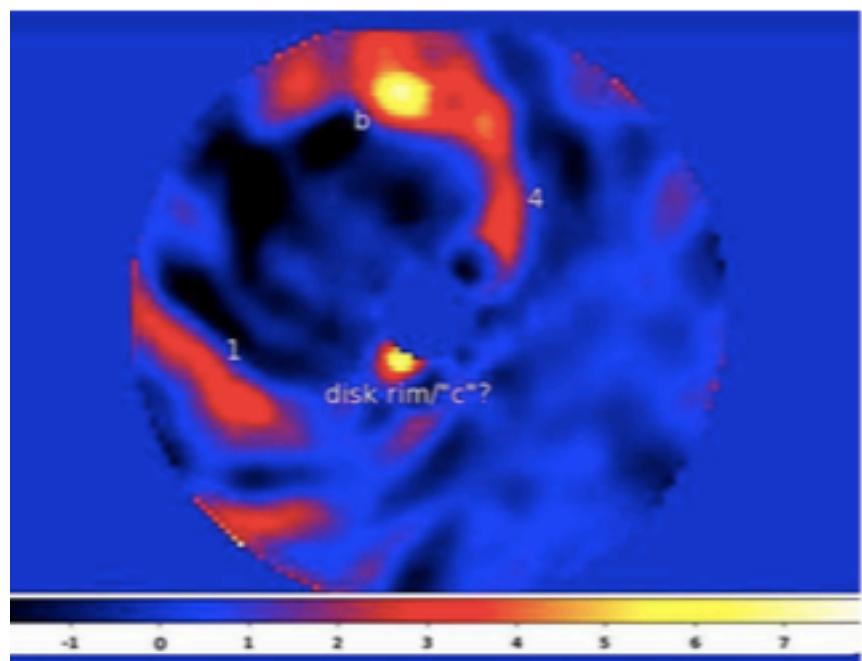
Benisty et al. (2015)

# The strange case of thermal emission in HD 100546

L band ( $\sim 3.5 \mu\text{m}$ )

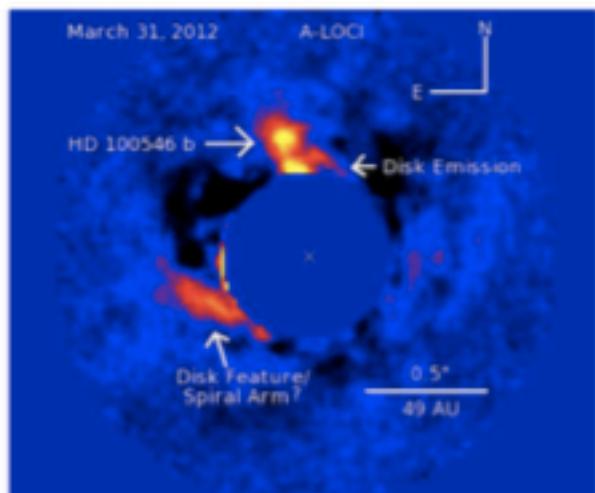


H band ( $\sim 1.6 \mu\text{m}$ )

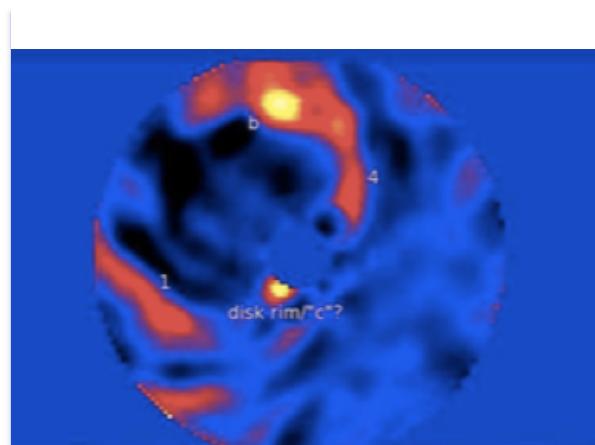


Currie et al. (2014), Currie et al. (2015)

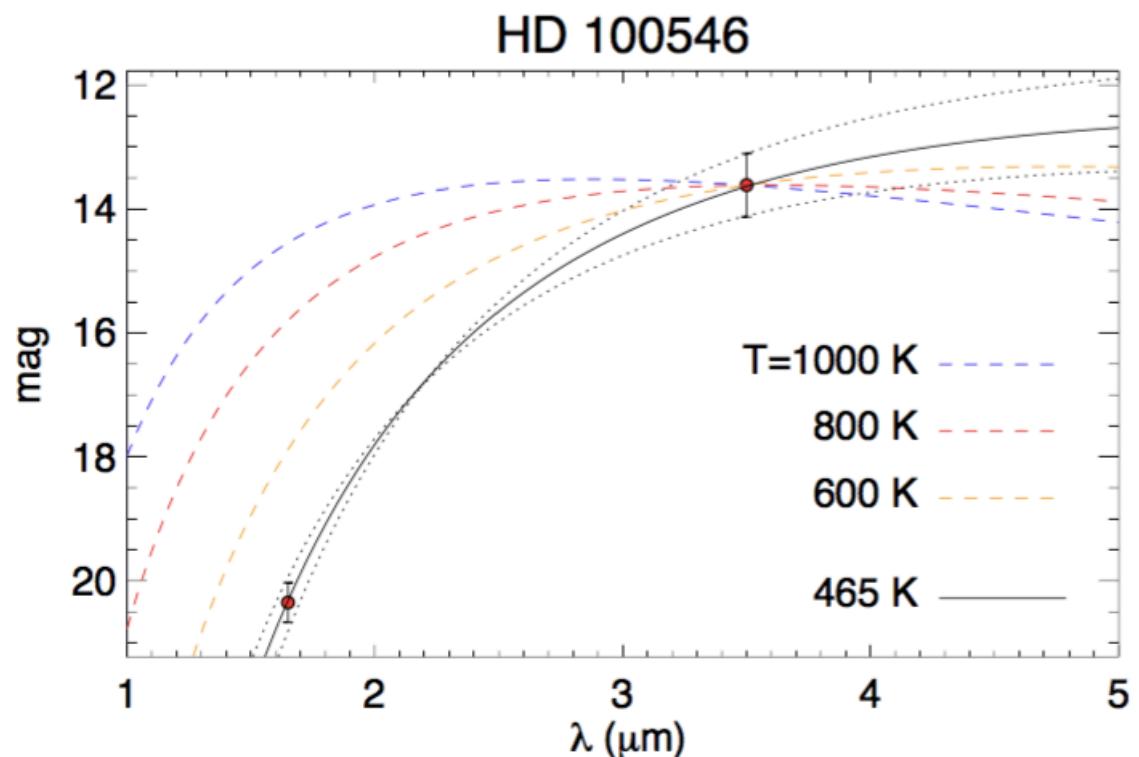
## Pinning down the temperature



L band

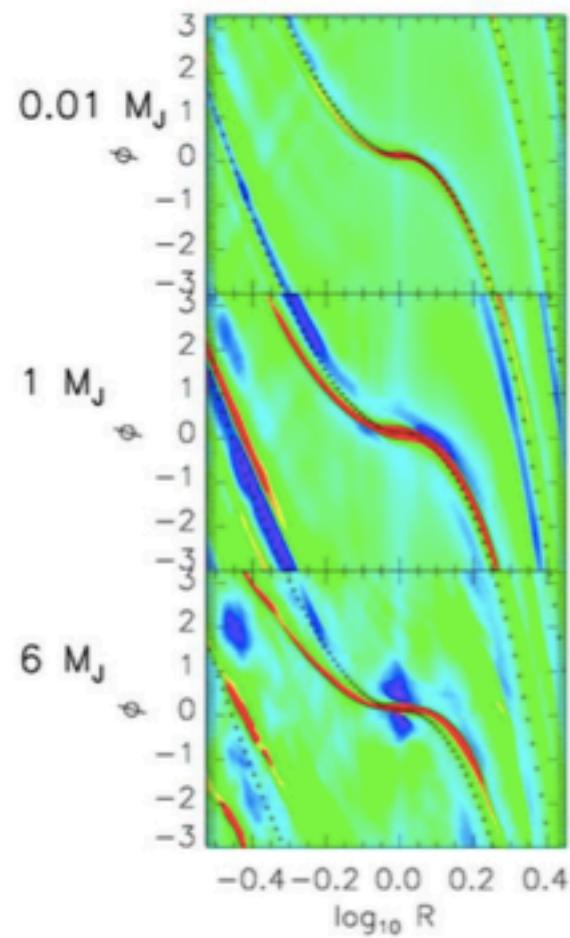


H band



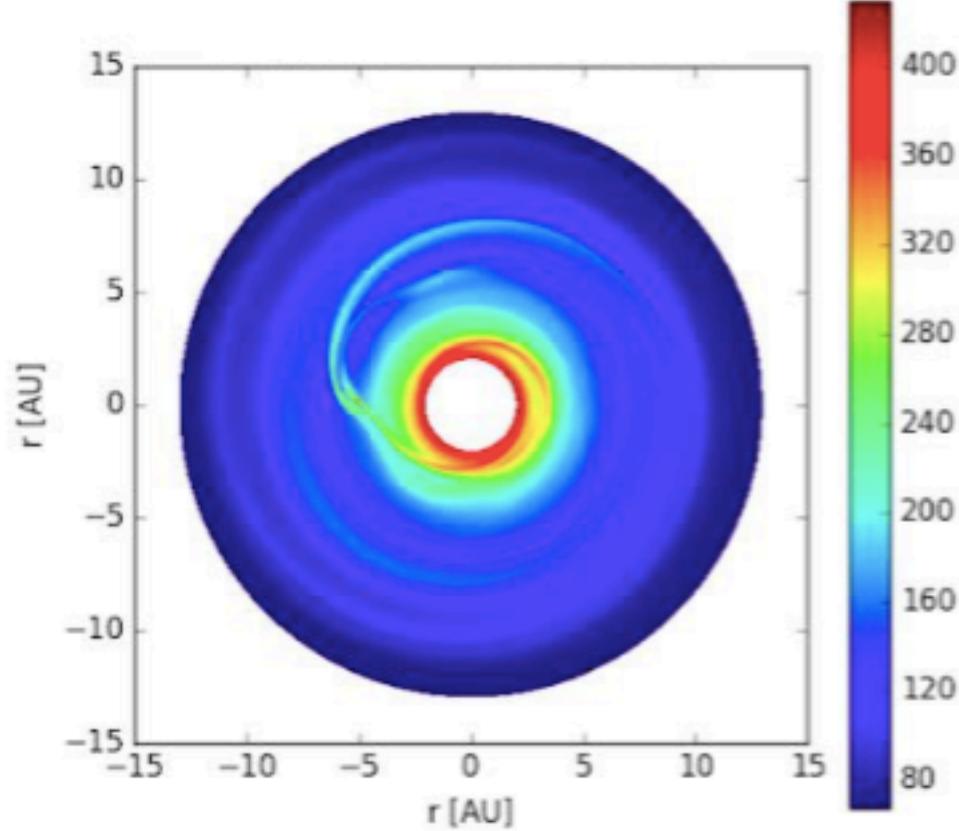
Lyra et al. (2016)

# Supersonic Wakes of High Mass Planets



Density

Zhu et al. (2015)

