







**S**T**A**T**E** 











## **Evolution of Circumstellar Disks and Planet Formation**

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**New Mexico State University** 

**Current Funding** 



AAG - 2020, 2021



XRP - 2023 EW – 2021, 2022, 2023 TCAN - 2020

**Computational Facilities** 



NRAO - Apr 11<sup>th</sup>, 2024

## Outline

- Planet Formation
- Hydrodynamical Instabilities
- Disk observations

## TCAN-2020 (Theoretical and Computational Astrophysics Network) Planet Formation Collaboration



#### ΡΙ

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# Planet Formation





## **Circumstellar/Protoplanetary Disks**





**Planet Formation** 

"Planets form in disks of gas and dust"



## A miracle happens —



### **Dust evolution**



## **Headwind and Dust Drift**



The gas has some pressure support (sub-Keplerian).

The pebbles do not feel gas pressure (Keplerian).

**Dust coagulation and drift** 

Dust particle coagulation and radial drift

F.Brauer, C.P. Dullemond Th. Henning

Brauer et al. (2008)

## **Streaming Instability**

The dust drift is hydrodynamically unstable



Youdin & Goodman '05, Johansen & Youdin '07, Youdin & Johansen+ '07, Kowalik+ '13, Lyra & Kuchner '13, Schreiber+ '18, Klahr & Schreiber '20, Simon+ '16, '17, Carrera+ '15, '17, '20, Gole+ '20, Li+ '18, '19, Abod+ '19, Nesvorny+ '19

#### Gravitational collapse into planetesimals



Johansen et al. (2007)

**Turbulence and Accretion in 3D Global MHD Simulations of Stratified Protoplanetary Disk** 

#### Turbulence





#### Turbulence concentrates solids mechanically in pressure maxima

Lyra et al. (2008a)

## **Dead zones**





Lyra et al. (2008b, 2009a); After Lovelace & Hohlfeld 1978, Toomre 1981, Papaloizou & Pringle (1984), Hawley (1987), Lovelace et al (1999), Li et al. (2000), Varniere & Tagger (2005).







Magnetized inner disk + resistive outer disk Lyra & MacLow (2012)

#### Rossby wave instability (Kelvin-Helmholtz Instability in rotating disks)











## Vortices – an ubiquitous fluid mechanics phenomenon



#### Von Kármán vortex street





## **Vortex Trapping**



Grains do not feel the pressure gradient. They sink towards the center, where they accumulate.

#### Aid to planet formation

(Barge & Sommeria 1995, Tanga et al. 1996, Adams et al. 1996)

#### Speeds up planet formation enormously

(Lyra et al. 2008b, 2009ab, Raettig et al. 2012, 2021)

## **Vortex Trapping**



Time

## **Vortex Trapping – Initial Mass Function**



## **Initial Mass Function - Convergence**



Lyra et al. (submitted)

## Take home message

• Two routes for planet formation



**Streaming Instability** 



#### Vortex Trapping



Lyra+08,18 Raettig+Lyra 12,15,21

- Planet formation and turbulence.
  - Does turbulence help (concentration at large scales) or hinder (diffusion at small scales)?

# **Disk Instabilities**

## **Dead zones**



There should be a magnetized, active zone and a non-magnetic, dead zone







Magnetized inner disk + resistive outer disk Lyra & MacLow (2012)

#### **Instability Map**



Lesur et al. (2022)

#### Vertical shear instability

Angular velocity not constant in cylinders: unstable

#### Buoyancy stabilizes. The most unstable mode is isothermal.



#### Vertical shear instability



#### **Convective Overstability**



Lesur & Papaloizou (2010) Lyra & Klahr (2011) Klahr & Hubbard (2014) Lyra (2014) Latter (2016) Volponi (2016) Reed & Latter (2021) Raettig et al. (2021)

#### Resonant Buoyant Instability (née Zombie Vortex Instability)

 $\infty_{z}$  at x-y plane z=0.40431 t=0 2 2 0 - | -2 2 3 0 х

Cascade of baroclinic critical layers

## **Hydrodynamical Instabilities**



$\Omega \tau << 1$	$\Omega \tau \sim 1$	$\Omega \tau >> 1$
$(\kappa < 1 \text{ cm}^2/\text{g})$	$(\kappa \sim 1-50 \text{ cm}^2/\text{g})$	$(\kappa > 50 \text{ cm}^2/\text{g})$

#### Take-home message



# **Disk Observations**




# **Disk lifetime**





(Ribas et al. 2014)

