

# Hydrodynamical Instabilities in the Ohmic dead zone in circumstellar disks and dwarf novae

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California State University Northridge

## *Collaborators*

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Hubert Klahr (MPIA)

Colin McNally (QMUL)

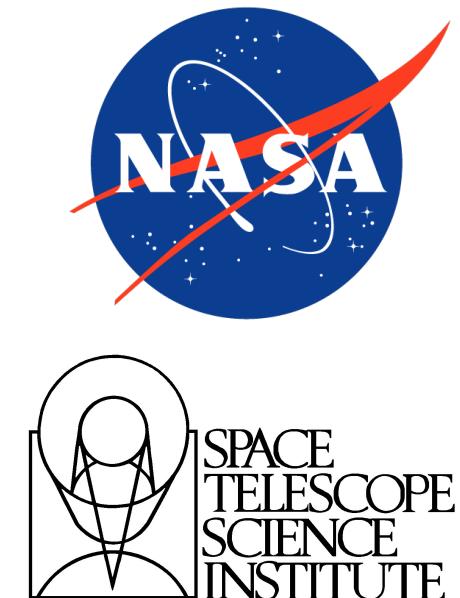
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Nikolai Piskunov (Uppsala)

Natalie Raettig (MPIA)

Orkan Umurhan (NASA Ames)

Neal Turner (JPL)



KITP, Feb 28<sup>th</sup>, 2017



# Outline

- The MRI active and dead zones
  - Rossby wave instability
- Hydrodynamical instabilities
  - Convective overstability
  - Zombie vortex instability
  - Vertical shear instability
- Secondary mechanisms
  - Elliptic and Magneto-Elliptic Instability
- Applications to dwarf novae

# Protoplanetary Disks



## PP disk fact sheet

Density:  $10^{13} - 10^{15} \text{ cm}^{-3}$   
(Air  $\sim 10^{21} \text{ cm}^{-3}$ )

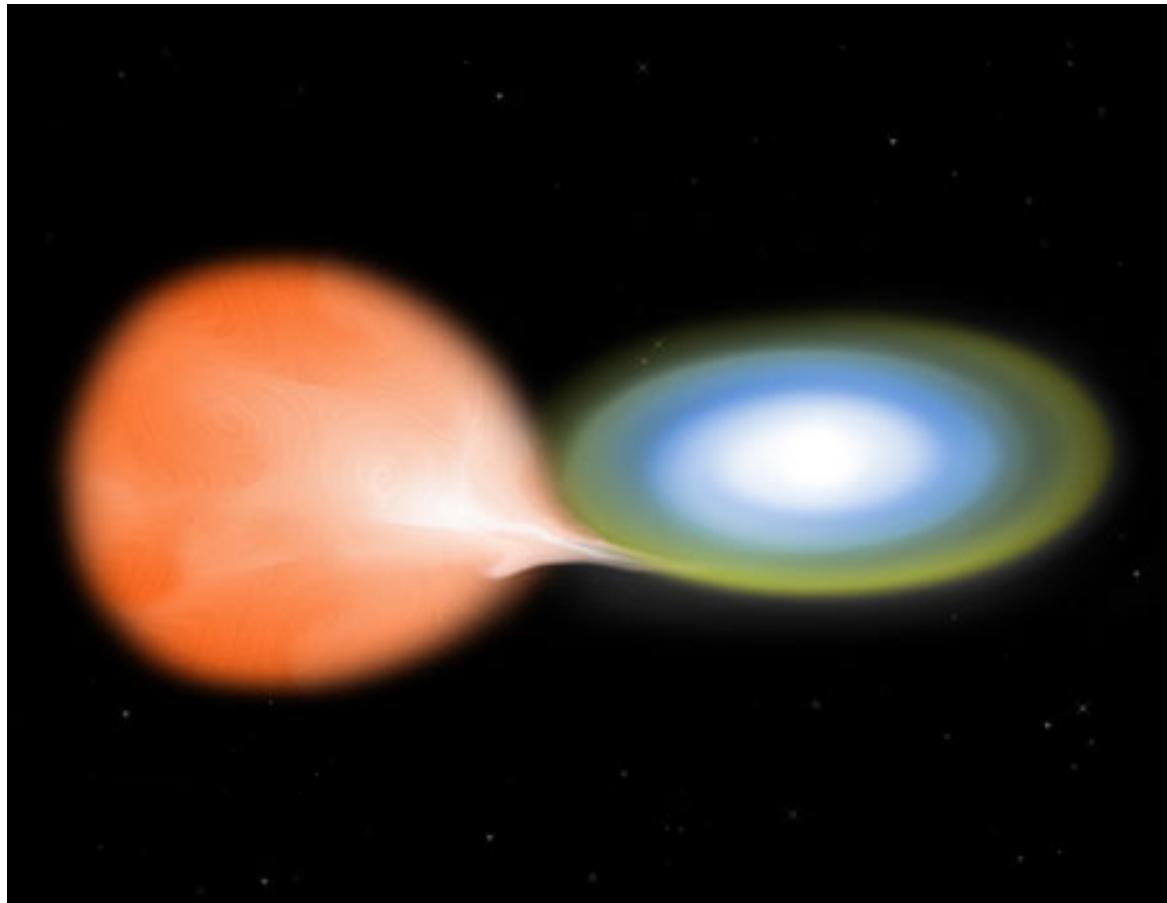
Temperature: 10-1000 K

Scale: 0.1-100AU  
(1 AU  $\sim 1.5 \times 10^{13} \text{ cm}$ )

Mass:  $10^{-3} - 10^{-1} M_{\text{sun}}$   
( $1 M_{\text{sun}} \sim 2 \times 10^{33} \text{ g}$ )

Timescales :  
Dynamical: 0.01 – 1000 yr  
Lifetime  $\sim 1\text{-}10 \text{ Myr}$

# Dwarf Novae Disks



## DN disk fact sheet

Density:  $10^{17} - 10^{19} \text{ cm}^{-3}$   
(Air  $\sim 10^{21} \text{ cm}^{-3}$ )

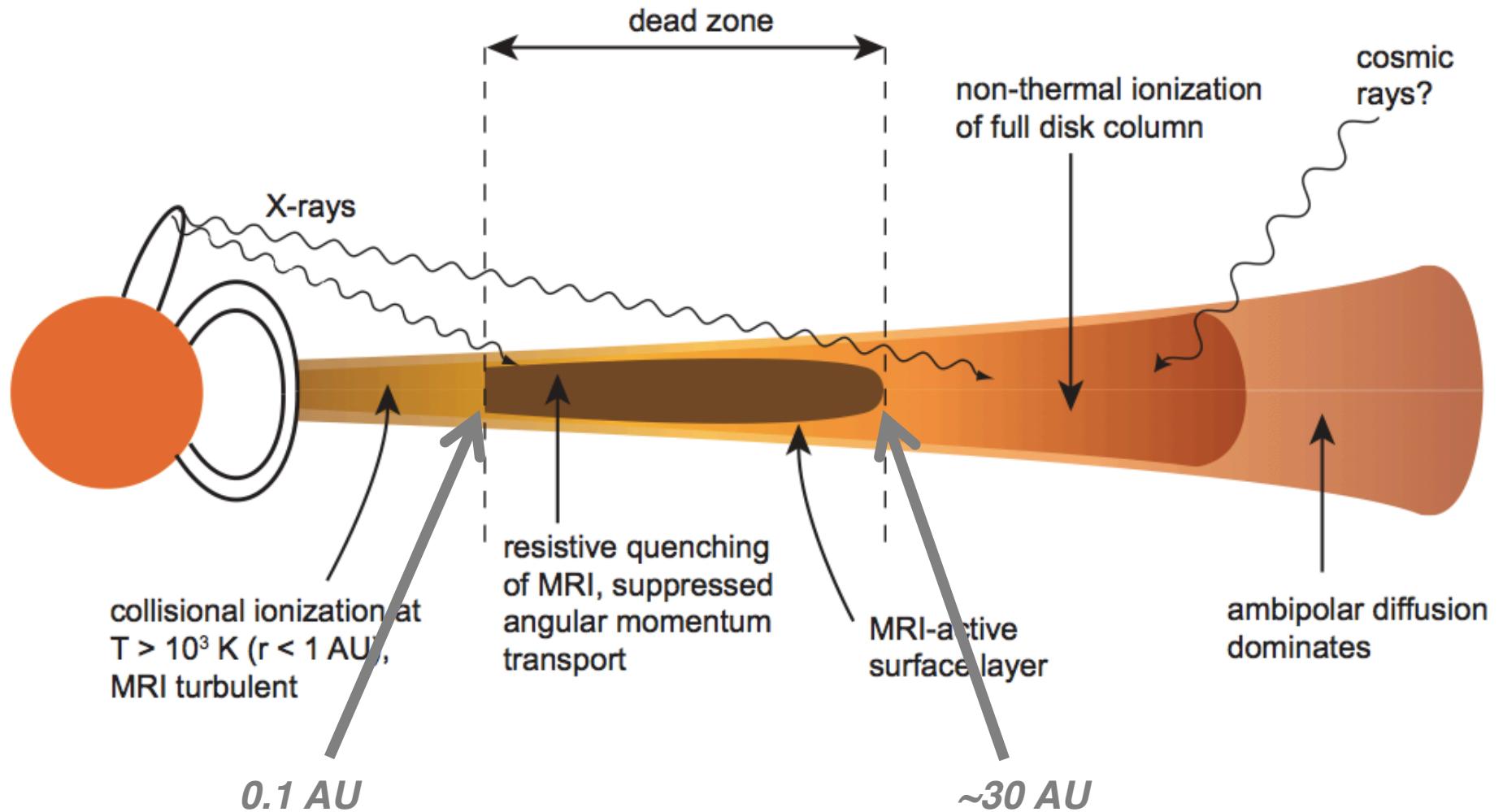
Temperature:  $10^3 - 10^5 \text{ K}$

Scale:  $10^9 - 10^{11} \text{ cm}$   
(1 AU  $\sim 1.5 \times 10^{13} \text{ cm}$ )

Mass:  $10^{-9} - 10^{-8} M_{\text{sun}}$   
( $1 M_{\text{sun}} \sim 2 \times 10^{33} \text{ g}$ )