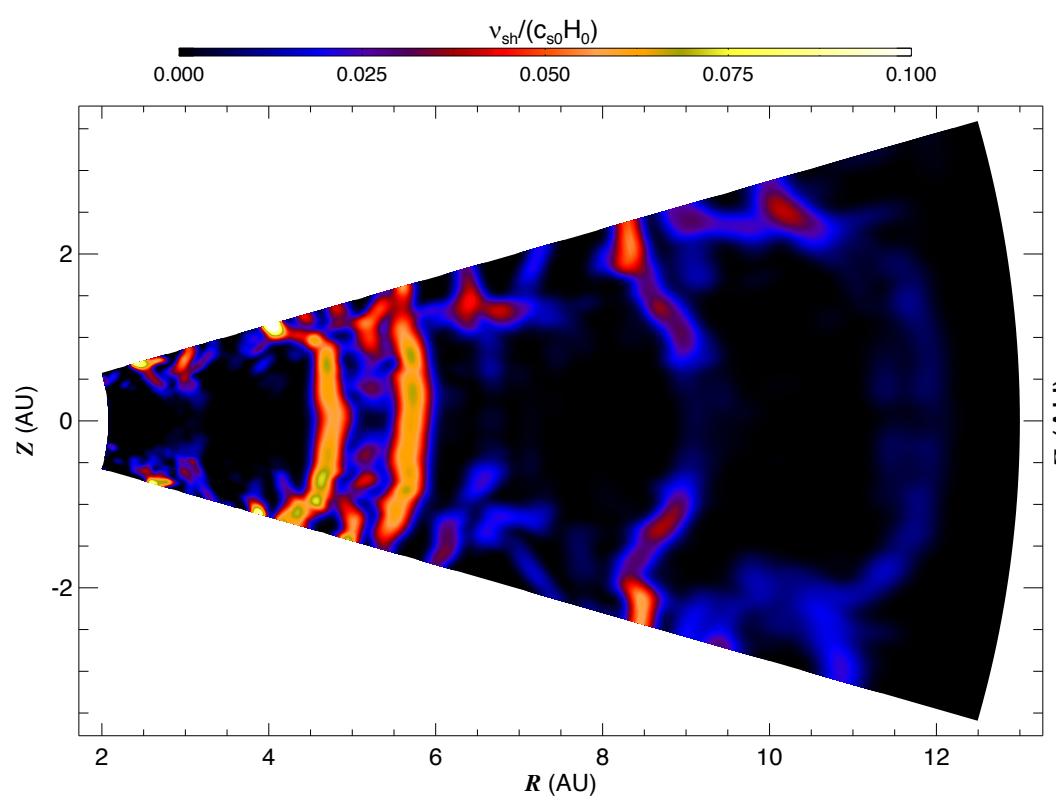


Temperature -  $5 M_J$

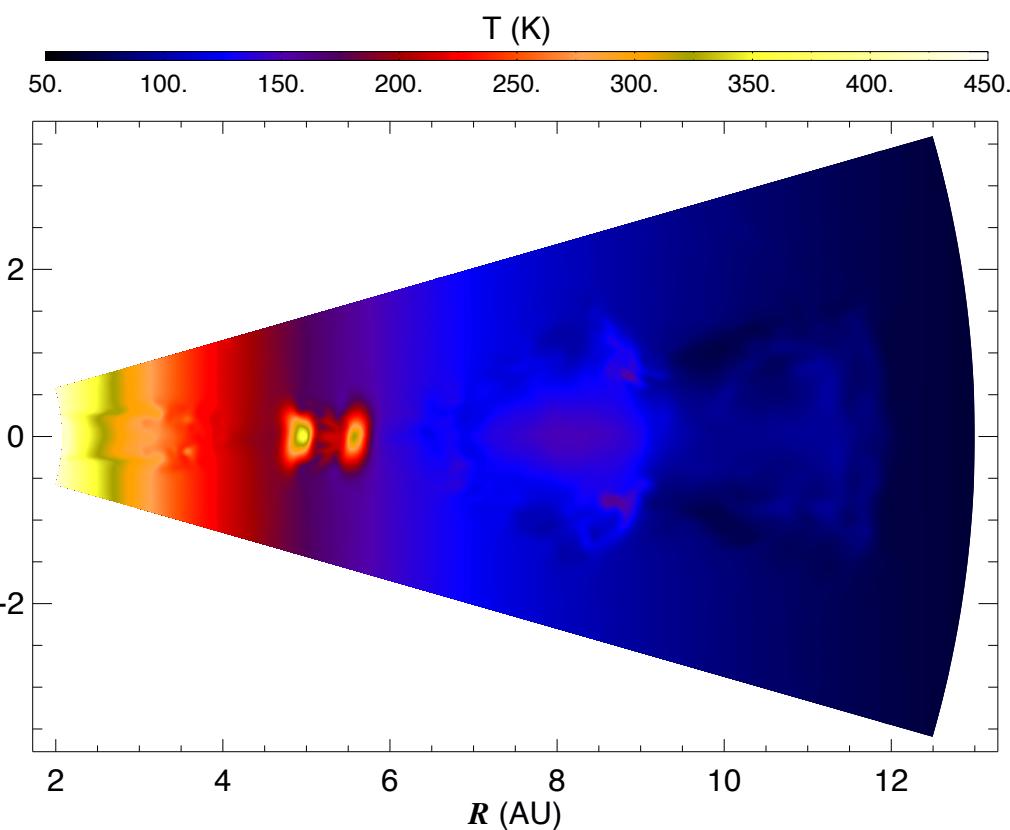
Lyra et al. (2016)

# Shock bores

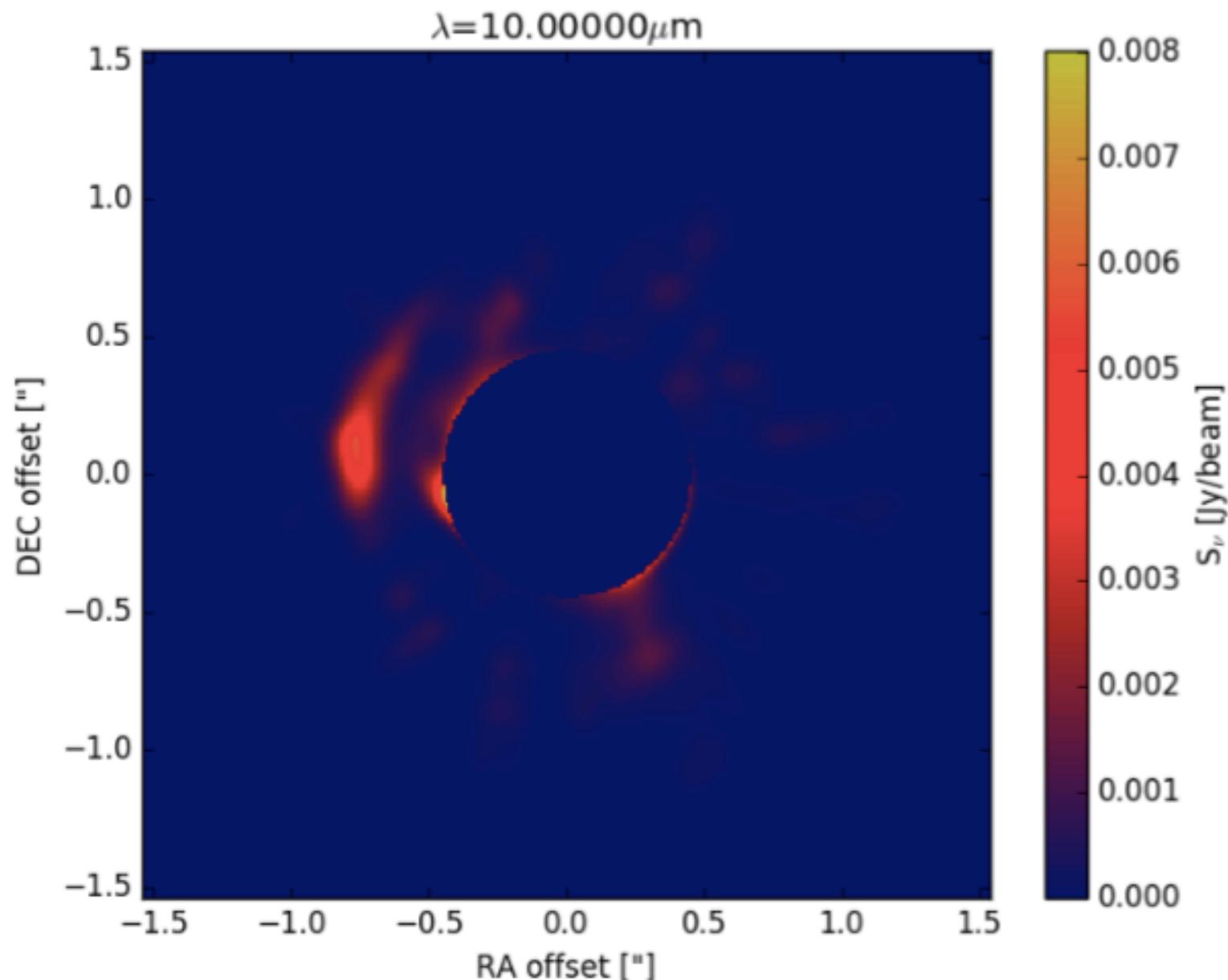
Shocks (velocity convergence)



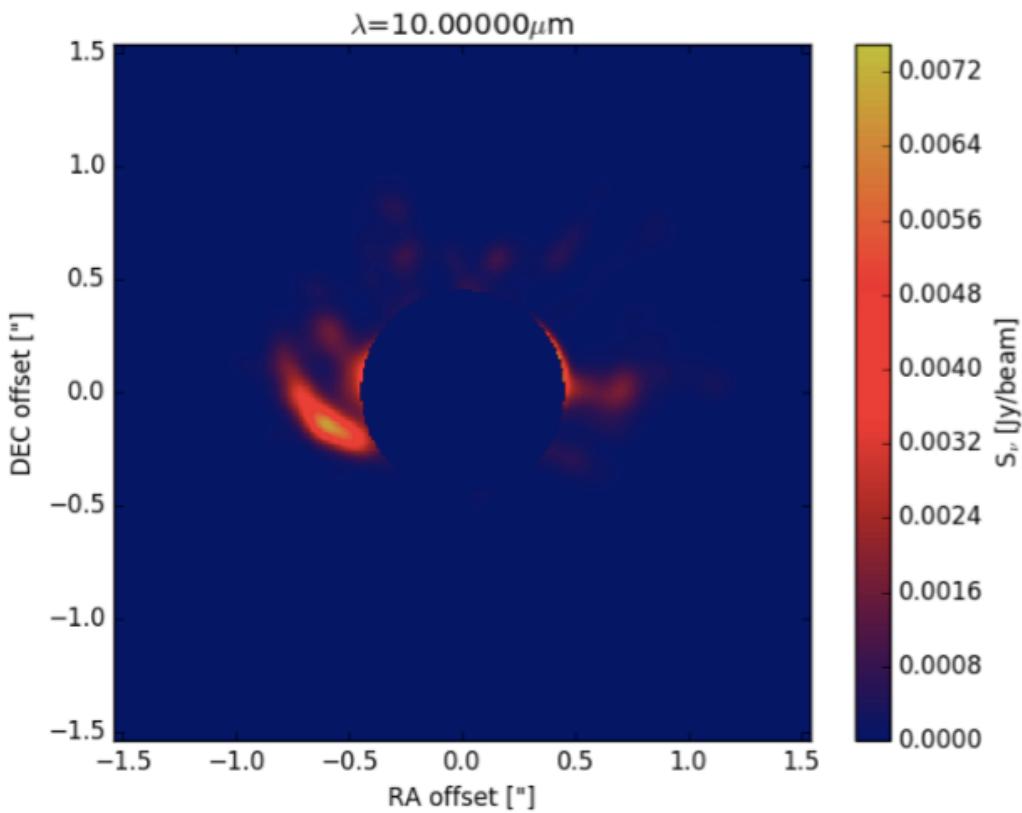
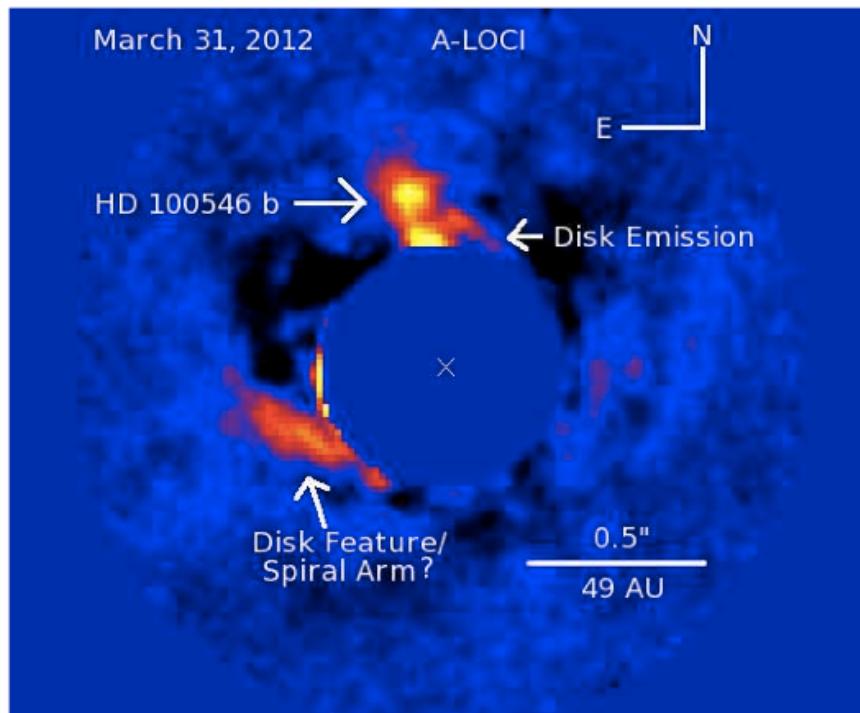
Temperature