Disks dissipate within ~10 Myr Mass accretion rates ~  $10^{-8} M_{\odot} \text{ yr}^{-1}$ 

# **Disk spectra**



# A class of disks with missing hot dust.



# Disks with missing hot dust.





☆



al line



# **Planetary companion**



#### These cavities may be the telltale signature of forming planets

PDS 70 and PDS 70b



A way to directly study planet-disk interaction

# Planet-disk interaction: gaps, spirals, and vortices.



#### **Observational evidence: gaps, spirals, and vortices**



The ALMA Partnership et al. (2015)

Muto et al. (2012)

van der Marel et al. (2013)

# ALMA

# VLA



# The ALMA Revolution



### **Before ALMA**

ALMA





#### Dust traps in disks: ALMA Cycle 0 (2012)





#### Oph IRS 48



#### A Major Asymmetric Dust Trap in a Transition Disk

Nienke van der Marel,<sup>1</sup>\* Ewine F. van Dishoeck,<sup>1,2</sup> Simon Bruderer,<sup>2</sup> Til Birnstiel,<sup>3</sup> Paola Pinilla,<sup>4</sup> Cornelis P. Dullemond,<sup>4</sup> Tim A. van Kempen,<sup>1,5</sup> Markus Schmalzl,<sup>1</sup> Joanna M. Brown,<sup>3</sup> Gregory J. Herczeg,<sup>6</sup> Geoffrey S. Mathews,<sup>1</sup> Vincent Geers<sup>7</sup>

The statistics of discovered exoplanets suggest that planets form efficiently. However, there are fundamental unsolved problems, such as excessive inward drift of particles in protoplanetary disks during planet formation. Recent theories invoke dust traps to overcome this problem. We report the detection of a dust trap in the disk around the star Oph IRS 48 using observations from the Atacama Large Millimeter/submillimeter Array (ALMA). The 0.44-millimeter/wavelength continuum map shows high-contrast crescent-shaped emission on one side of the star, originating from millimeter-sized grains, whereas both the mid-infrared image (micrometer-sized dust) and the gas traced by the carbon monoxide 6-5 rotational line suggest rings centered on the star. The difference in distribution of big grains versus small grains/gas can be modeled with a vortex-shaped dust trap triggered by a companion.

A lthough the ubiquity of planets is confirmed almost daily by detections of new exoplanets (1), the exact formalong-standing problem in astrophysics (2). In

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#### van der Marel+ '13

#### A huge vortex observed with ALMA



# **Drag-Diffusion Equilibrium**



Trapped particle

#### **Analytical Solution for dust in Drag-Diffusion Equilibrium**



Steady-state solution  

$$\rho_d(a,z) = \epsilon \rho_0 (S+1)^{3/2} \exp \left\{ -\frac{\left[a^2 f^2(\chi) + z^2\right]}{2H^2} (S+1) \right\}$$
Lyra & Lin '13

$$S = \frac{St}{d}$$
$$\delta = v_{\rm rms}^2 / c_s^2,$$

- *a* = vortex semi-minor axis
- H = disk scale height (temperature)
- $\chi$  = vortex aspect ratio
- $\delta$  = diffusion parameter
- St = Stokes number (particle size)
- $f(\chi)$  = model-dependent scale function

#### **Disk Tomography** SPHERE-ALMA-VLA overlay of MWC 758



Marino+Lyra '15

#### Pebble trapping







Casassus+Lyra '19

### **Model vs Observation**



Raettig+Lyra '15

#### Take home message

• Vortex-trapped dust in drag-diffusion equilibrium explains the observations

$$\rho_d(a,z) = \varepsilon \rho_0 (S+1)^{3/2} \exp\left\{-\frac{\left[a^2 f^2(\chi) + z^2\right]}{2H^2} (S+1)\right\}$$







# The future

After 10 years of ALMA...

Nearly all nearby disks observed at <0.1" (< 20-30AU) show substructures.