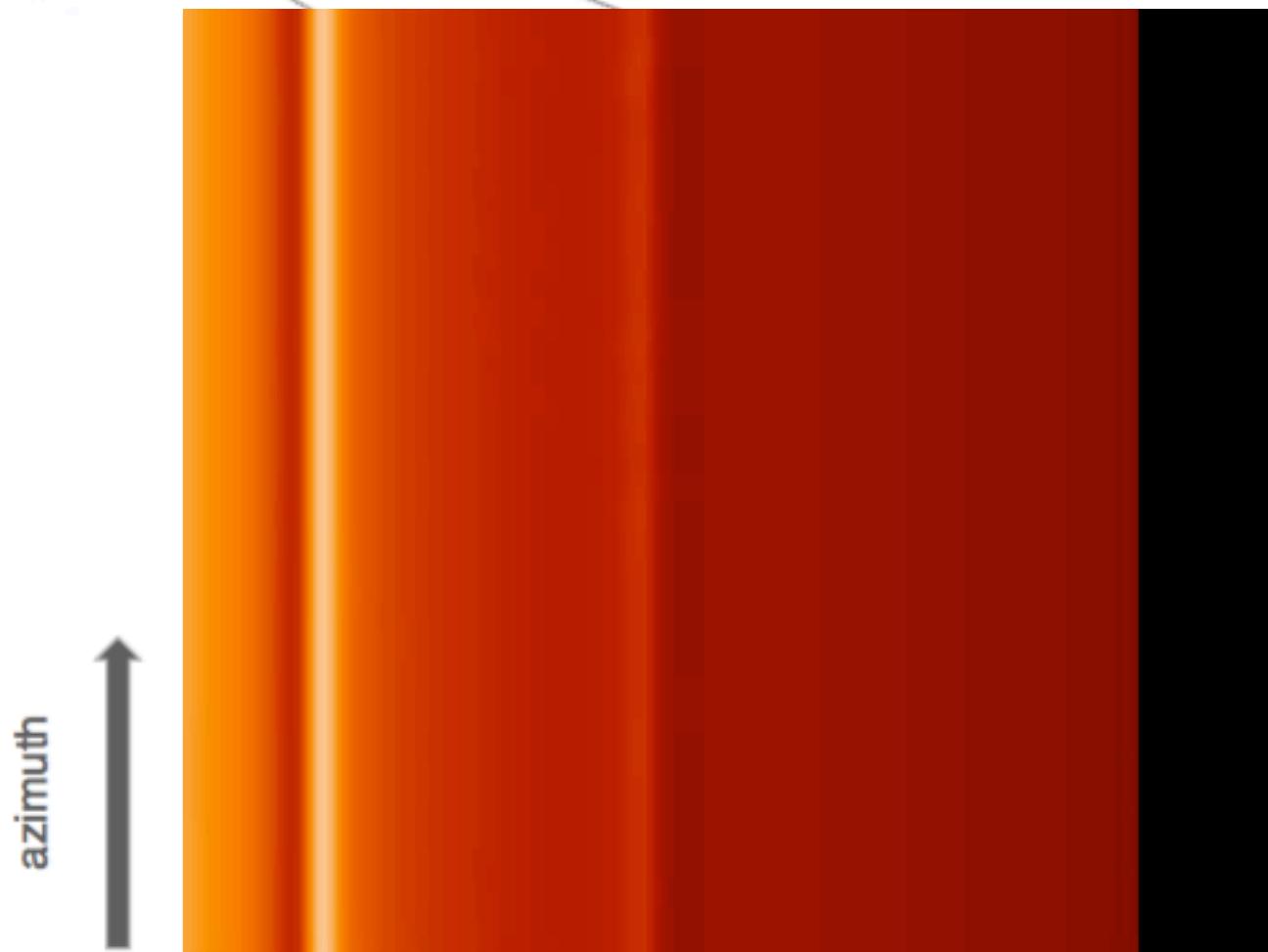
Timescales :  
Dynamical  $\sim$  hours  
Viscous time  $\sim$  weeks - months

# Dead zones





## A simple dead zone model

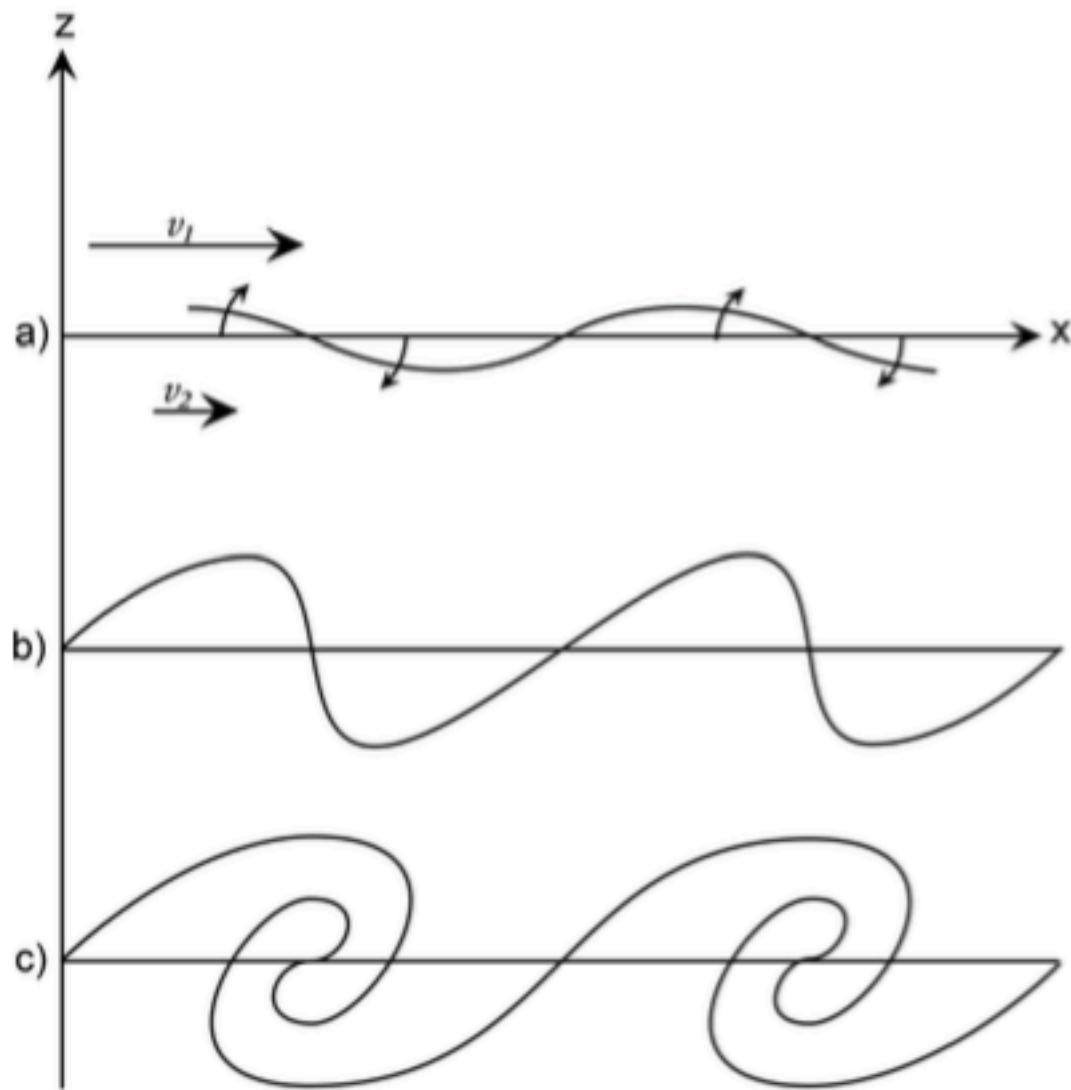


radius

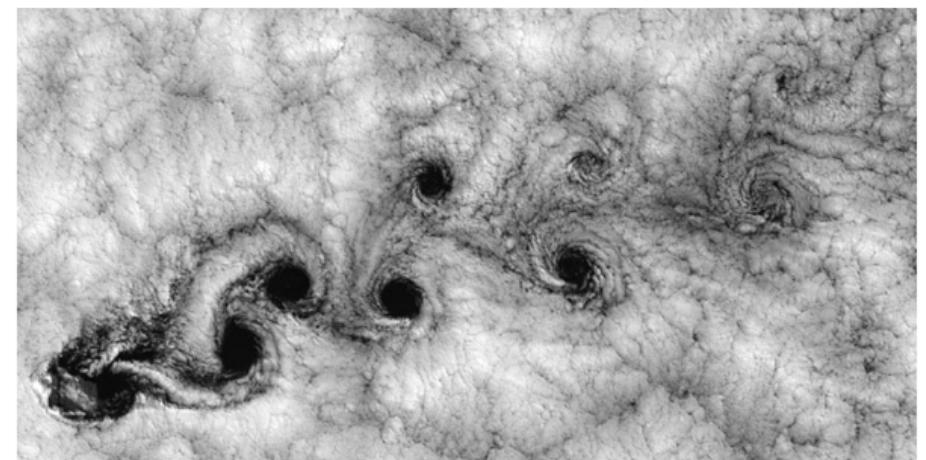
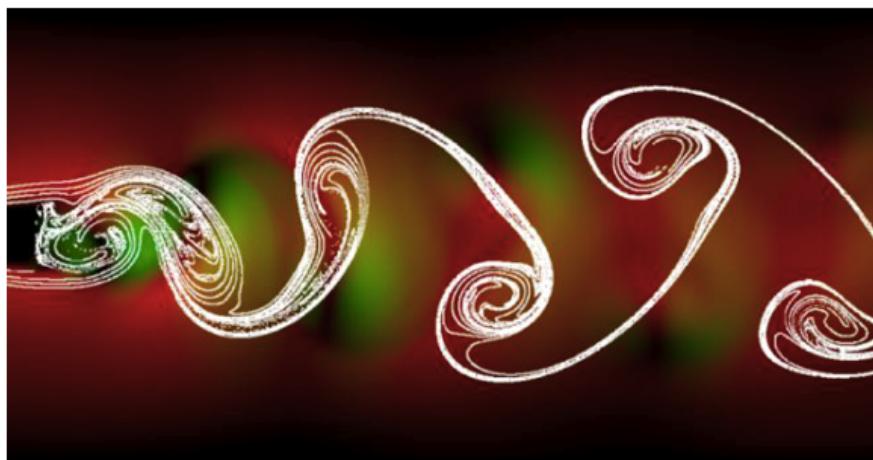
Lyra et al. (2008b, 2009a);  
See also Varniere & Tagger (2006)

# Rossby wave instability

(or... Kelvin-Helmholtz in differentially rotating disks)



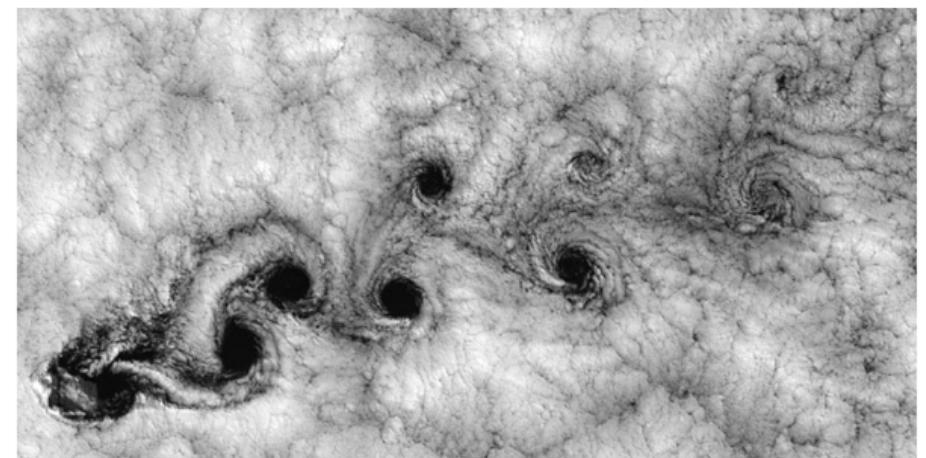
# Vortices – an ubiquitous fluid mechanics phenomenon