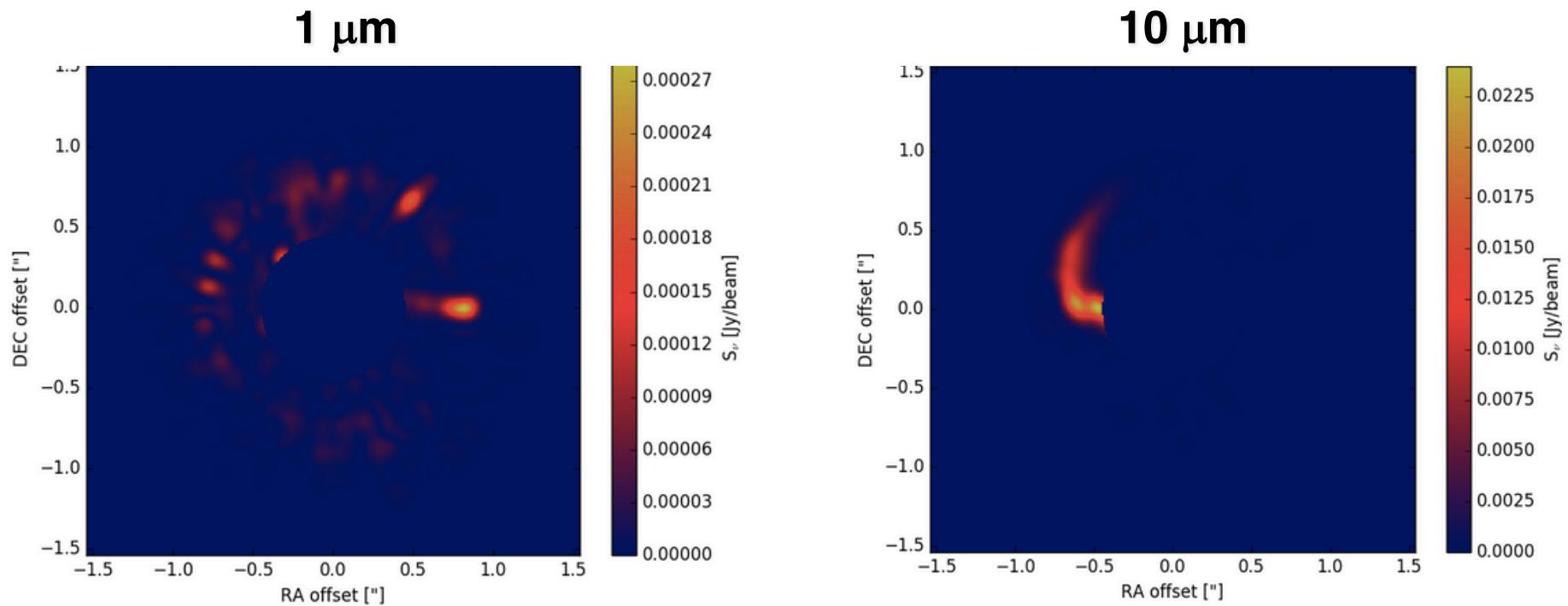
# Synthetic image



# Observation vs Synthetic Image

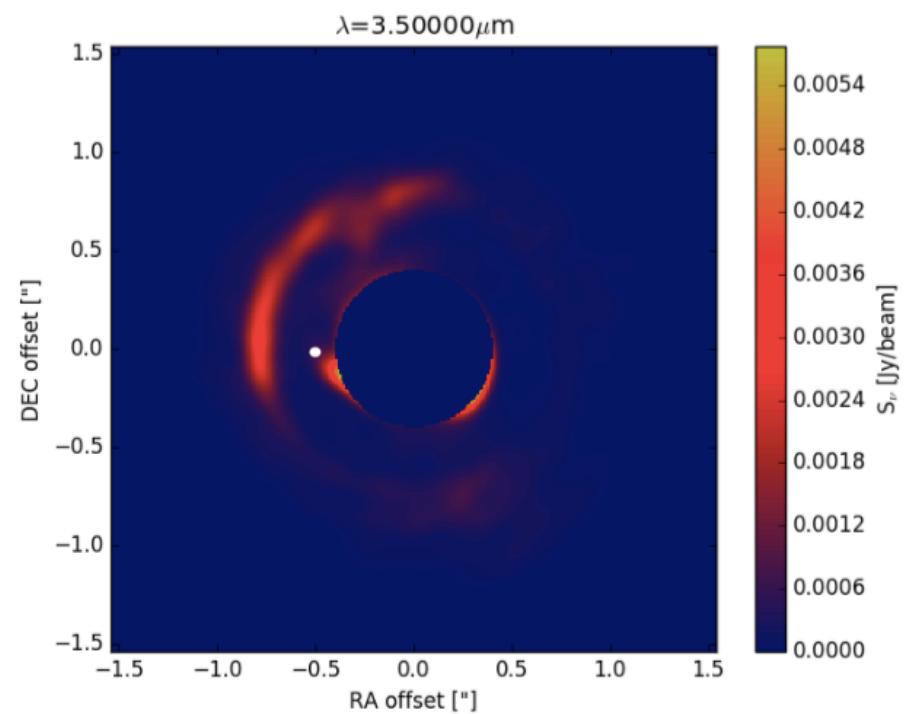
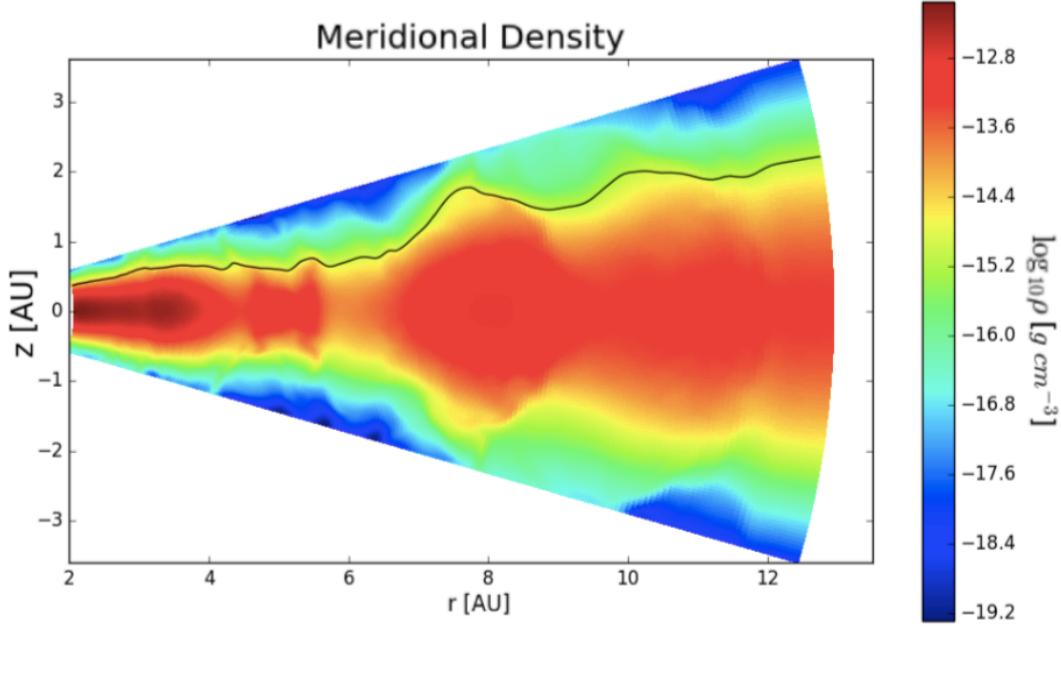