3 main types of substructures

- Crescent-shaped
- Spiral arms
- Rings/Gaps



# ngVLArocks!





![](_page_58_Picture_3.jpeg)

# **Planets at 5AU**

ALMA @ 0.87mm

ngVLA @ 3mm

5 mas = 0.7 AU

![](_page_59_Figure_1.jpeg)

Ricci et al. 2018

ngVLA identifies gaps/substructures down to ~5-10 M<sub>Earth</sub>

# ngVLA: Proper motions

Jupiter at 5 AU

![](_page_60_Picture_2.jpeg)

#### Conclusions

- Two routes for planet formation (streaming instability and vortices, complementary)
- Does turbulence help (concentration at large scales) or hinder (diffusion at small scales)?
- Three dynamical instabilities in the Ohmic dead zone
  - Different regimes of opacity, operate in different regions
  - Saturate into vortices
  - Dust trapped in drag-diffusion equilibrium explains the observations
- Issues:
  - Are the dynamical instabilities (chiefly the Vertical Shear Instability) responsible for the observed crescents?
  - Overlap between instabilities unclear
  - Global model of Convective Overstability needed
  - Relevance of Resonant Buoyant Instability ("zombie vortex") unclear/unlikely.
  - Planet formation properties / Synergy with streaming instability

	ZVI	COV	VSI
Global model	$\bigotimes$	$\bigotimes$	
Vertical Stratification		$\bigotimes$	$\bigcirc$
Boundaries with other instabilities	$\bigotimes$	$\bigotimes$	$\bigotimes$
Interaction with dust	$\bigotimes$	$\bigcirc$	$\bigcirc$
Observational Validation/Rule out	$\bigotimes$	$\bigotimes$	$\bigotimes$
Planet Forming Properties	$\bigotimes$	$\bigotimes$	$\bigotimes$

![](_page_61_Picture_14.jpeg)

#### Convergence

![](_page_62_Figure_1.jpeg)

![](_page_63_Figure_0.jpeg)

Fig. 7.— Left: Upper limits to turbulent velocities in HD 163296 as a function of radius R and mid-plane height z. The colors denote the species from which there is a majority of emission; the CO transitions are at large z whereas C<sup>18</sup>O and DCO+ are from lower in the disk. Also included on the plot are lines of constant  $\Sigma$ , H and 3H (H being the gas scale height) and where CO is ice. In this source, turbulence is at most a few percent of the sound speed. From Flaherty et al. (2017). Right: The turbulent velocity as a function of radius as measured from CO emission in DM Tau (hashed lines), compared to the results from previous work using CS (Guilloteau et al. 2012; grey shaded region). As opposed to HD 163296, DM Tau exhibits strong turbulence, consistent with theoretical predictions Flaherty et al. (2020).

#### **Vortices and Planet Formation**

![](_page_64_Figure_1.jpeg)

## Pebble trapping in vortices in LOCAL models

![](_page_65_Figure_1.jpeg)

Raettig et al (2015)

#### Vortex destruction at high dust load?

![](_page_66_Figure_1.jpeg)

Dust to gas ratio 10<sup>-4</sup>

Dust to gas ratio 10<sup>-2</sup>

![](_page_66_Figure_4.jpeg)

67

Raettig et al (2015)

# Pebble trapping does not destroy vortices

![](_page_67_Figure_1.jpeg)

# Vortex column disrupted only around the midplane

![](_page_68_Figure_1.jpeg)

# Pebble drift: follows vortex

![](_page_69_Figure_1.jpeg)

# **Pebble trapping in 3D vortices**

![](_page_70_Figure_1.jpeg)

![](_page_70_Figure_2.jpeg)

Raettig et al. (2021)

St=0.05

 $\begin{array}{c} \epsilon_{0}=10^{-2} \\ \epsilon_{0}=10^{-3} \\ \epsilon_{0}=10^{-4} \end{array}$ 

 $\begin{array}{c} \text{St=1.0} \\ \epsilon_0 = 10^{-2} \\ \epsilon_0 = 10^{-3} \end{array}$ 

 $\epsilon_0 = 10^{-4}$ 

30

40

20

# nature astronomy

Fingerprints of streaming instability

How can we verify the streaming instability hypothesis?
### **Planetesimal Formation**



Initial mass function consistent with mass distribution of asteroid belt. Slope 1.6





### **Structure of the Kuiper Belt**



Gladman+ '08, Lacerda '09, Batygin+ '10, Dawson & Murray-Clay '12

### **Structure of the Kuiper Belt**



# Arrokoth (MU<sub>69</sub>)



New Horizons Flyby, Jan 2019

# **Cold Classical KBOs: Preference for Prograde**



# **Counting binaries: Preference for Prograde (~80%)**



## The density dichotomy of Kuiper Belt objects



# **Possible Solution?**

- Assumptions
  - Constant composition at birth and growth
  - Porosity removal by gravitational compaction



Bierson & Nimmo 2019

- Problems
  - Low-mass objects need to be unreasonably porous
  - Timing! <sup>26</sup>Al would melt if formed within 4 Myr

# **Abandoning Constant Composition**

Heating and UV irradiation remove ice on Myr timescales

- Small grains lofted in the atmosphere lose ice
- Big grains are shielded and remain icy.