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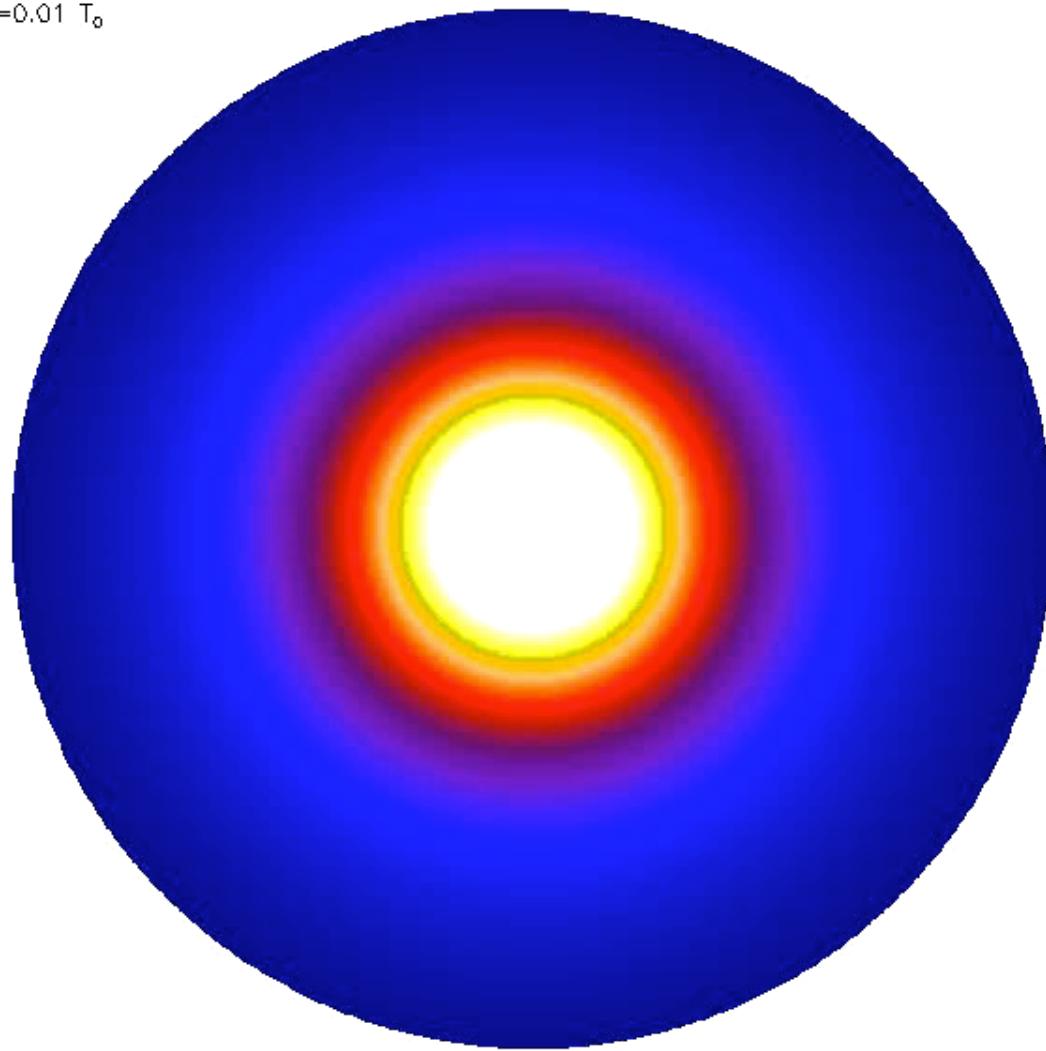
## Von Kármán *vortex street*





## Inner (0.1 AU) active/dead zone boundary

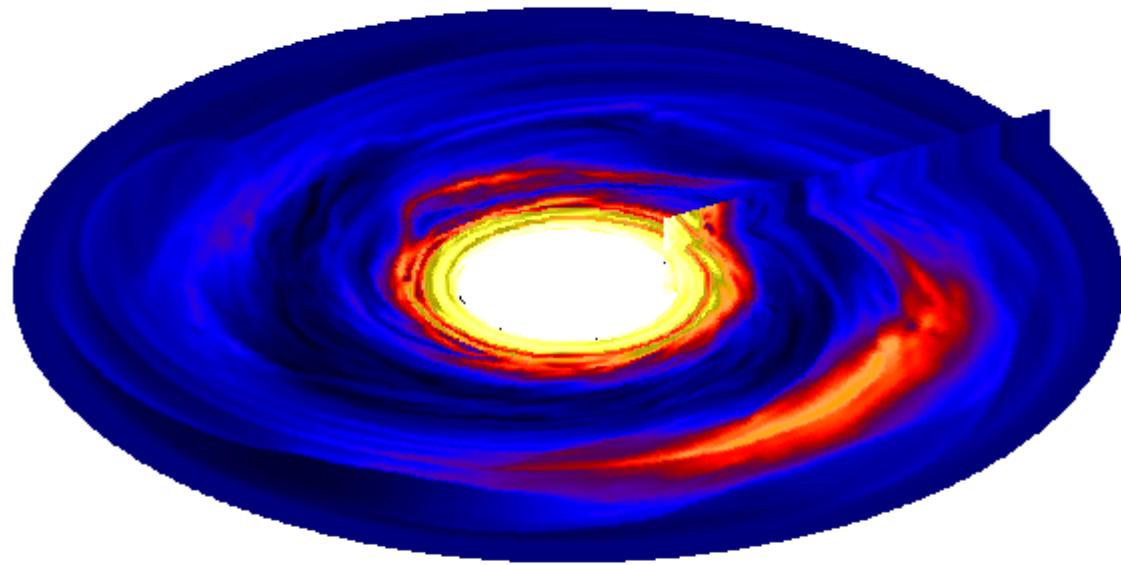
$t=0.01 T_0$



Magnetized inner disk + resistive outer disk  
Lyra & Mac Low (2012)

# Inner (0.1AU) active/dead zone boundary

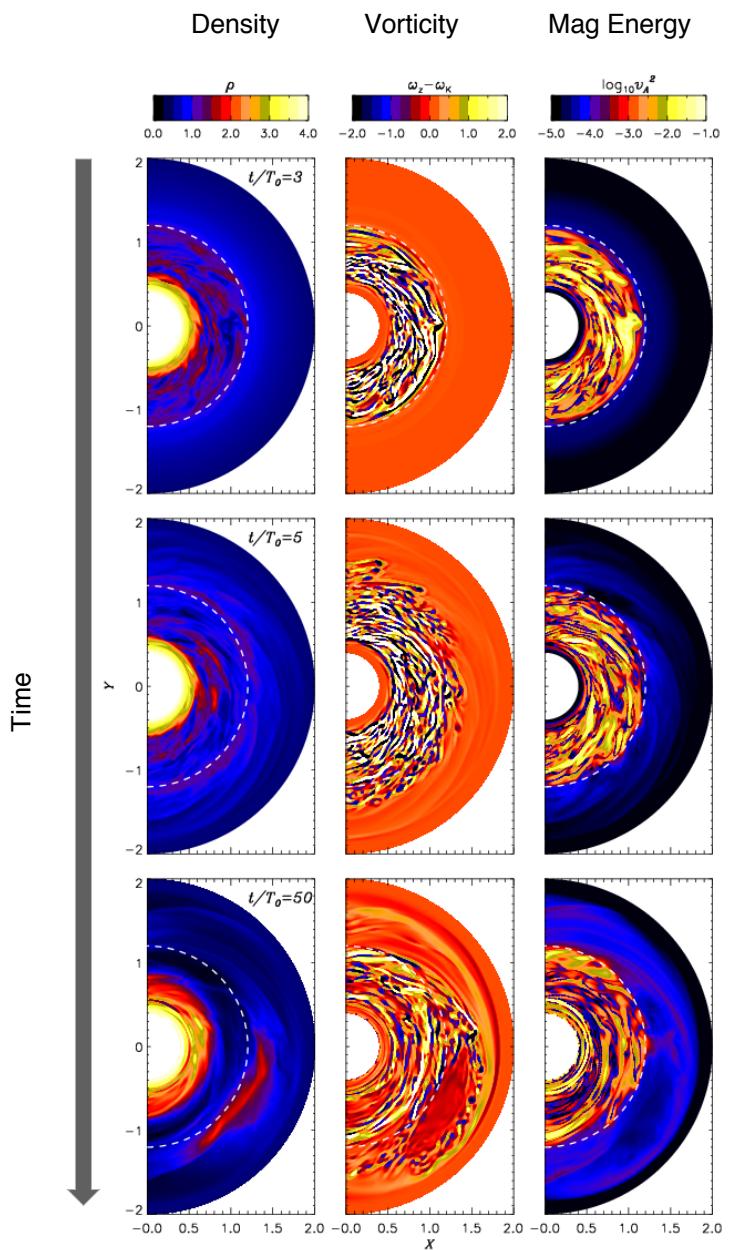
$t=22.28 T_0$



0.00 2.00 4.00  
 $\rho$

Magnetized inner disk + resistive outer disk

Lyra & Mac Low (2012)