# Effect of shocks alone



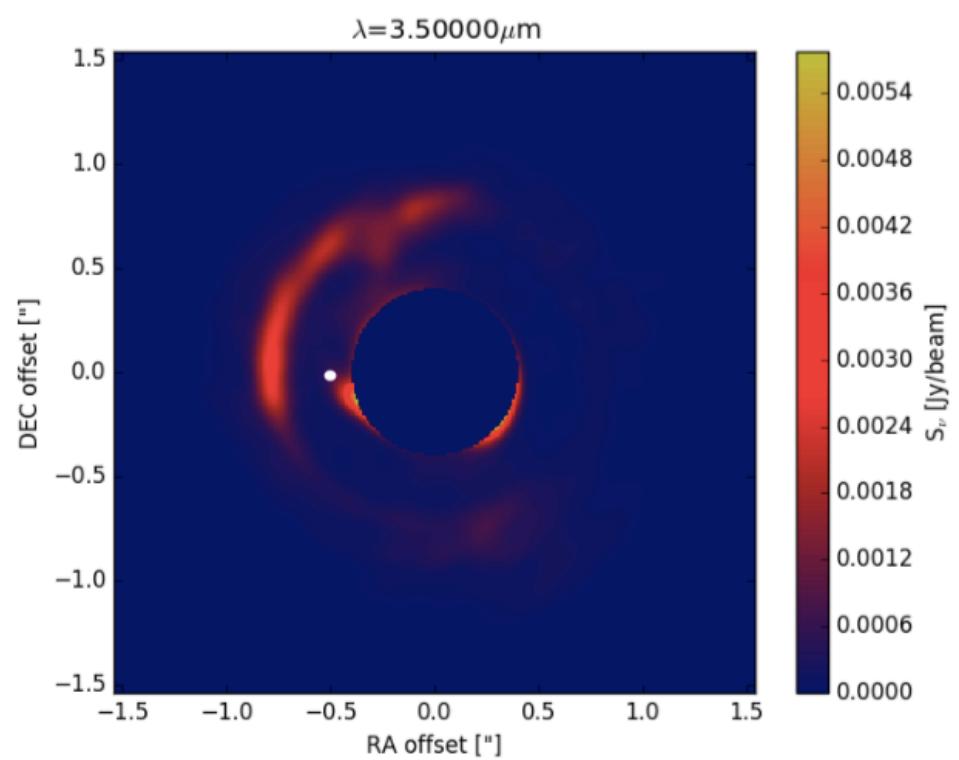
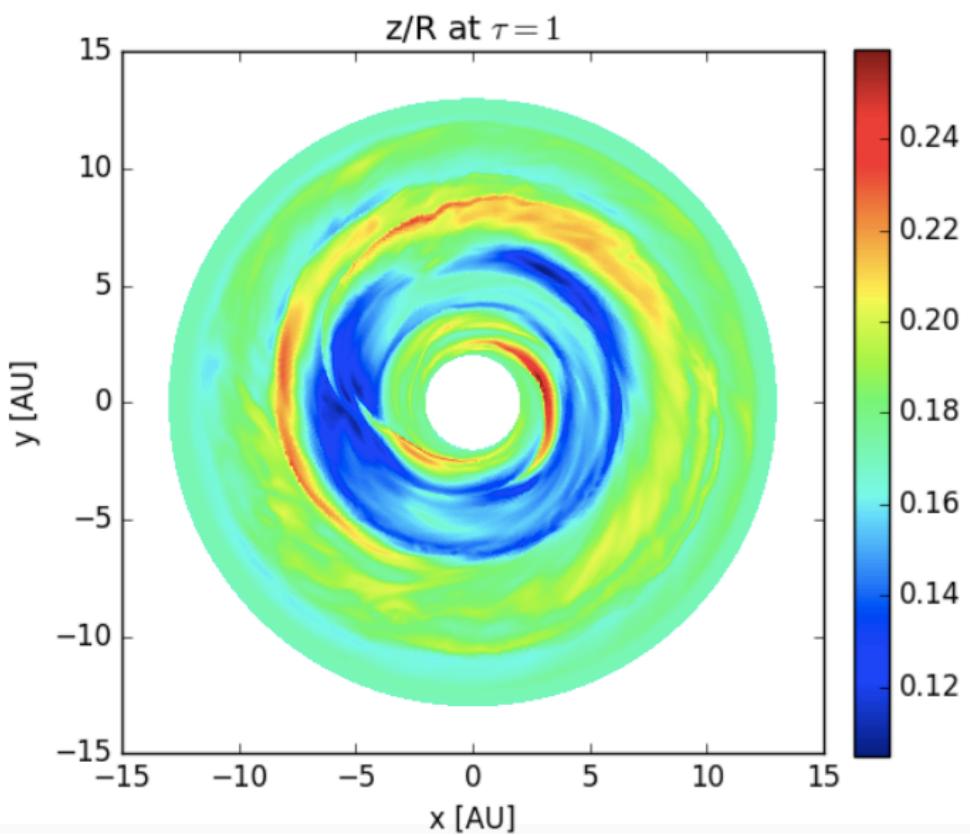
Hord et al. (2017)

# Scattering – A puffed up outer gap



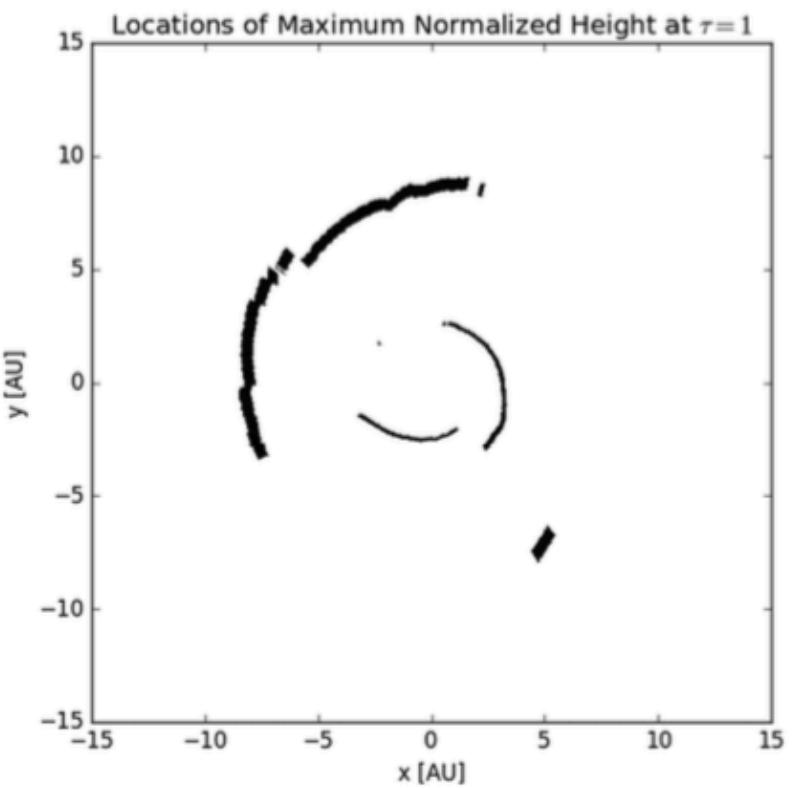
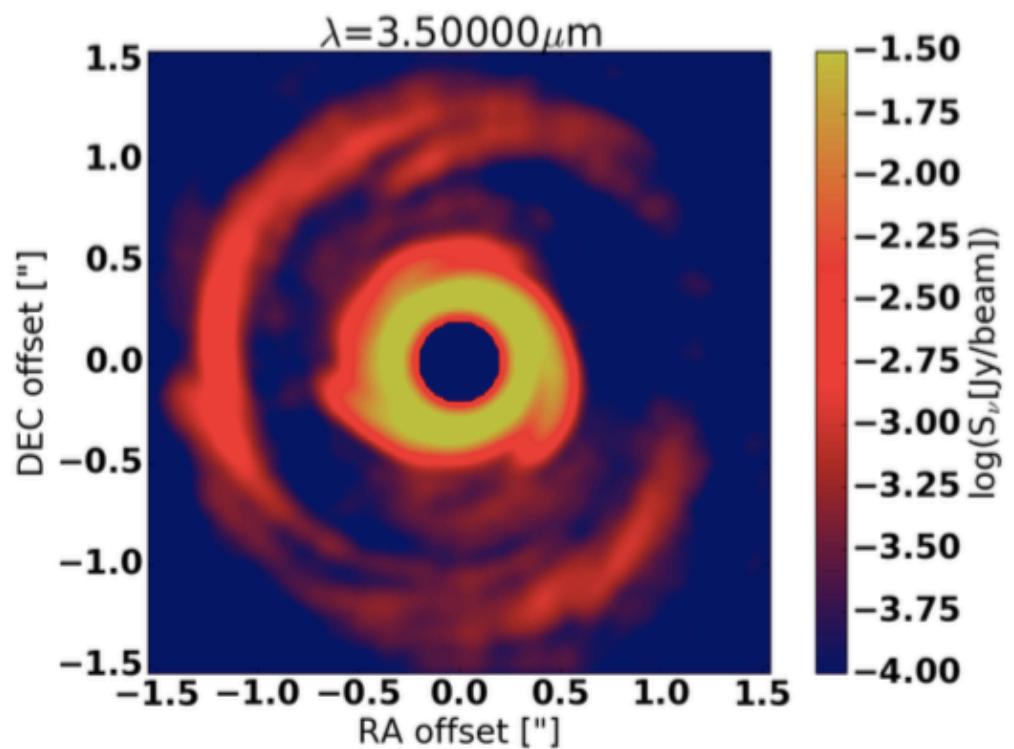
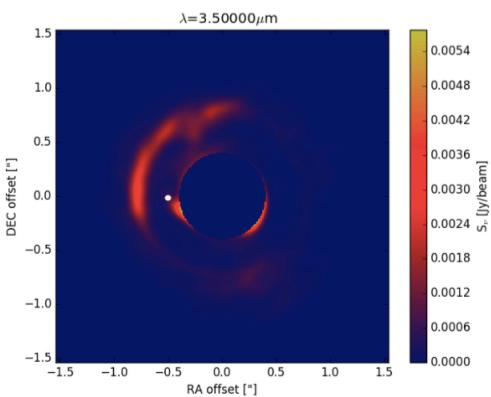
Hord et al. (2017)

# Scattering



Hord et al. (2017)

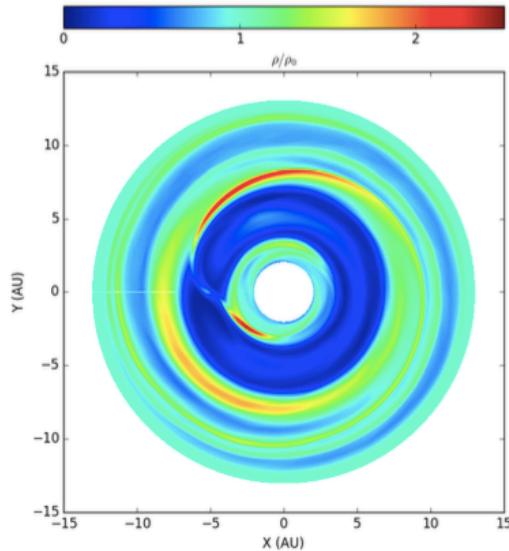
We see what is not in the  
shadow of the inner disk spirals



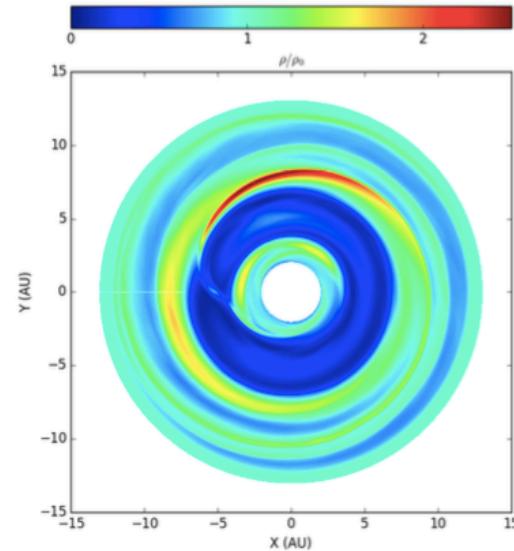
Hord et al. (2017)

# The pattern is stationary

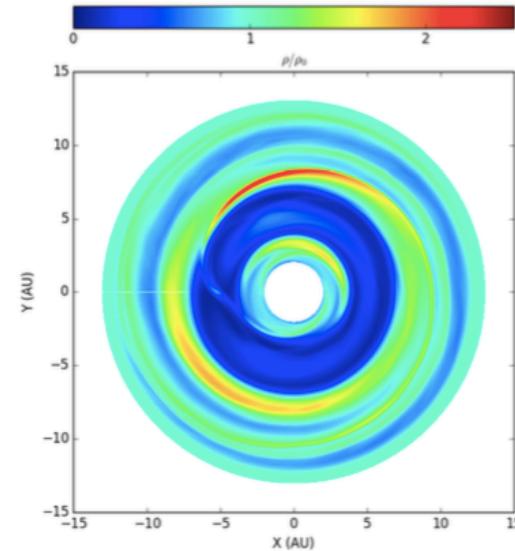
$T = 39$  orbits



$T = 40$  orbits

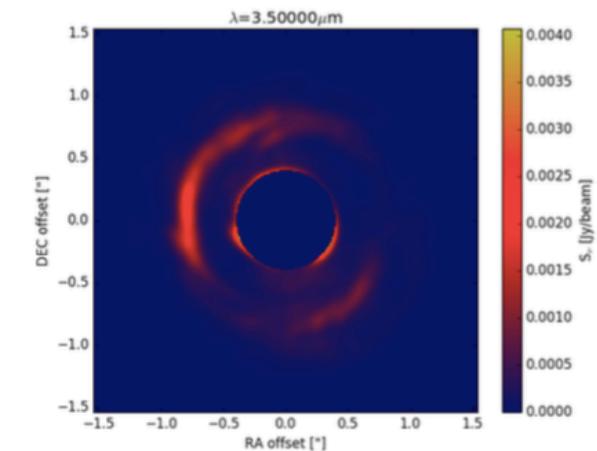
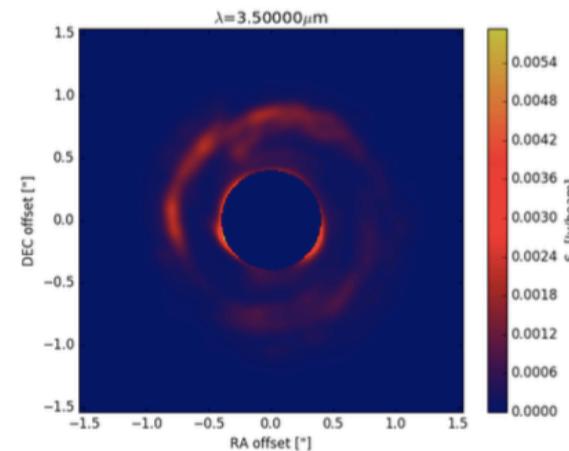
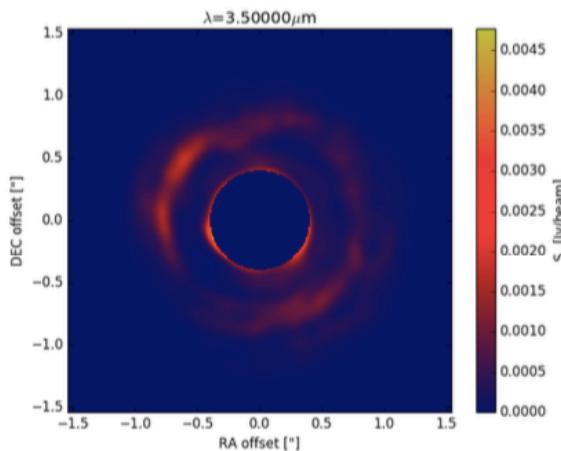