# Split into icy and silicate pebbles



# The first planetesimals are icy



#### **Pebble Accretion**



Lyra+ '08, '09, '23, Ormel & Klahr '10, Lambrechts & Johansen '12 See Johansen & Lambrechts '17 for a review

### Pebble Accretion: Geometric, Bondi, and Hill regime

Bondi accretion - Bound against thermal (dynamic) kinetic energy Hill accretion - Bound against stellar tide

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Bondi radius

Capture radius

$$\xi \equiv \left(\frac{R_{\rm acc}}{2H_d}\right)^2 \qquad \dot{M}_{\rm 3D} = \lim_{\xi \to 0} \dot{M} = \pi R_{\rm acc}^2 \rho_{d0} \delta v,$$
$$\dot{M}_{\rm 2D} = \lim_{\xi \to \infty} \dot{M} = 2R_{\rm acc} \Sigma_d \delta v,$$



Johansen & Lambrechts '17

### Integrate pebble accretion





#### Pebble Accretion: Pebbles of different size accrete differently





Drag time = Time to cross Bondi sphere

Drag time ~ Orbital Time

### **Accretion Rates**



$$\begin{split} \varepsilon &= \left(\frac{R_{\rm acc}}{2H_d}\right)^2 \quad \begin{array}{l} \dot{M}_{\rm 3D} = \lim_{\substack{\xi \to 0}} \dot{M} = \pi R_{\rm acc}^2 \rho_{d0} \delta v, \\ \dot{M}_{\rm 2D} &= \lim_{\substack{\xi \to \infty}} \dot{M} = 2R_{\rm acc} \Sigma_d \delta v, \\ \dot{M}_{\rm 2D} = \lim_{\substack{\xi \to \infty}} \dot{M} = 2R_{\rm acc} \Sigma_d \delta v, \\ &\text{Lambrechts & Johansen (2012)} \end{split} \\ \hline \\ \begin{array}{l} \dot{P} Olydisperse \ (multiple \ species) \\ \dot{M}_{\rm 2D,Hill} = \frac{6(1-p)}{14-5q-3k} \left(\frac{St_{\rm max}}{0.1}\right)^{2/3} \Omega R_H^2 Z \Sigma_g. \\ \dot{M}_{\rm 3D,Bondi} \approx C_1 \frac{\gamma_l \left(\frac{b_l+1}{s}, j_l a_{\rm max}^s\right)}{s j_1^{(b_l+1)/s}} + C_2 \frac{\gamma_l \left(\frac{b_2+1}{s}, j_2 a_{\rm max}^s\right)}{s j_2^{(b_2+1)/s}} + \\ &C_3 \frac{\gamma_l \left(\frac{b_3+1}{s}, j_3 a_{\rm max}^s\right)}{s j_3^{(b_3+1)/s}} + C_4 \frac{\gamma_l \left(\frac{b_4+1}{s}, j_4 a_{\rm max}^s\right)}{s j_4^{(b_4+1)/s}}, \end{split}$$

Lyra et al. (2023)



Lyra et al. 2023

### Growing Pluto by silicate pebble accretion





### Growing Pluto by silicate pebble accretion



#### **Resulting Densities vs Mass relations**





#### Take-home message

- Streaming Instability fits
  - slope of asteroid belt distribution,
  - prograde-retrograde distribution of Kuiper belt objects
  - Low density of small classical Kuiper belt objects
- Pebble accretion is a very efficient planetary growth mechanism
  - Polydisperse Bondi accretion 1-2 orders of magnitude more efficient than monodisperse
  - Silicate pebble accretion explains densities of high-mass Kuiper belt objects

- Polydisperse Bondi accretion 1-2 orders of magnitude more efficient than monodisperse
  - Best accreted pebbles are those of drag time ~ Bondi time, not the largest ones
  - The largest ones dominate the mass budget, but accrete poorly
- Onset of Bondi accretion 1-2 orders of magnitude lower in mass compared to monodisperse
  - Reaches 100-350km objects within Myr timescales
  - Bondi accretion possible on top of Streaming Instability planetary embryos within disk lifetime
- Analytical solution to
  - Polydisperse 2D Hill and 3D Bondi



Simon et al. 2023