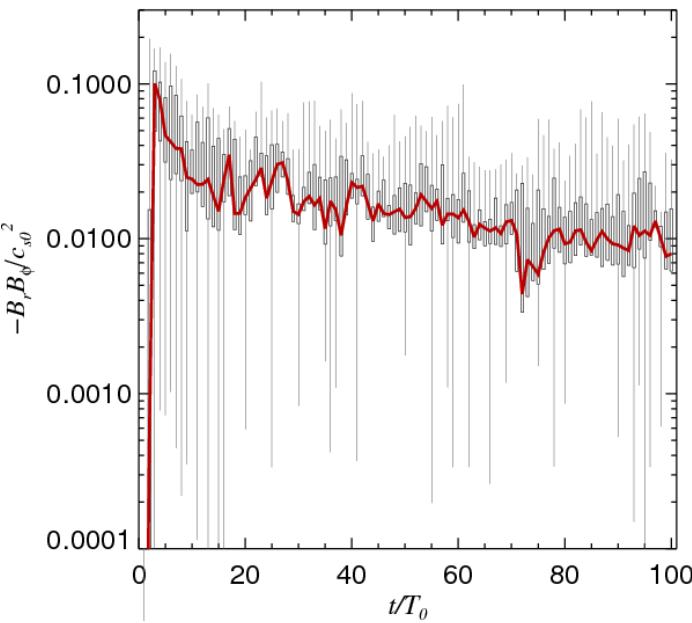


# Inner (0.1AU) active/dead zone boundary

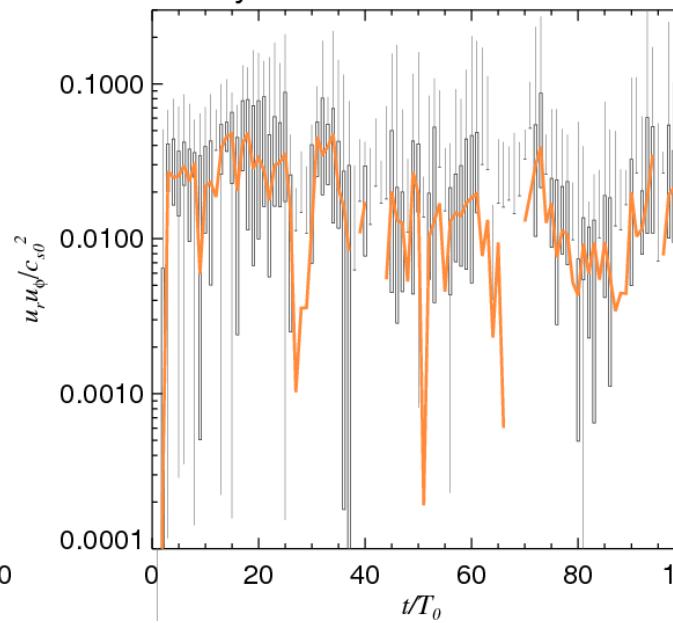
Active zone

Dead zone

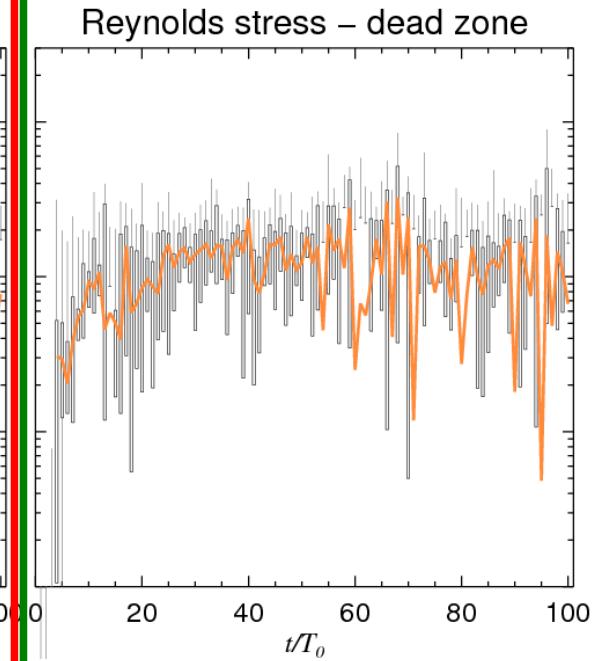
Maxwell stress – active zone



Reynolds stress – active zone



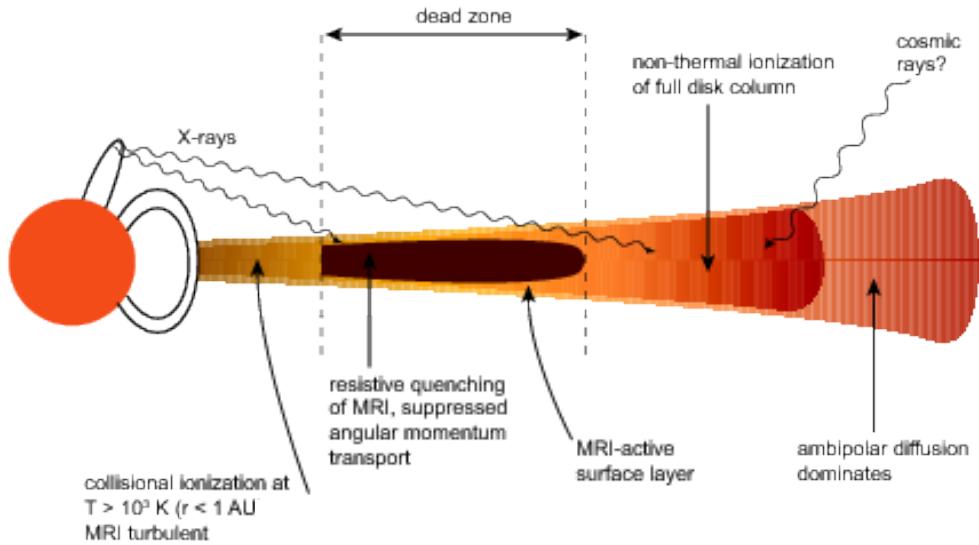
Reynolds stress – dead zone



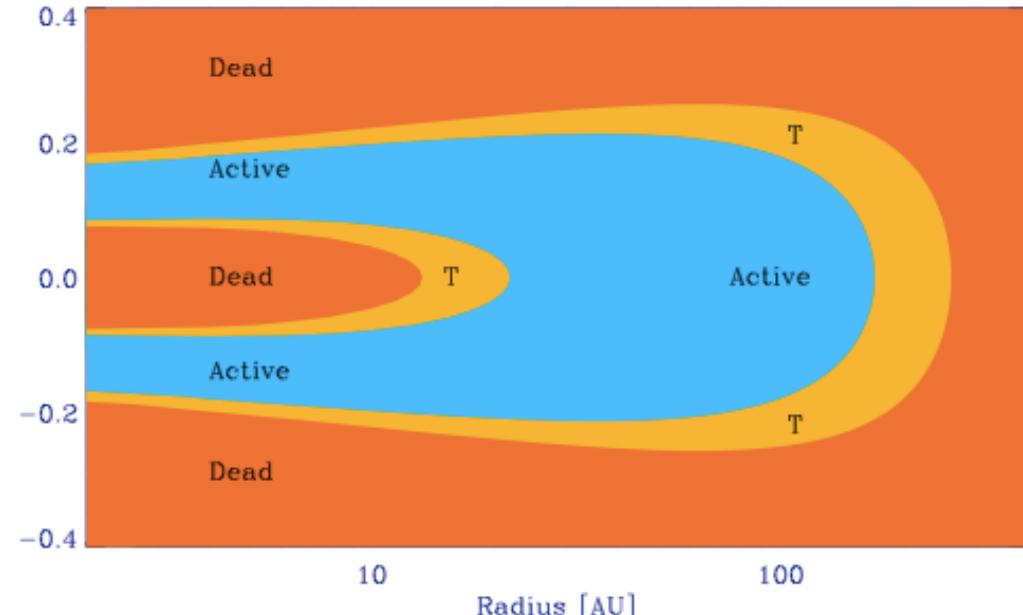
$$\alpha \sim 10^{-3} - 10^{-2}$$

Significant angular momentum transport

# Outer Dead/Active zone transition KHI



Armitage (2010)

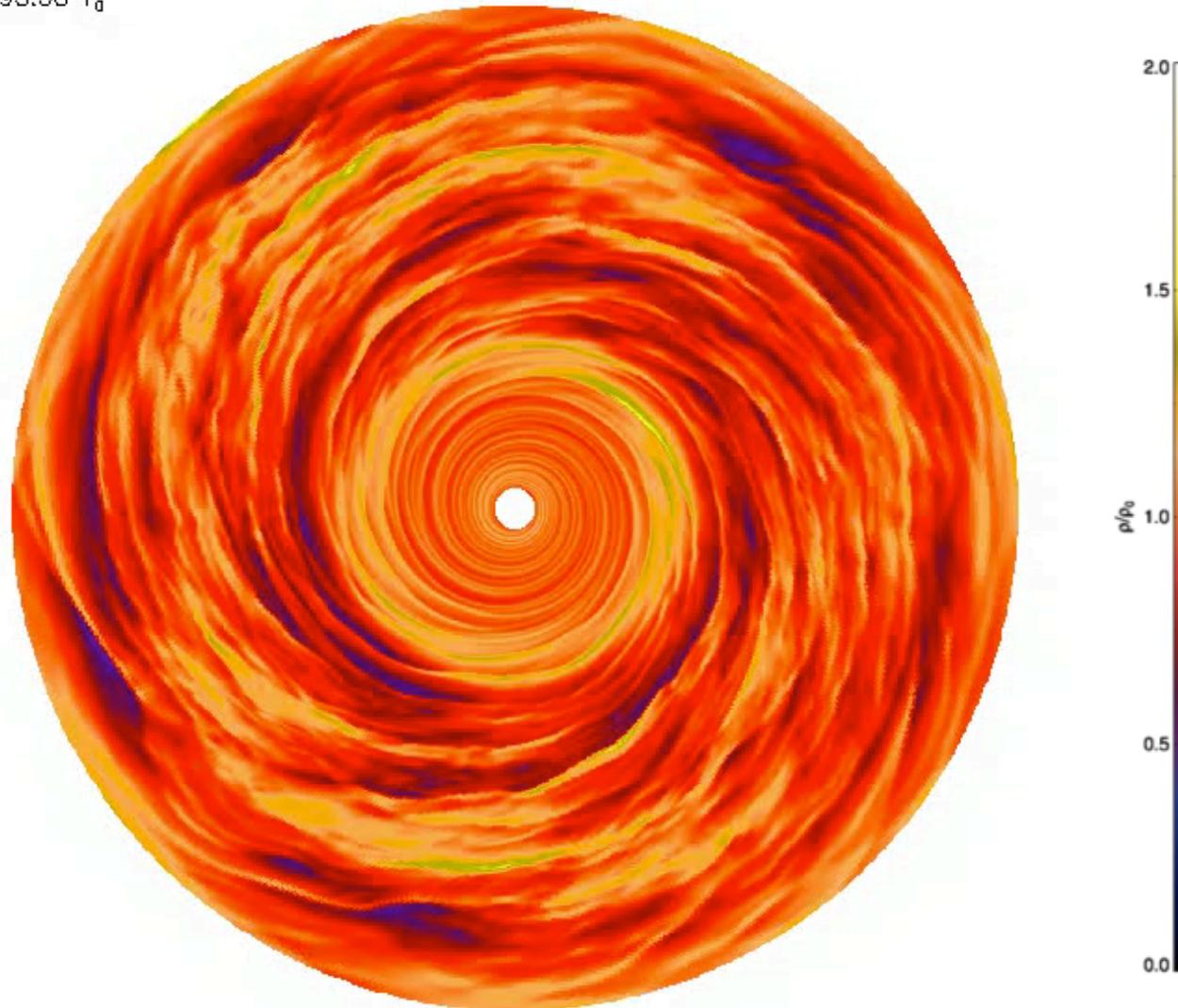


Dzyurkevitch et al (2013)

The **outer** dead zone transition in ionization supposed  
**TOO SMOOTH**  
to generate an KH-unstable bump.