$T = 41$  orbits

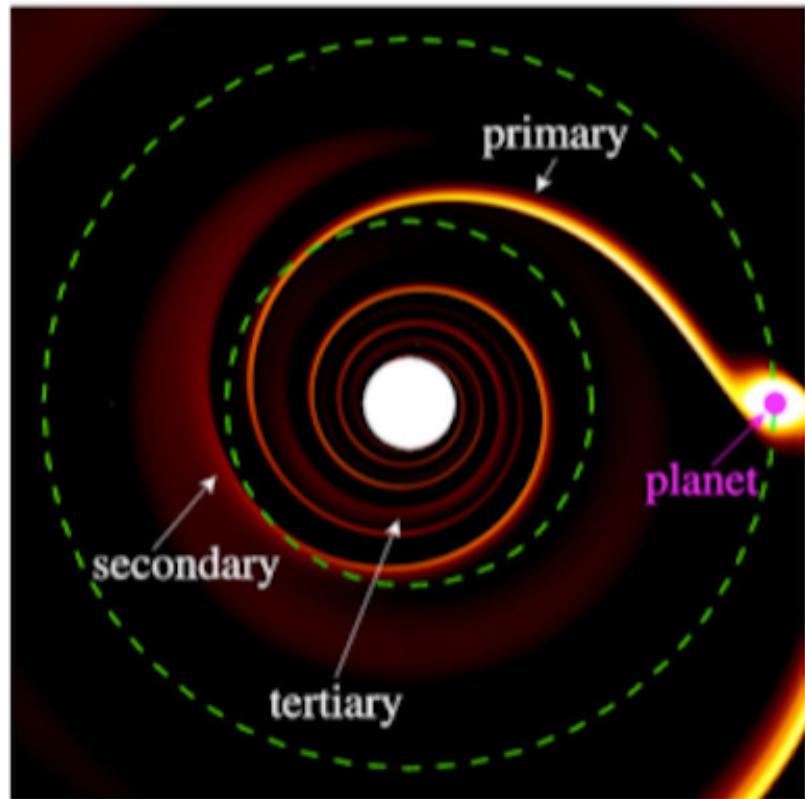


## Density

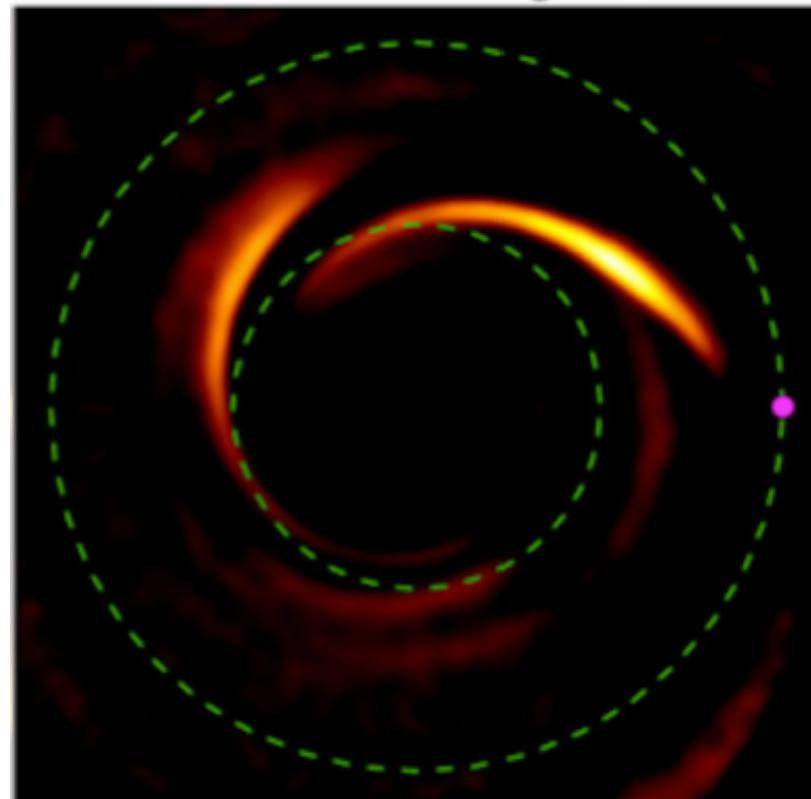
## Intensity



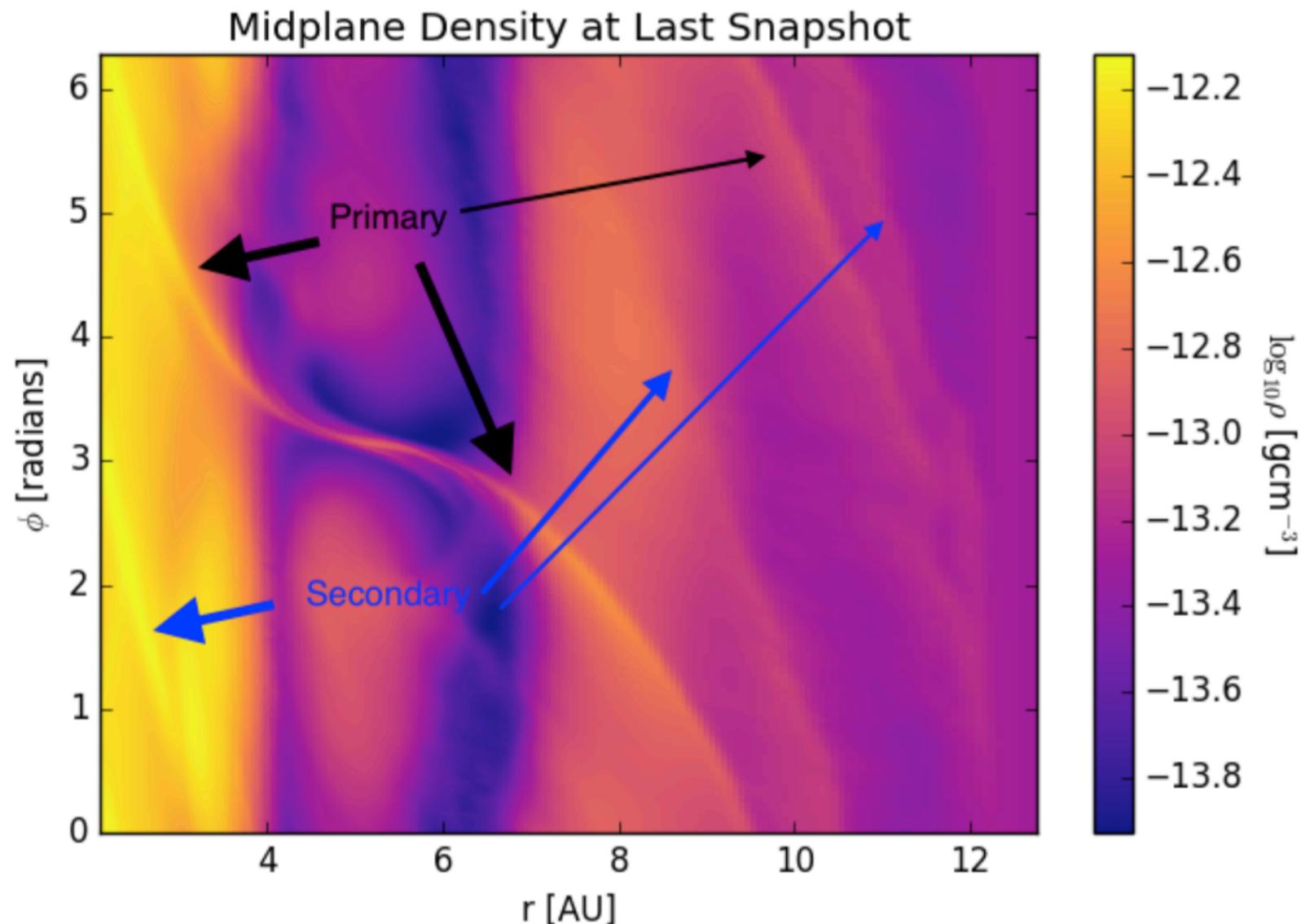
## Primary and Secondary spiral arms



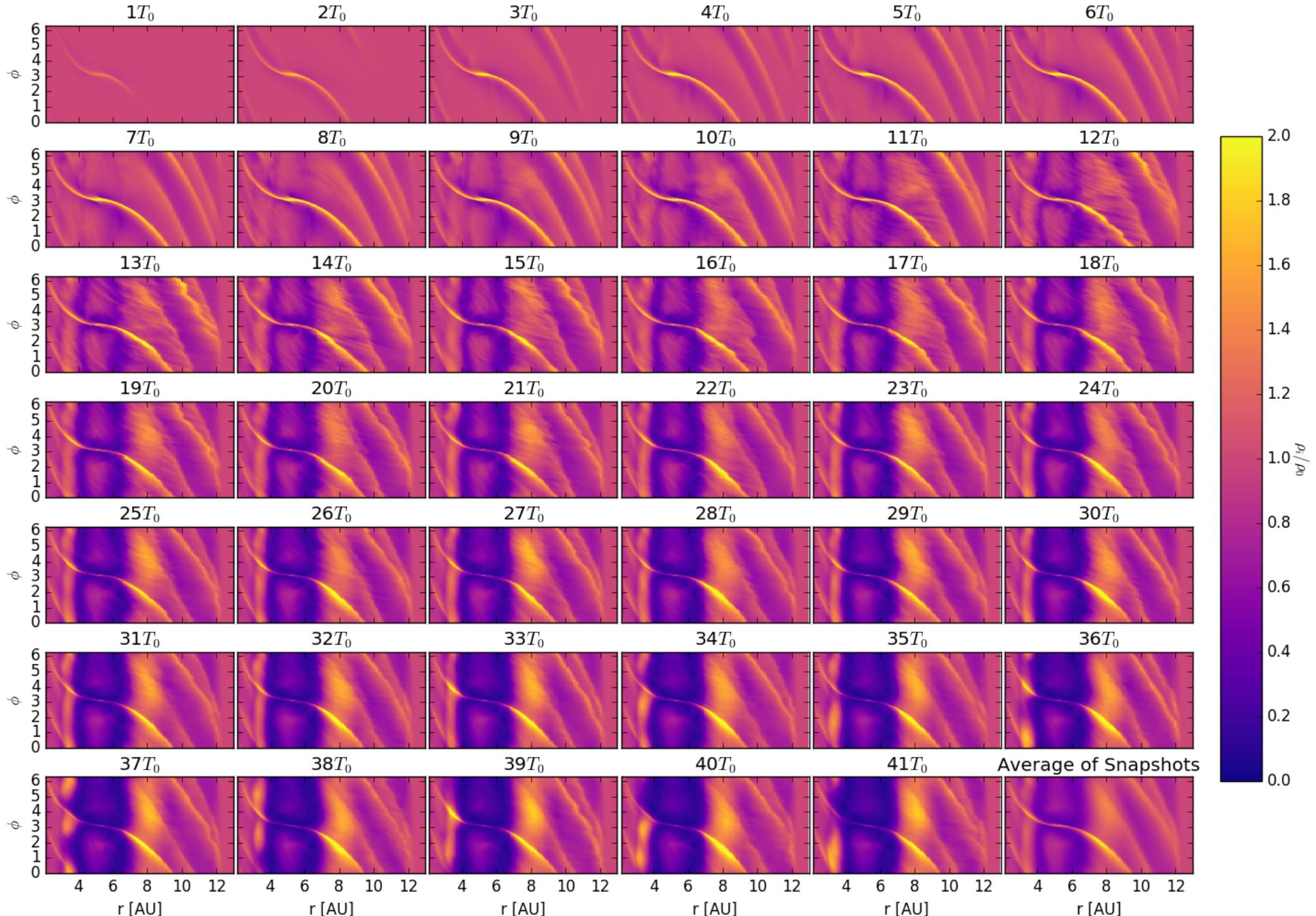
Scattered Light



## Primary and Secondary spiral arms

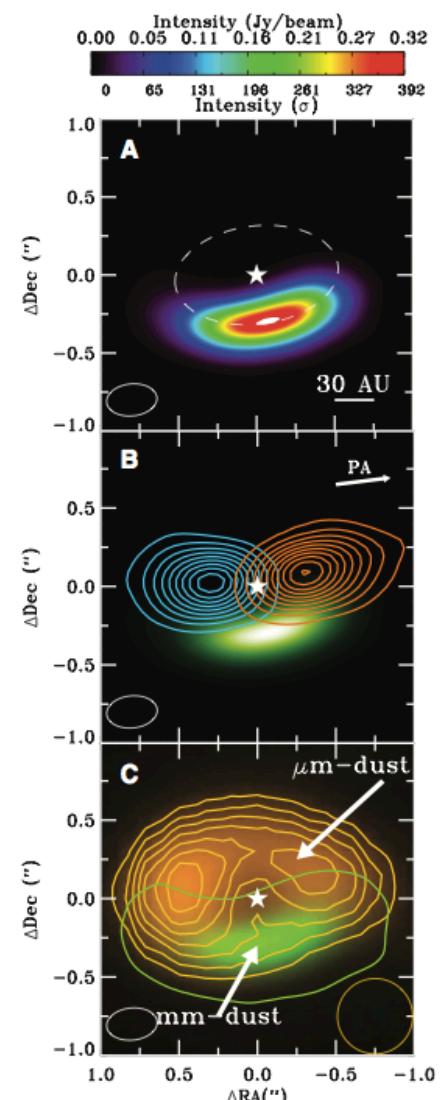


# The raised feature has its origins in a secondary spiral arm



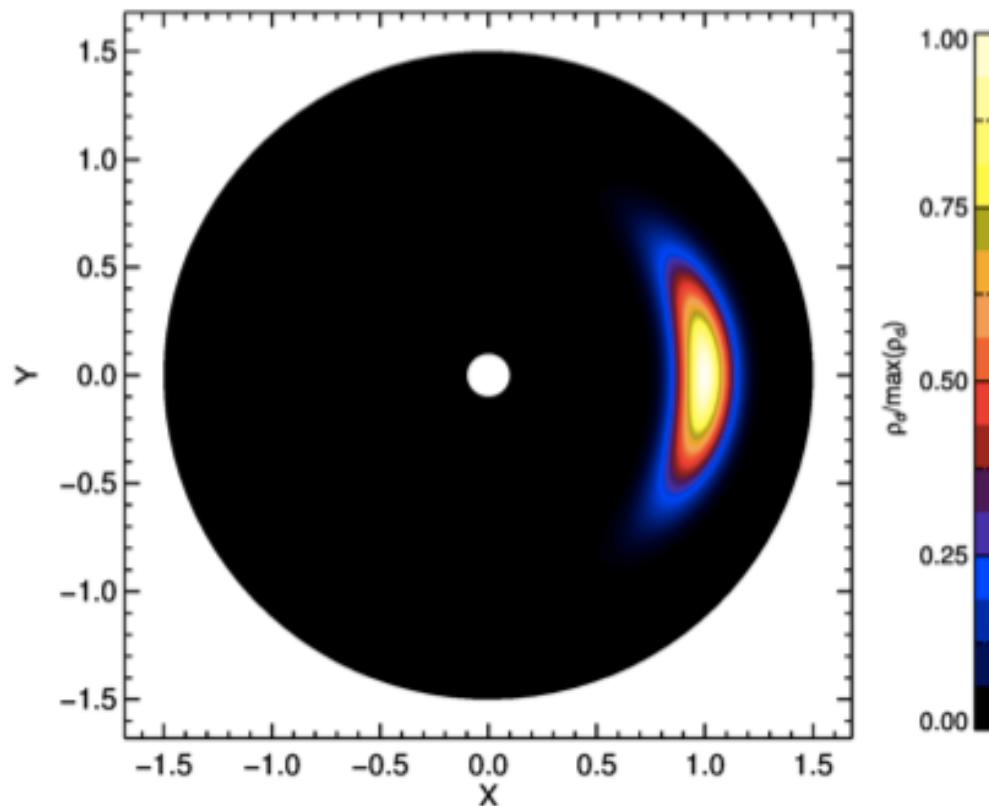
# Conclusions

- Disk vortices are a prime location for planet formation
  - Tea leaf effect
- Dust trapped in drag-diffusion equilibrium exp observations
- Hydrodynamical instabilities vs planet excitation
  - Vertical Shear Instability
    - *Vertical violation of Solberg-Hoiland criterion*
  - Convective Overstability