# Outer Dead/Active zone transition: 3D MHD

$t=95.58 T_0$

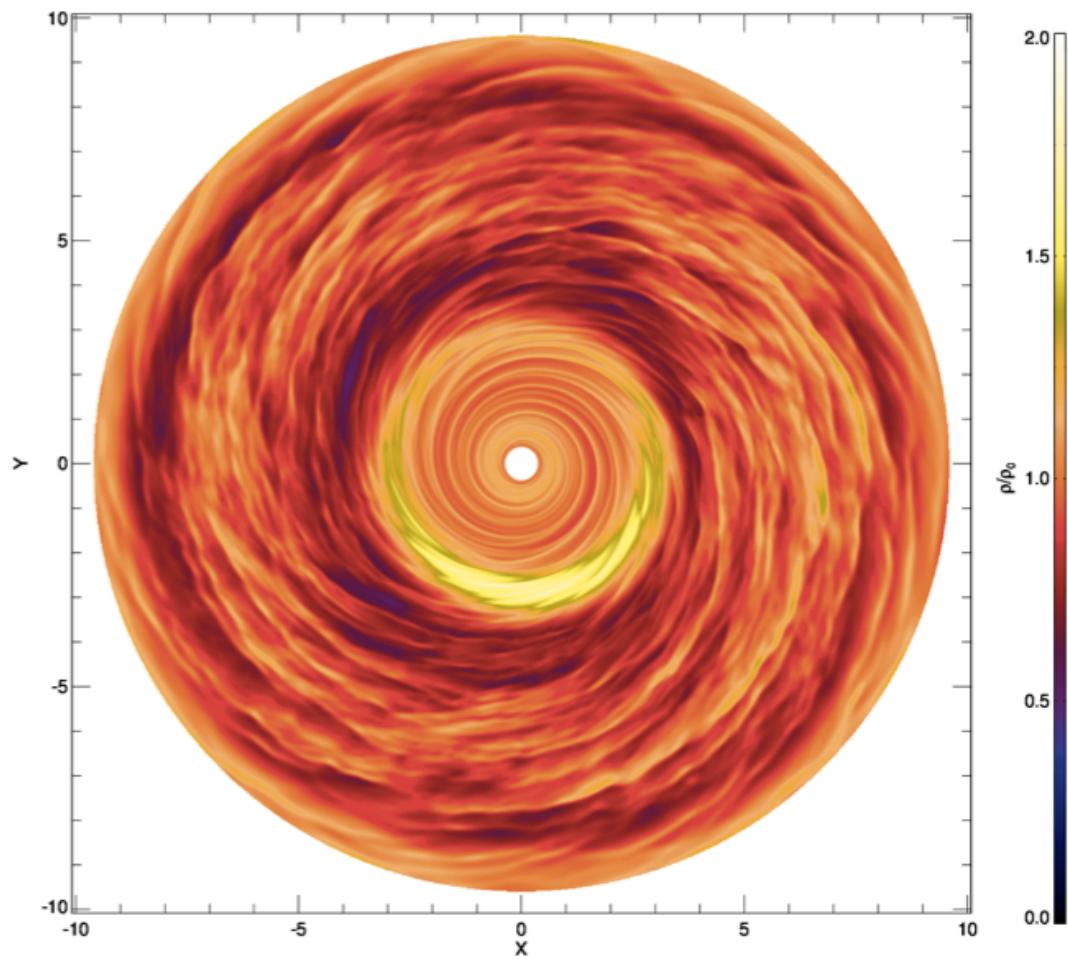
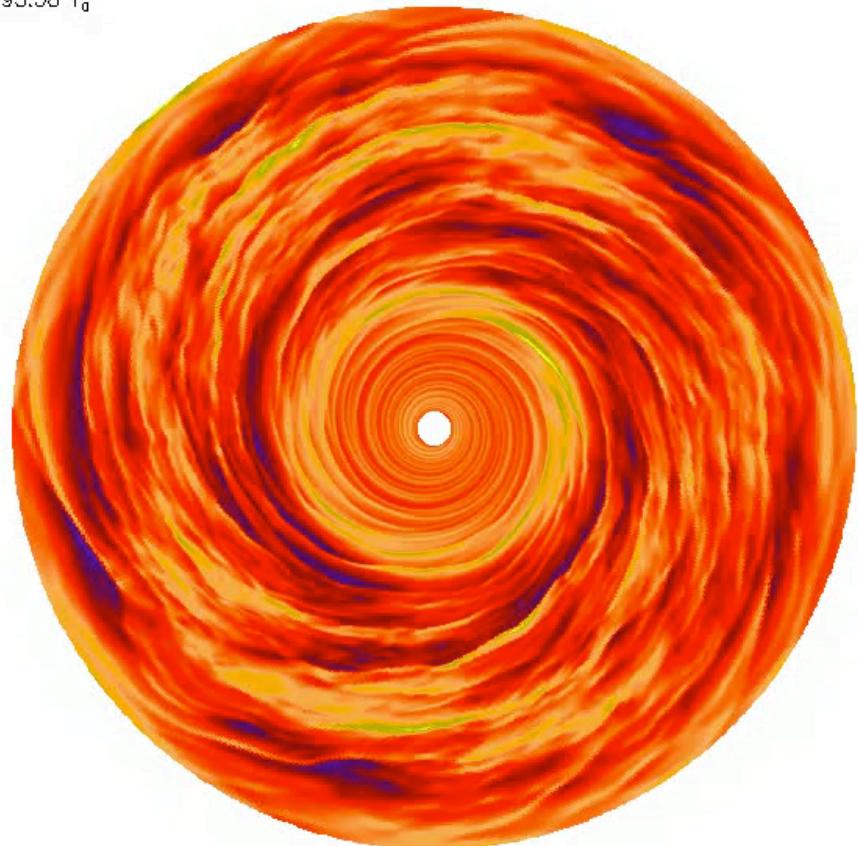


Resistive inner disk + magnetized outer disk

Lyra, Turner, & McNally (2015)

# Outer Dead/Active zone transition KHI

$t=95.58 T_0$

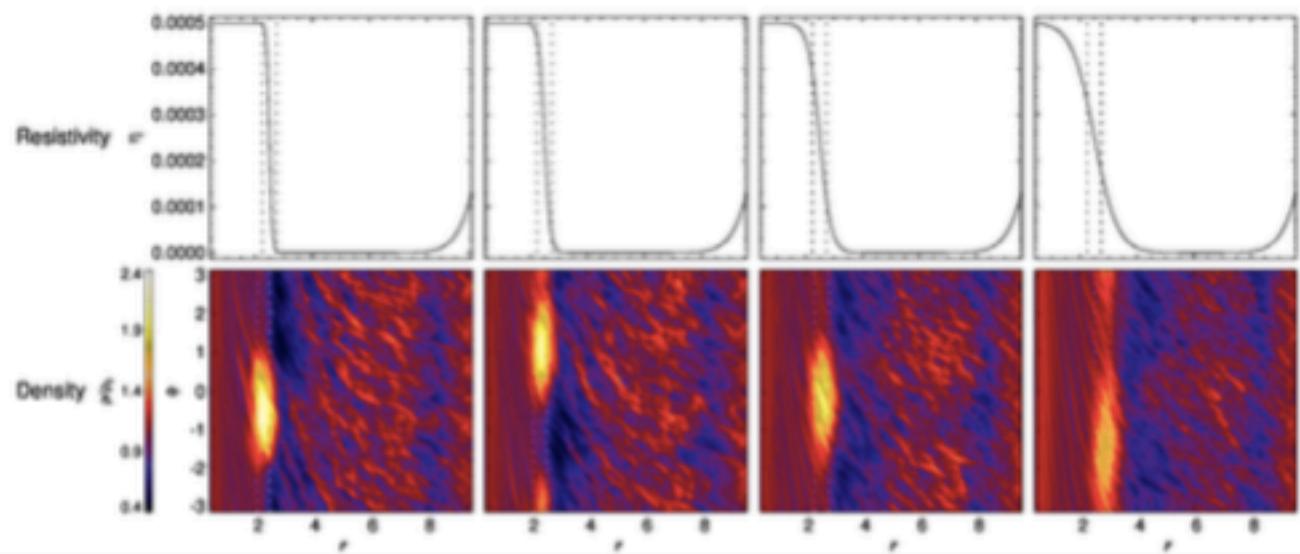


Resistive inner disk + magnetized outer disk

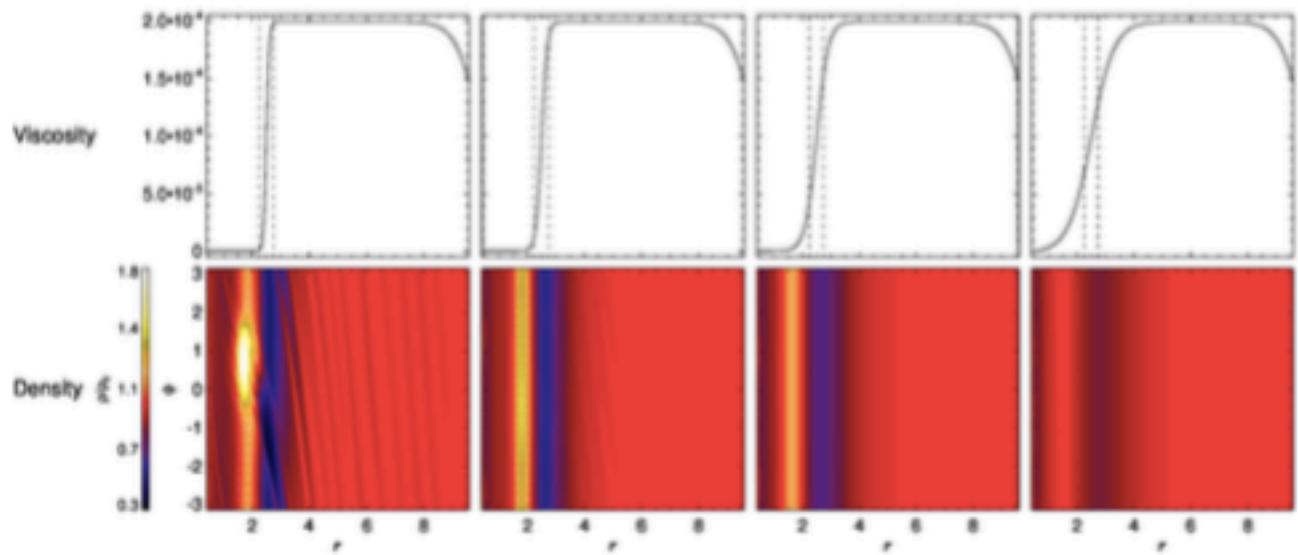
Lyra, Turner, & McNally (2015)

# Outer Dead/Active zone transition RWI

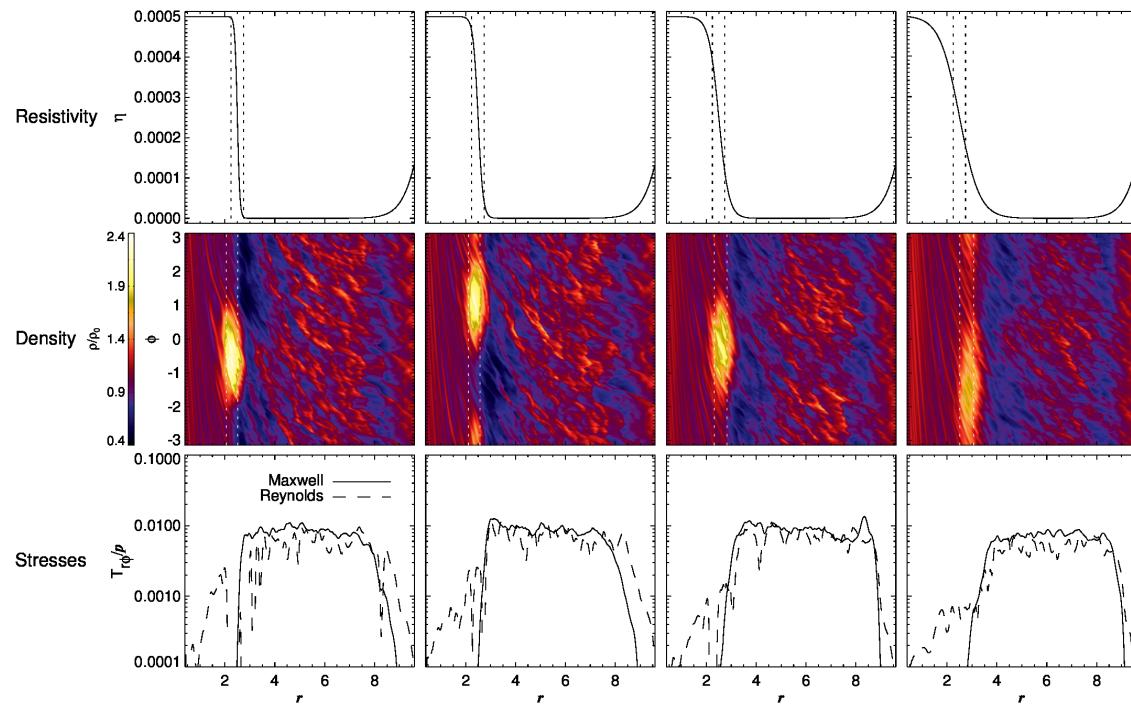
MHD



Hydro



# Outer Dead/Active zone transition RWI



Lyra, Turner, & McNally (2015)

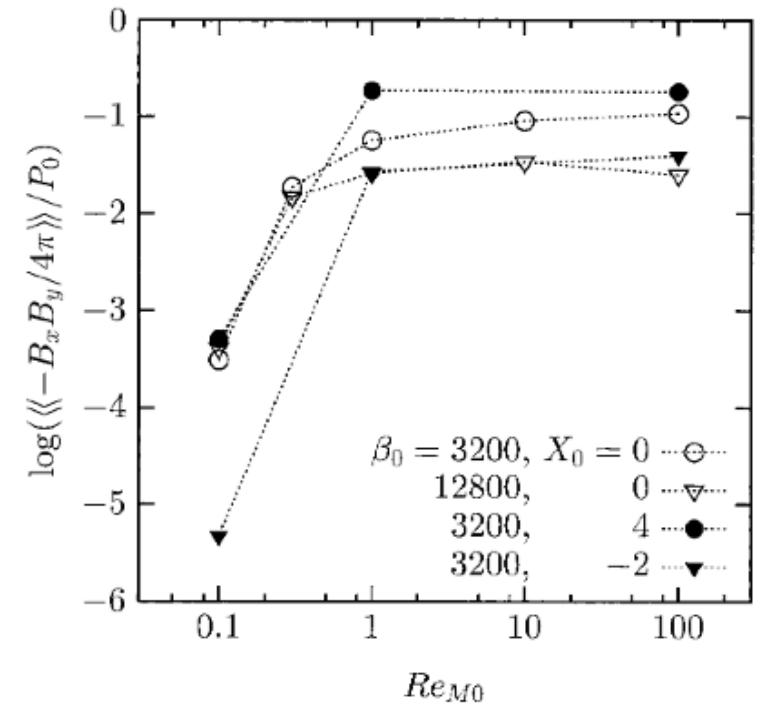
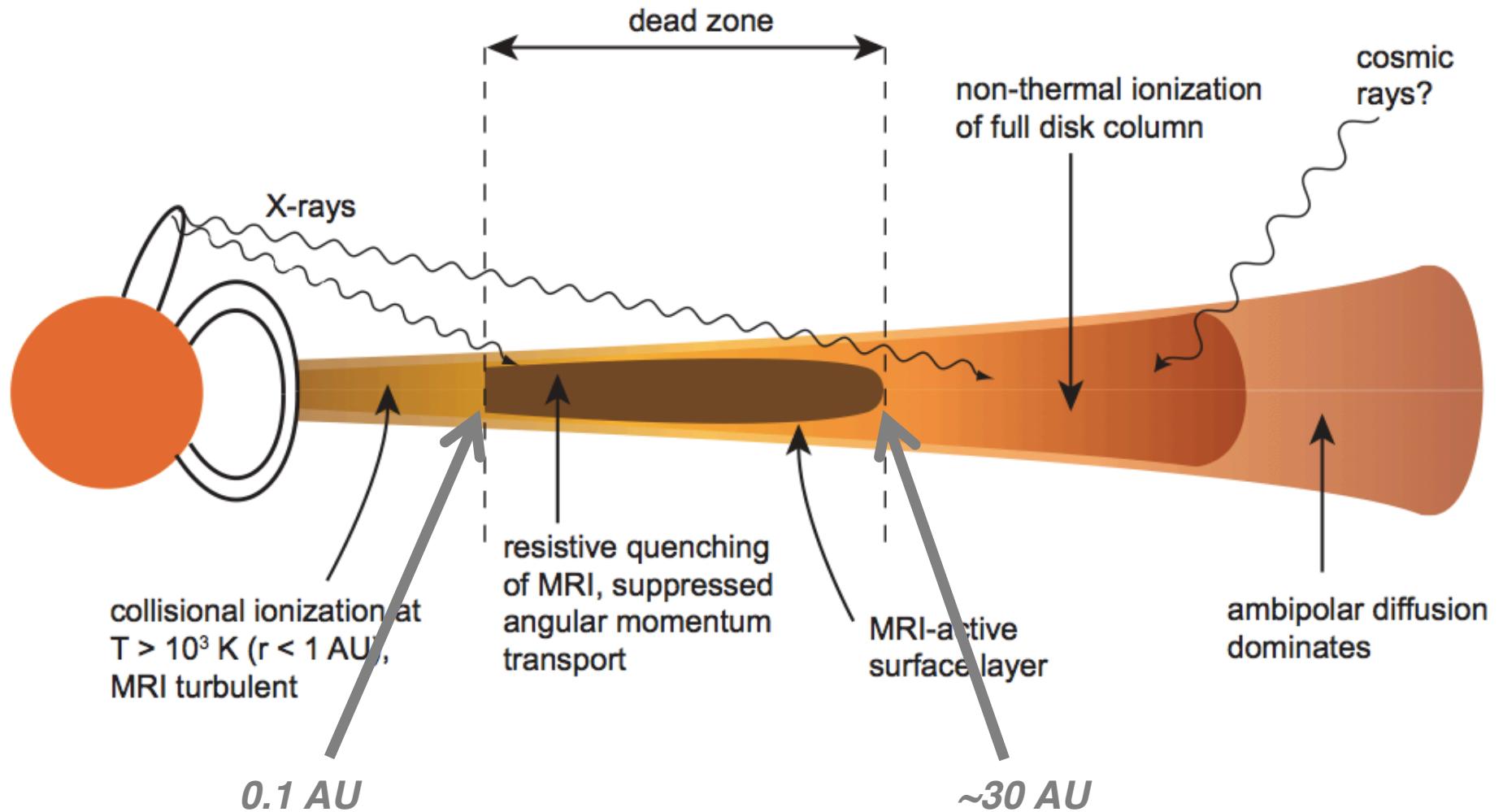


FIG. 9.—Saturation level of the Maxwell stress as a function of the magnetic Reynolds number  $Re_{M0}$ . Open circles and triangles denote the models without Hall term ( $X_0 = 0$ ) for  $\beta_0 = 3200$  and 12,800, respectively. The models including the Hall term are shown by filled circles ( $X_0 = 4$ ) and triangles ( $X_0 = -2$ ).

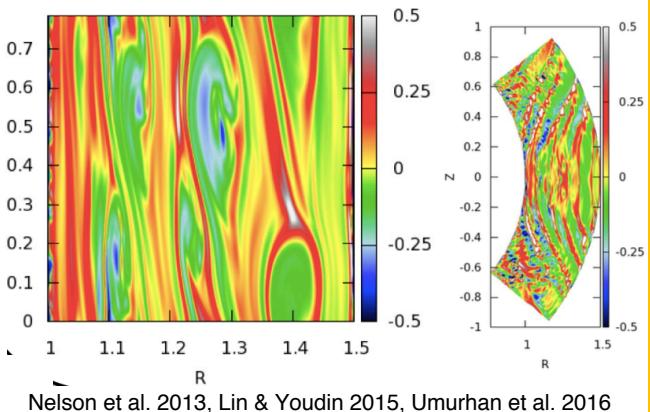
Sano and Stone (2002)

# Dead zones

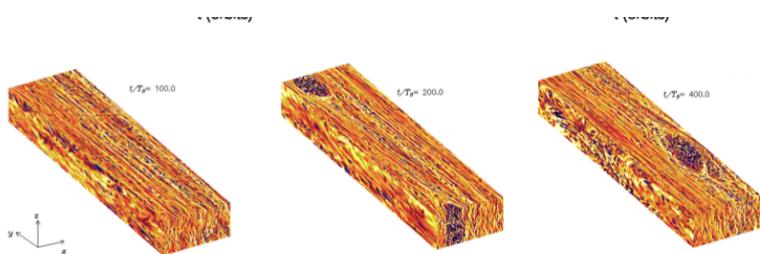


# Thermal Instabilities

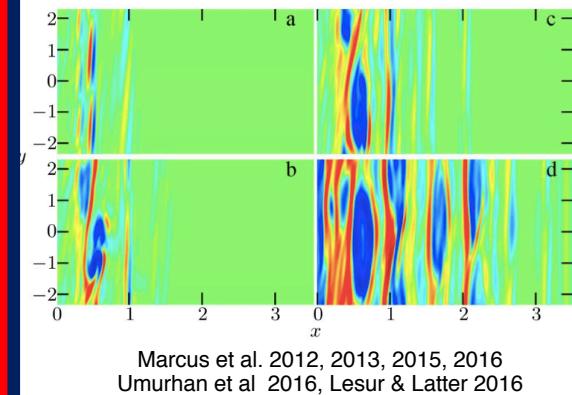
## Vertical Shear Instability



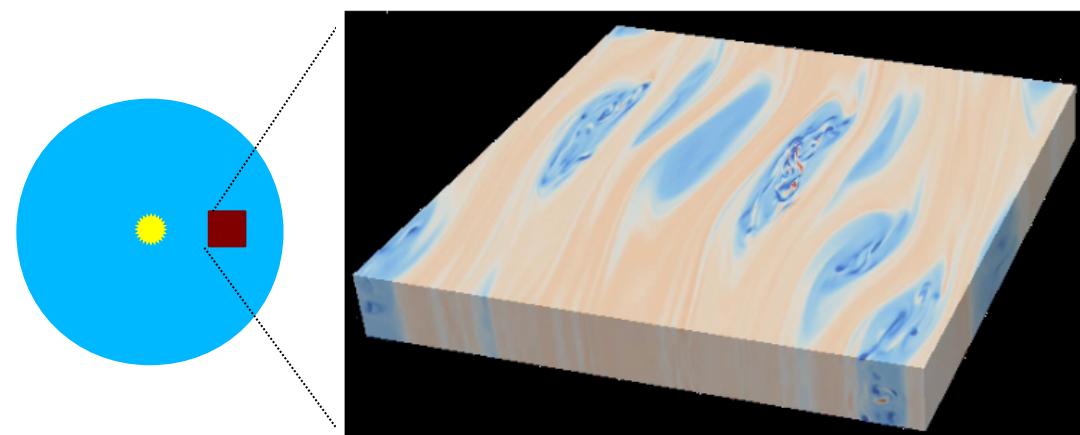
## Convective Overstability



## Zombie Vortex Instability

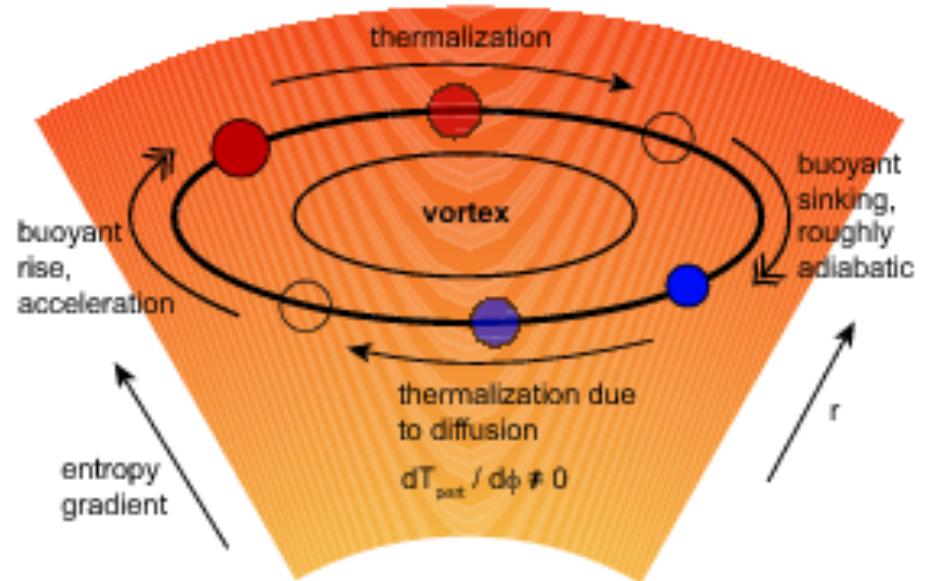


# Convective Overstability (née “Subcritic Baroclinic Instability”)



Lesur & Papaloizou (2010)