    - *Amplification of epicyclic motion by buoyancy*
  - Zombie Vortex Instability
    - *Resonance between epicyclic and buoyancy*
- Hot lobes next to high mass planets at high res
- Planets puff up their outer gap edges – visible



# Conclusions

- Disk vortices are a prime location for planet formation
  - Tea leaf effect
- Dust trapped in drag-diffusion equilibrium explains the observations



citation mechanism:

*criterion*

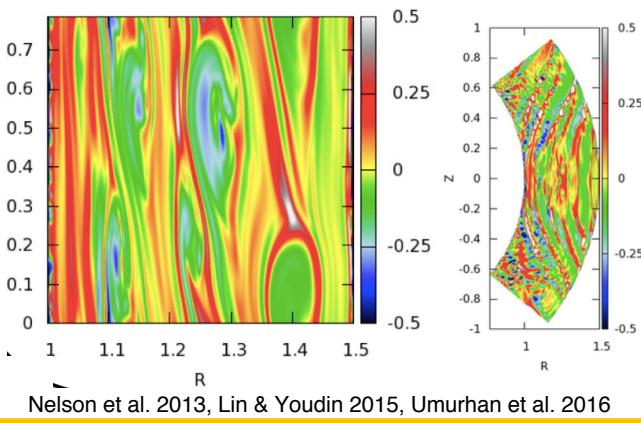
*buoyancy*

*buoyancy frequency*

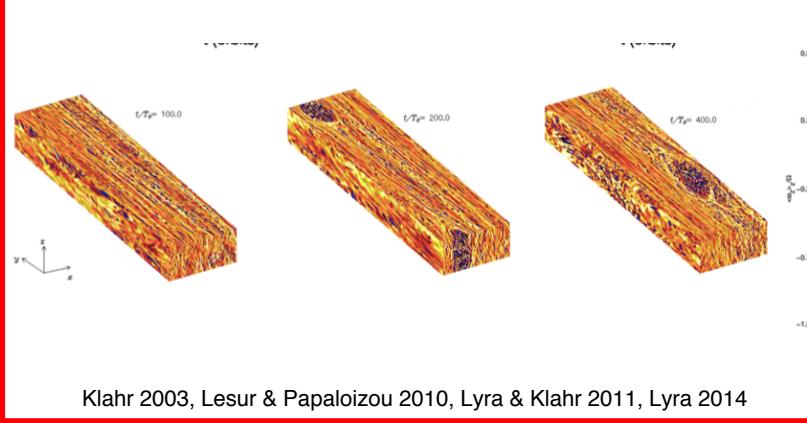
*high resolution*

sible in scattered light

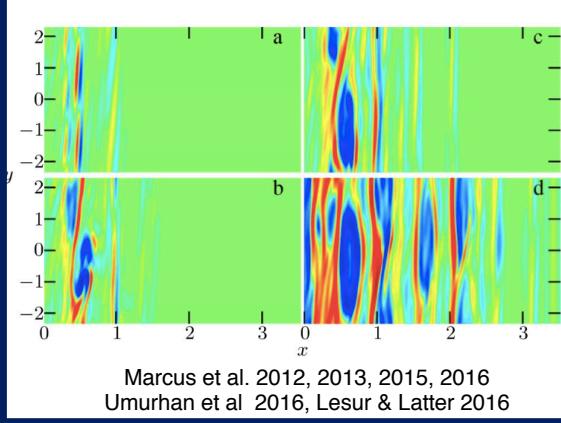
## Vertical Shear Instability



## Convective Overstability



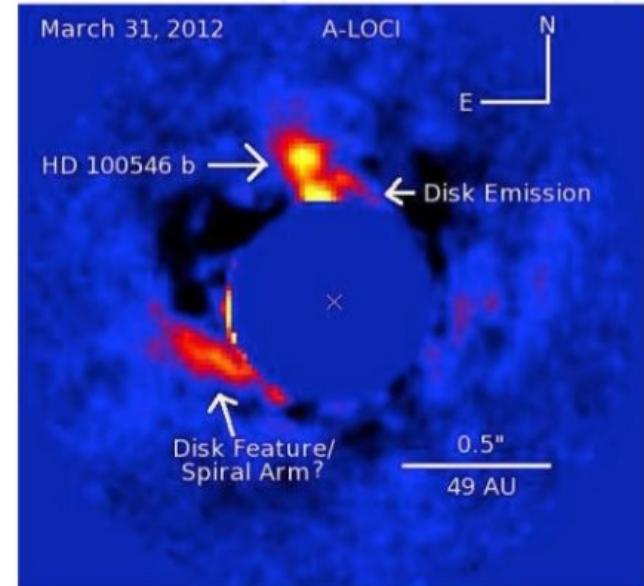
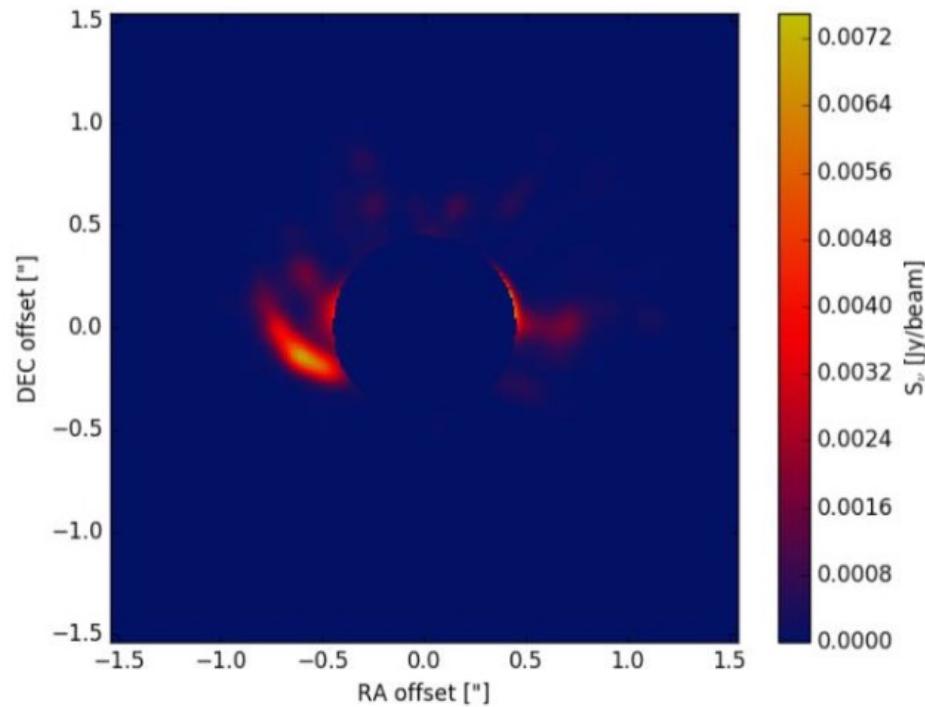
## Zombie Vortex Instability