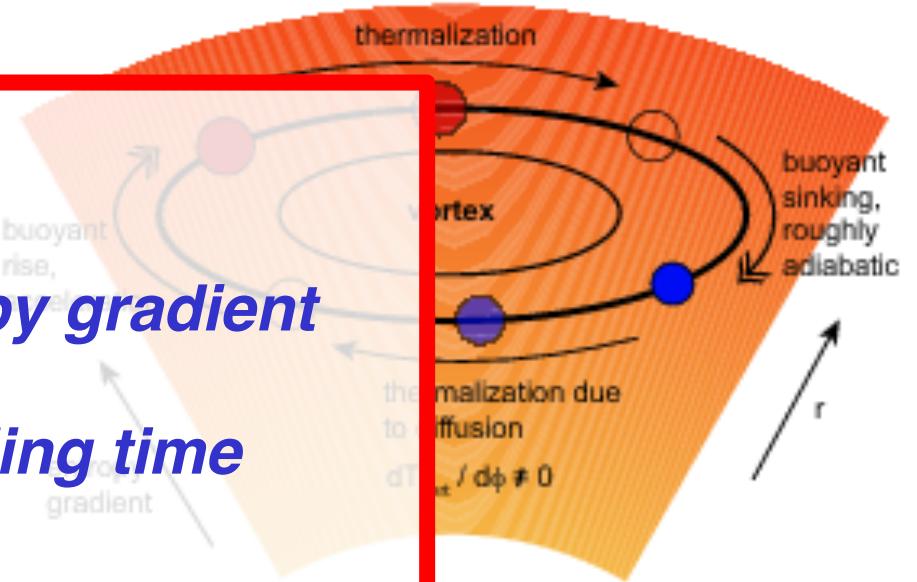
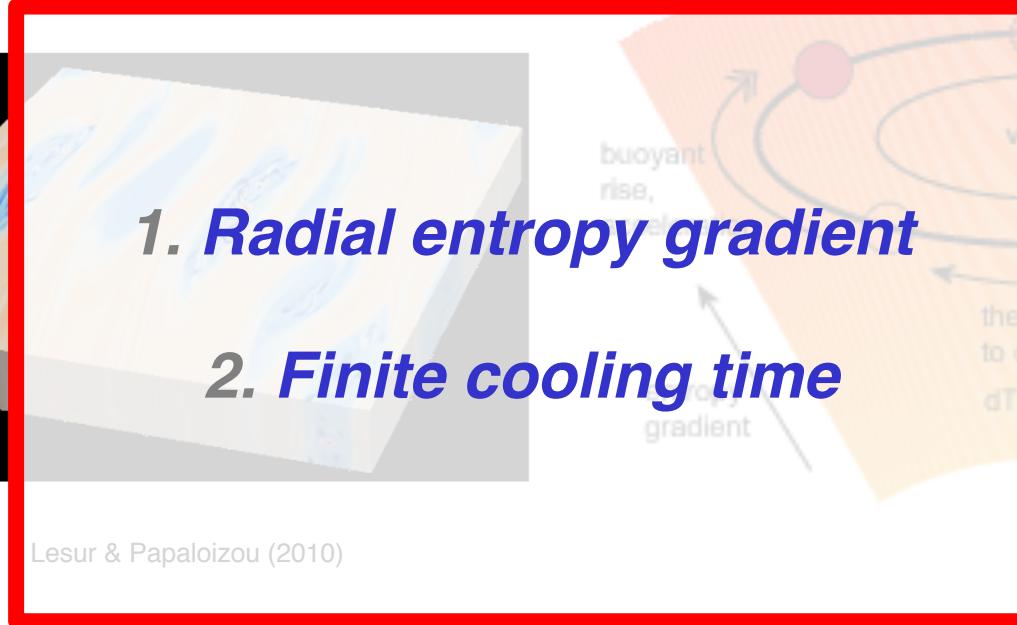
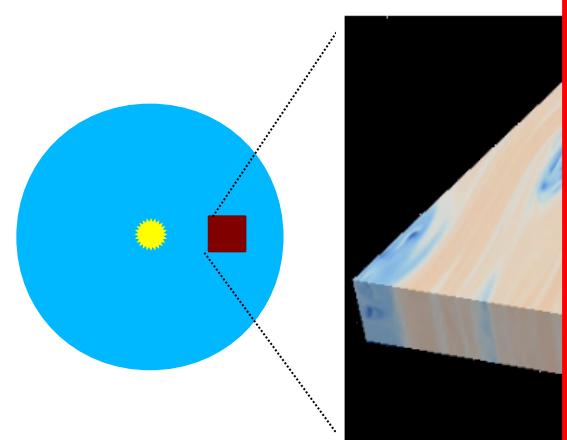
Sketch of the  
Subcritic Baroclinic Instability



Armitage (2010)

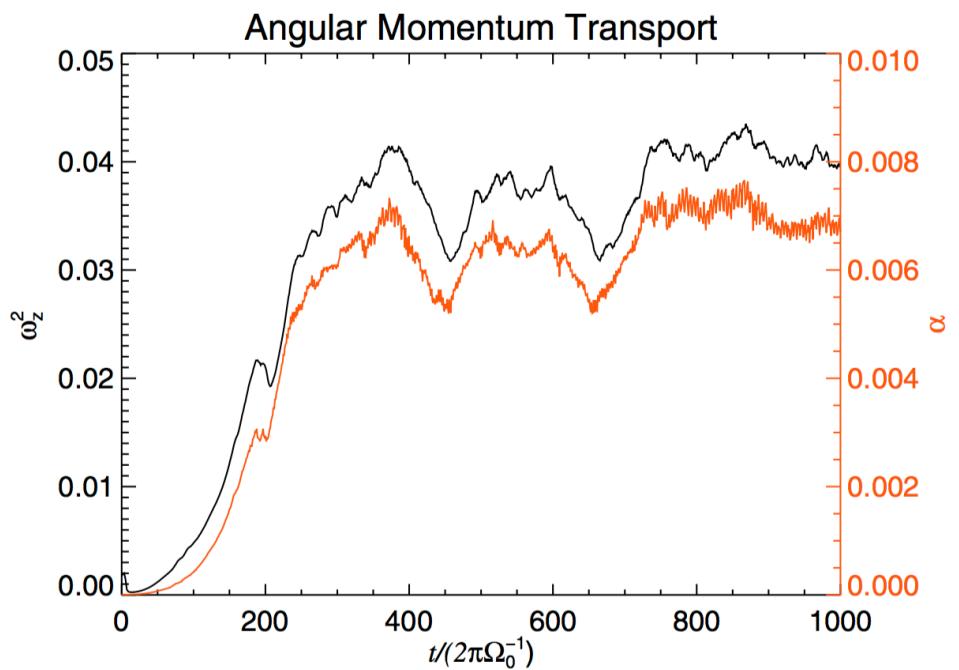
# Convective Overstability (née “Subcritic Baroclinic Instability”)

Sketch of the  
Subcritic Baroclinic Instability



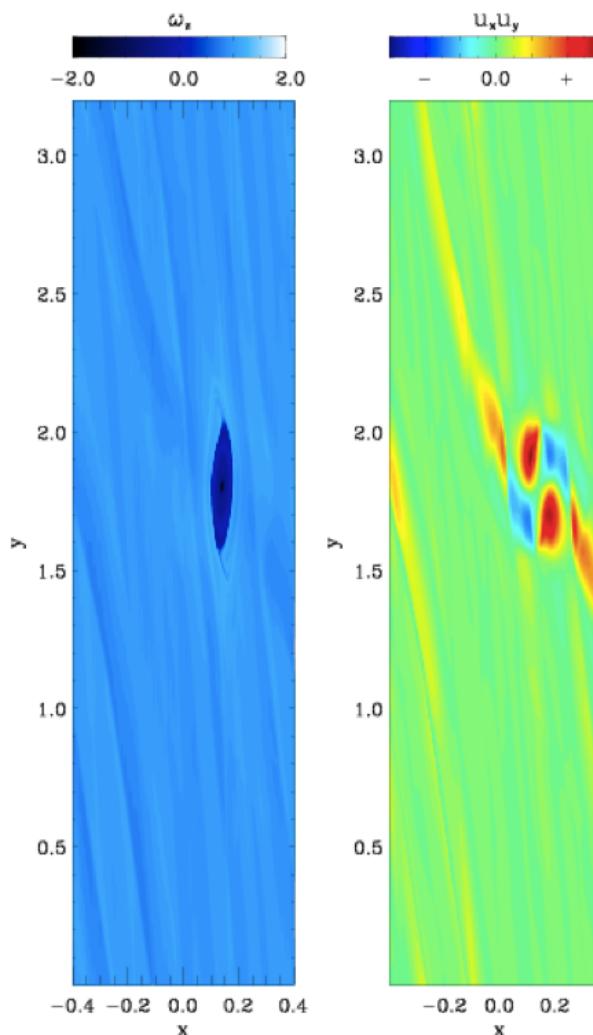
# Subcritic baroclinic instability and Accretion

$$\alpha \sim 0.005$$



Lyra & Klahr (2011)

Raettig, Lyra, & Klahr (2013)



The angular momentum is carried  
by **waves** excited by the vortex

(see also Heinemann & Papaloizou 2008, 2009)

# Convective Overstability

Klahr & Hubbard (2014), Lyra (2014), Latter (2015)