- Hydrodynamical instabilities vs planet excitation mechanism:
  - Vertical Shear Instability
    - *Vertical violation of Solberg-Hoiland criterion*
  - Convective Overstability
    - *Amplification of epicyclic motion by buoyancy*
  - Zombie Vortex Instability
    - *Resonance between epicyclic and buoyancy frequency*
- Hot lobes next to high mass planets at high resolution
- Planets puff up their outer gap edges – visible in scattered light

# Conclusions

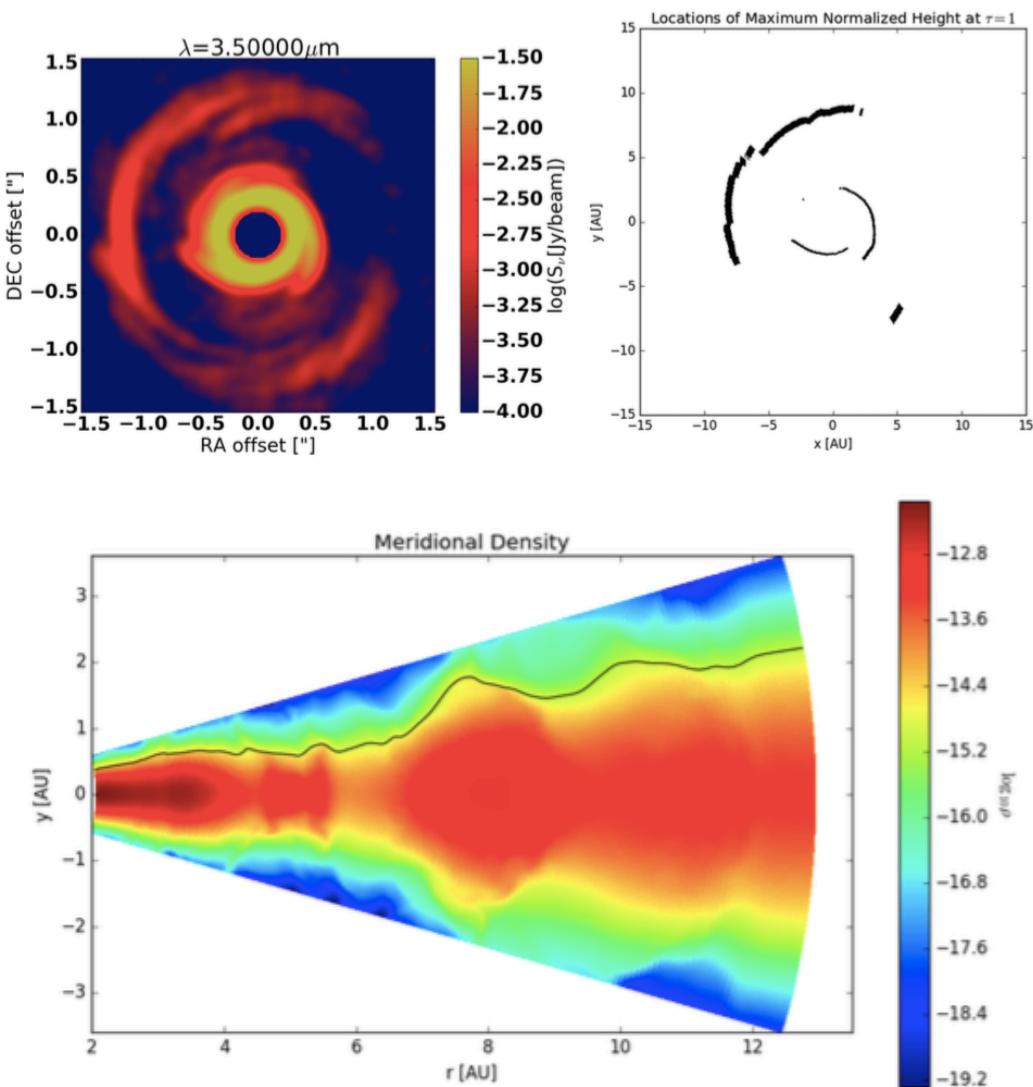
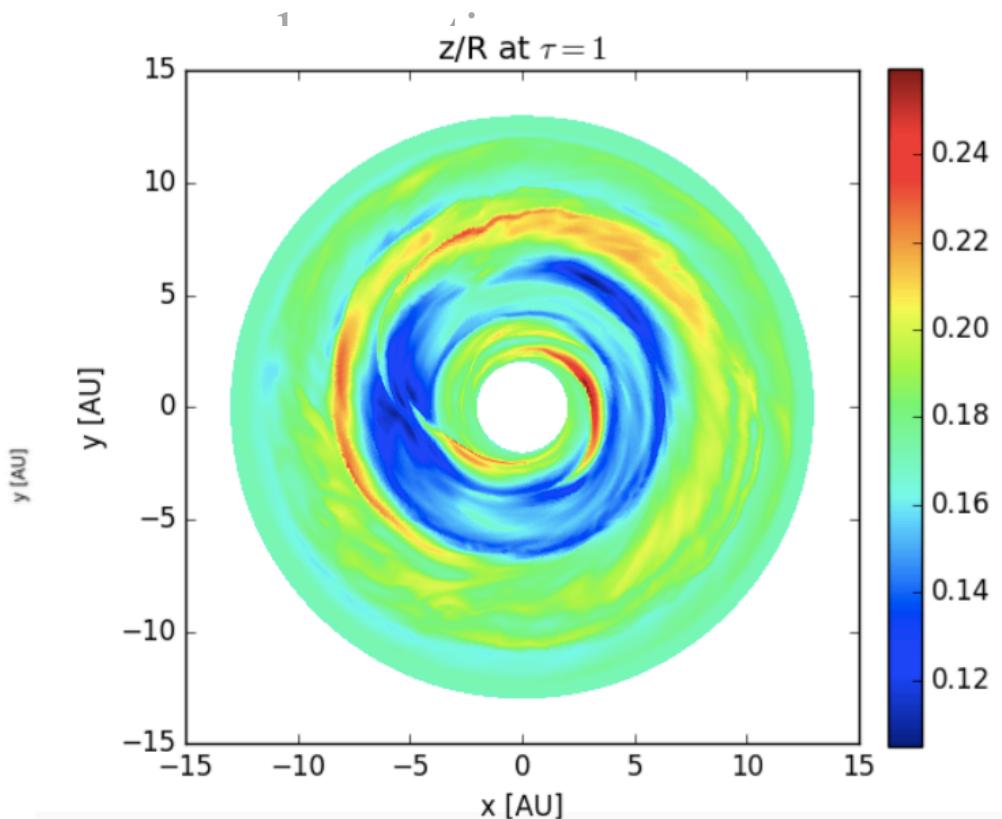
- Disk vortices are a prime location for planet formation
  - Tea leaf effect



- Hot lobes next to high mass planets at high resolution
- Planets puff up their outer gap edges – visible in scattered light

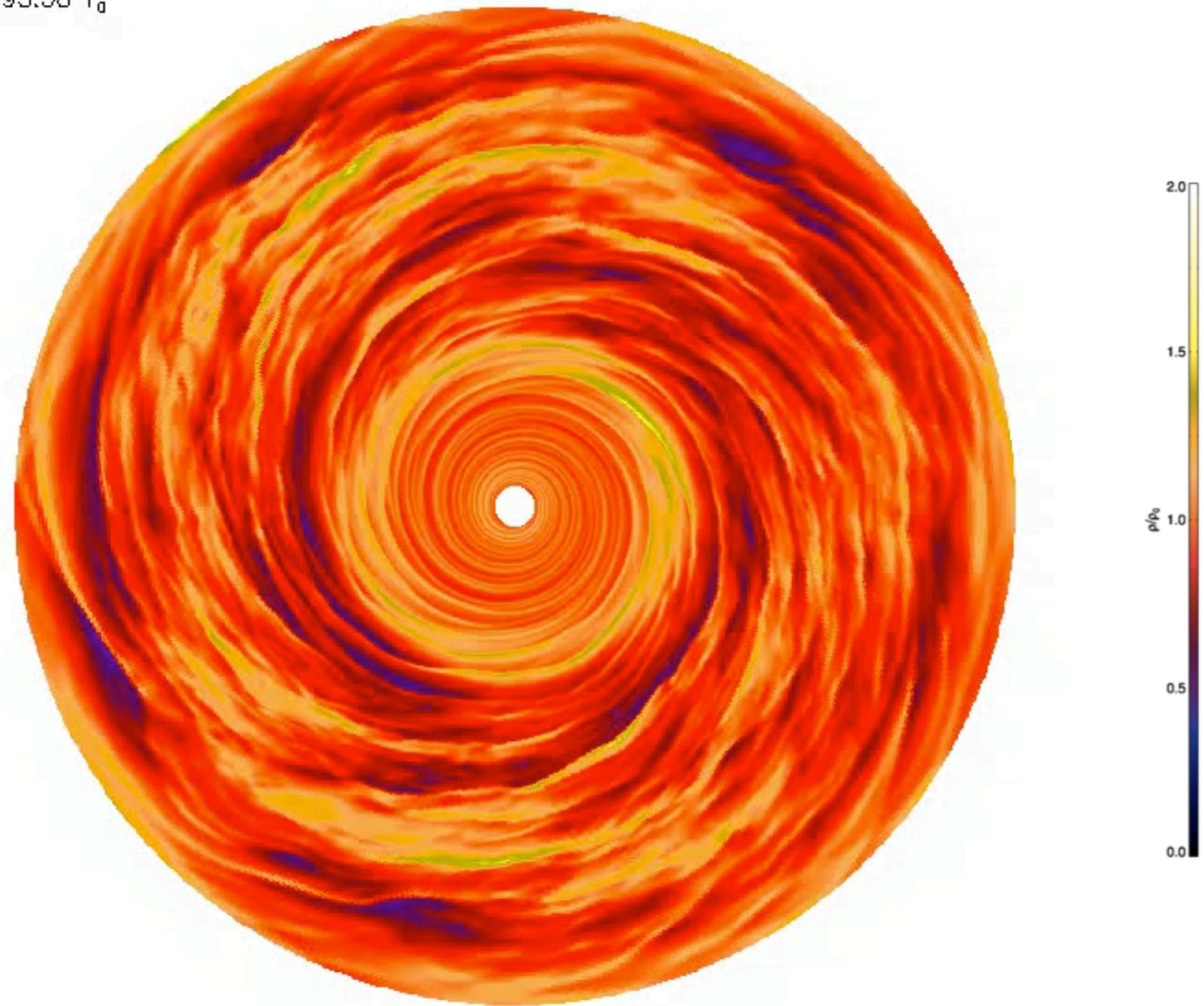
# Conclusions

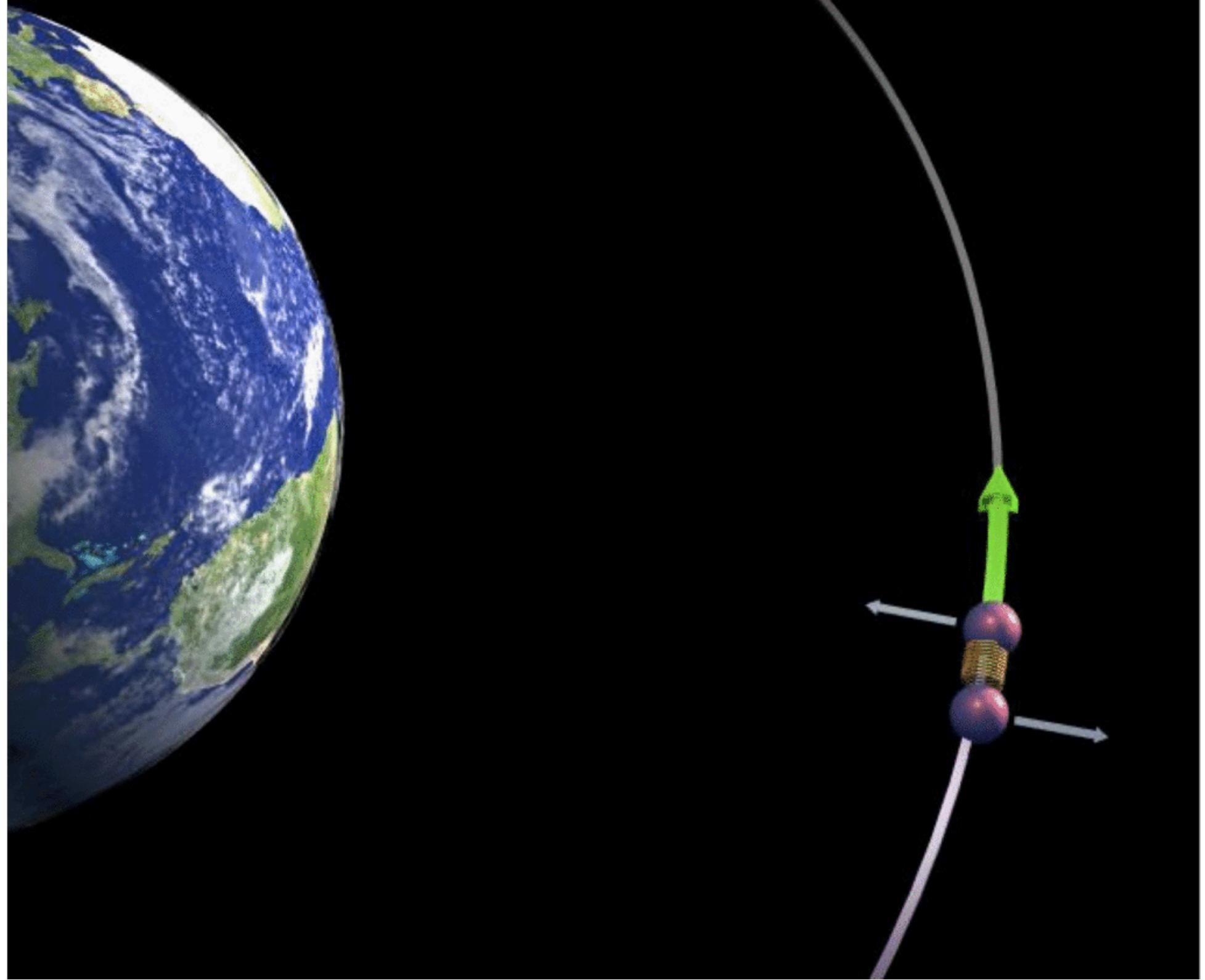
- Disk vortices are a prime location for:
  - Tea leaf effect
- Dust trapped in drag-diffusion



- Planets puff up their outer gap edges – visible in scattered light

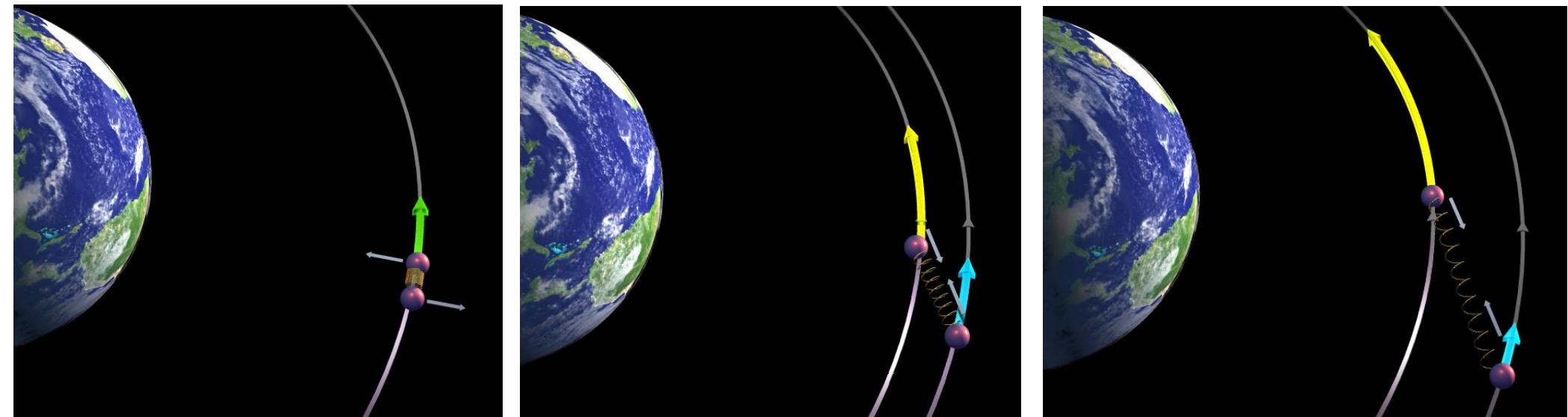
$t=95.58 T_0$





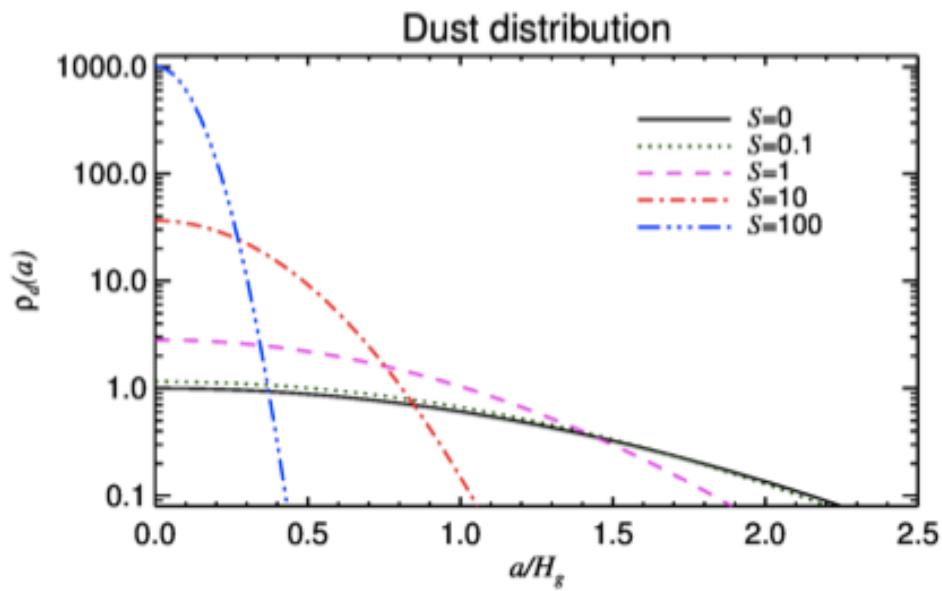
**Stretching builds up tension**

**Tension resists shear**



**Beads exchange angular momentum**

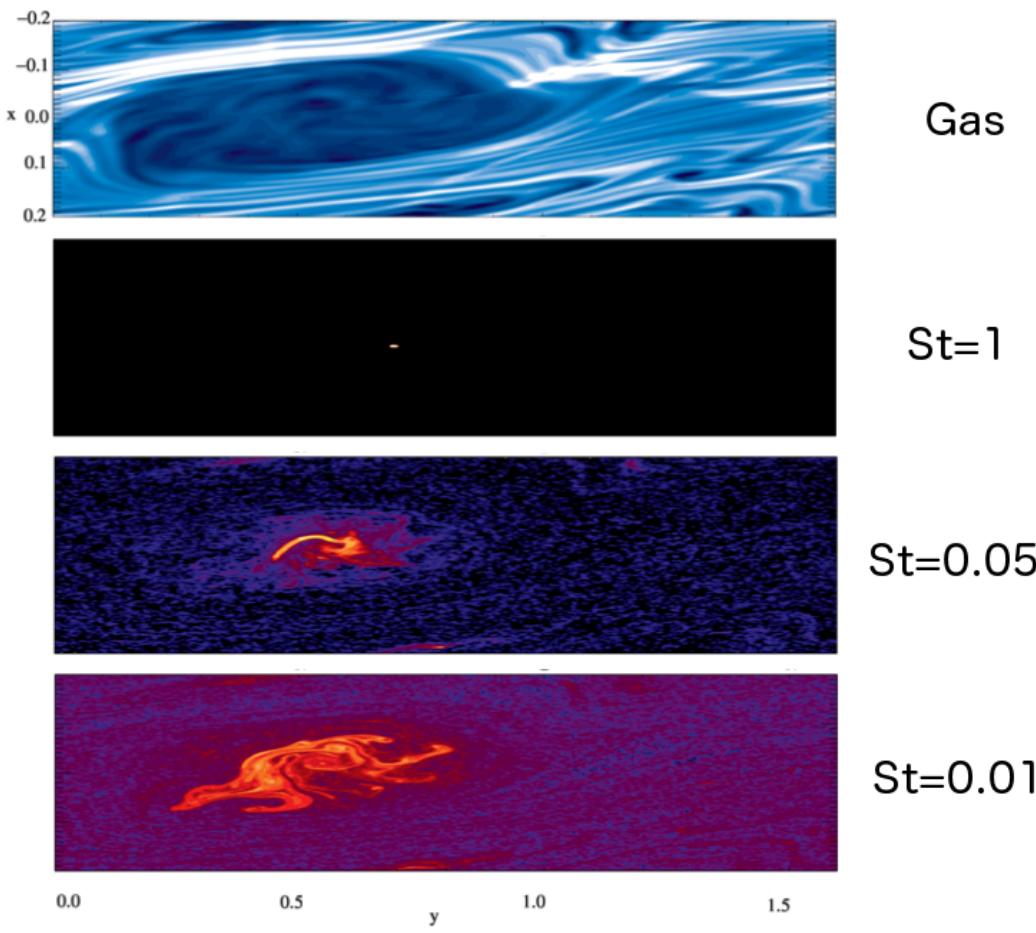
# Analytical vs Numerical



$$S = \text{St}/\delta$$

$$\delta = v_{\text{rms}}^2 / c_s^2,$$

Lyra & Lin (2013)



Raettig et al (2015)

# SPHERE-ALMA-VLA overlay of MWC 758

SPHERE ( $\mu\text{m}$ )  
ALMA ( $\sim \text{mm}$ )  
VLA (cm-m)