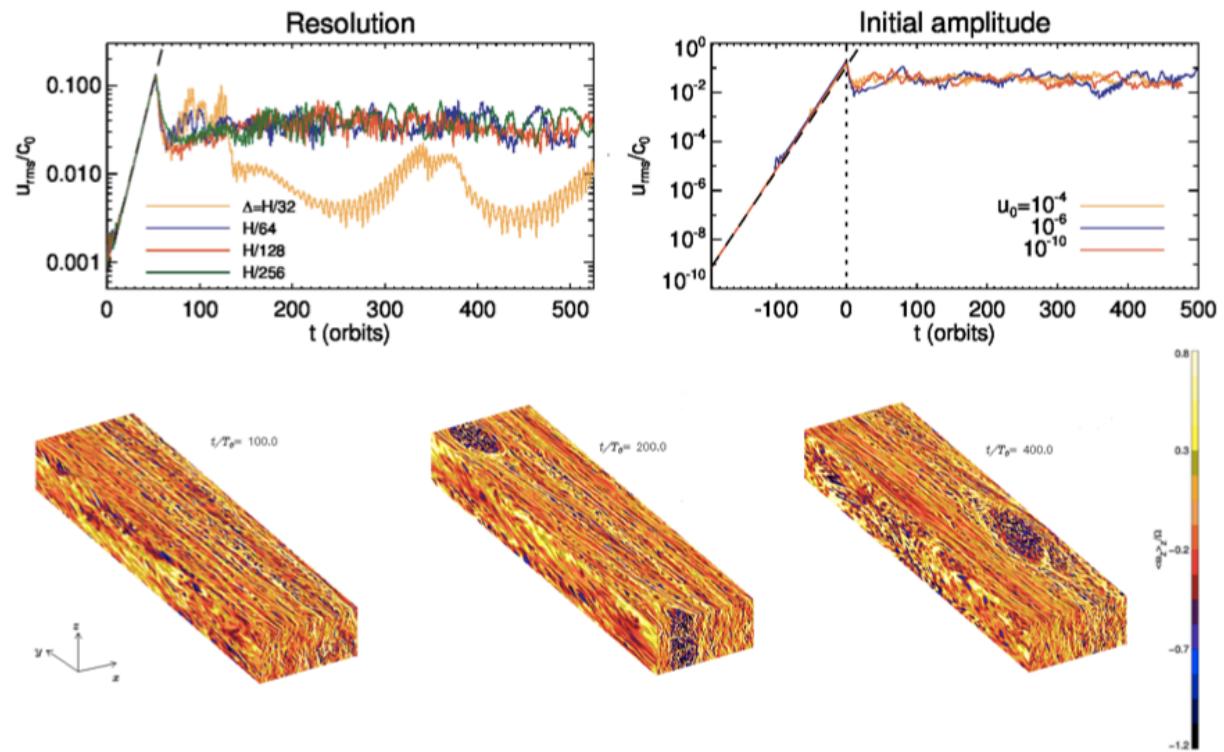
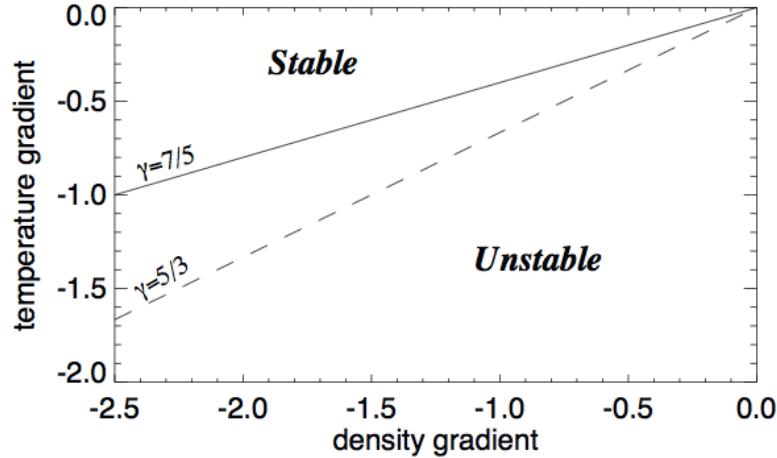
$$\begin{aligned}\frac{\partial \rho}{\partial t} + (\mathbf{u} \cdot \nabla) \rho &= -\rho \nabla \cdot \mathbf{u}, \\ \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} &= -\frac{1}{\rho} \nabla p + \mathbf{g}, \\ \frac{\partial p}{\partial t} + (\mathbf{u} \cdot \nabla) p &= -\gamma p \nabla \cdot \mathbf{u} - \frac{p}{T} \frac{(T - T_0)}{\tau},\end{aligned}$$

$$\bar{\omega}^3 + i\zeta\bar{\omega}^2 - \bar{\omega}\mu^2(\kappa^2 + N^2) - i\zeta\kappa^2\mu^2 = 0,$$

$$\zeta = 1/\gamma\tau \quad \mu^2 = k_z^2/k^2.$$

$$\tau_{\max} = \frac{1}{\gamma} \left| \frac{k}{k_z} \right| \frac{1}{\sqrt{\kappa^2 + N^2}}$$

$$\sigma_{\max} = -\frac{1}{4} \left| \frac{k_z}{k} \right| \frac{N^2}{\sqrt{\kappa^2 + N^2}}$$



Lyra (2014)

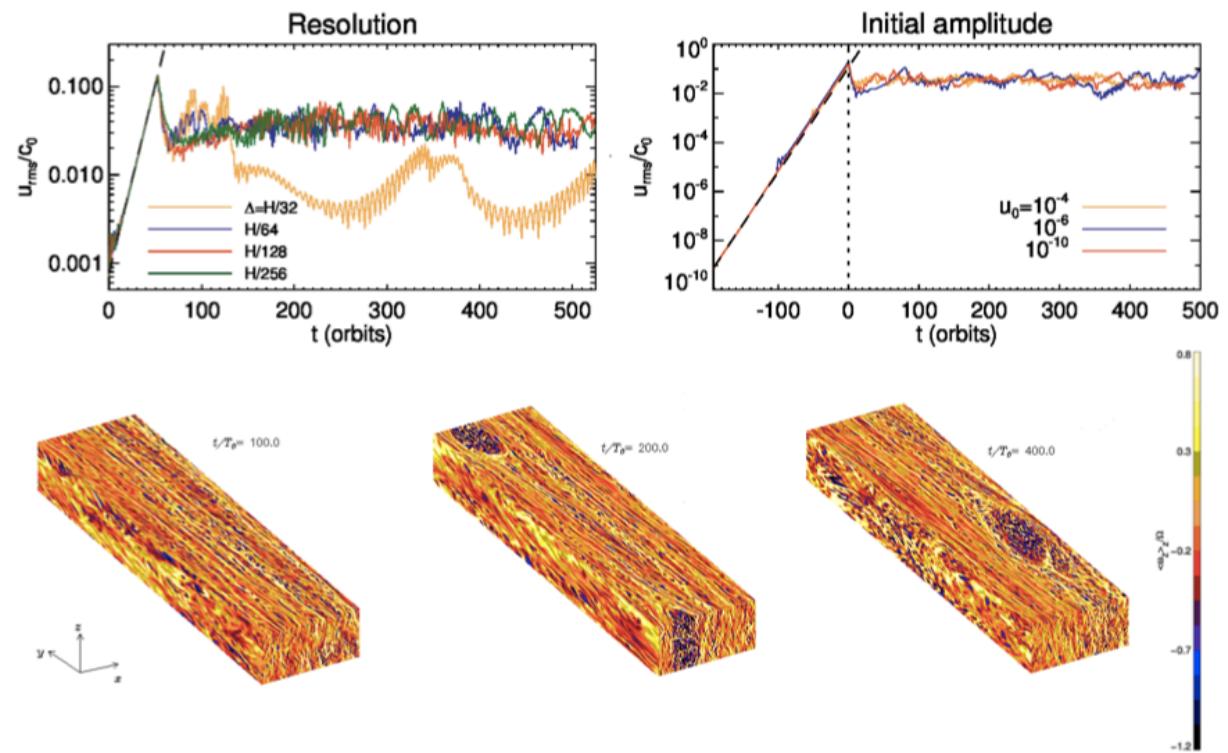
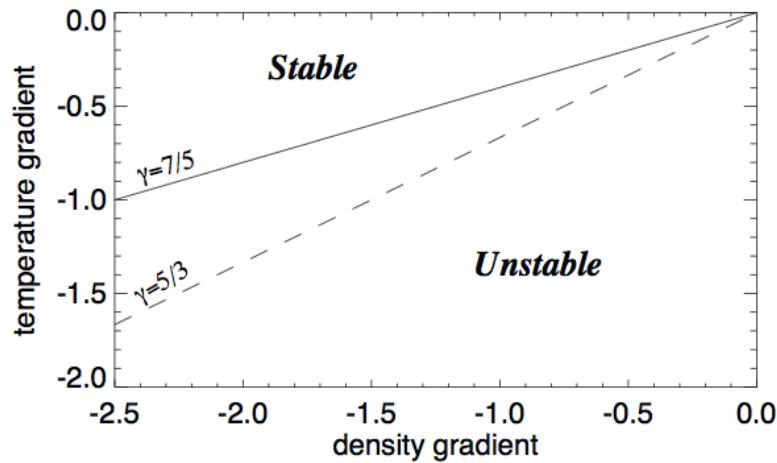
# Convective Overstability

Klahr & Hubbard (2014), Lyra (2014), Latter (2015)

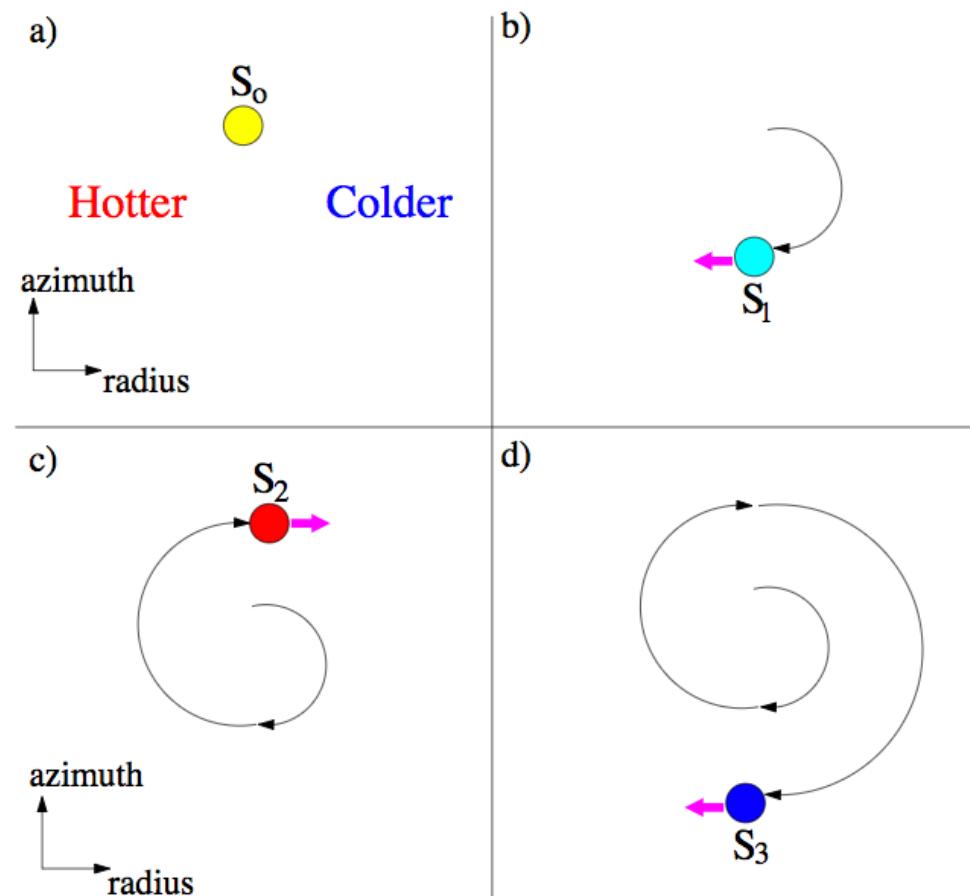
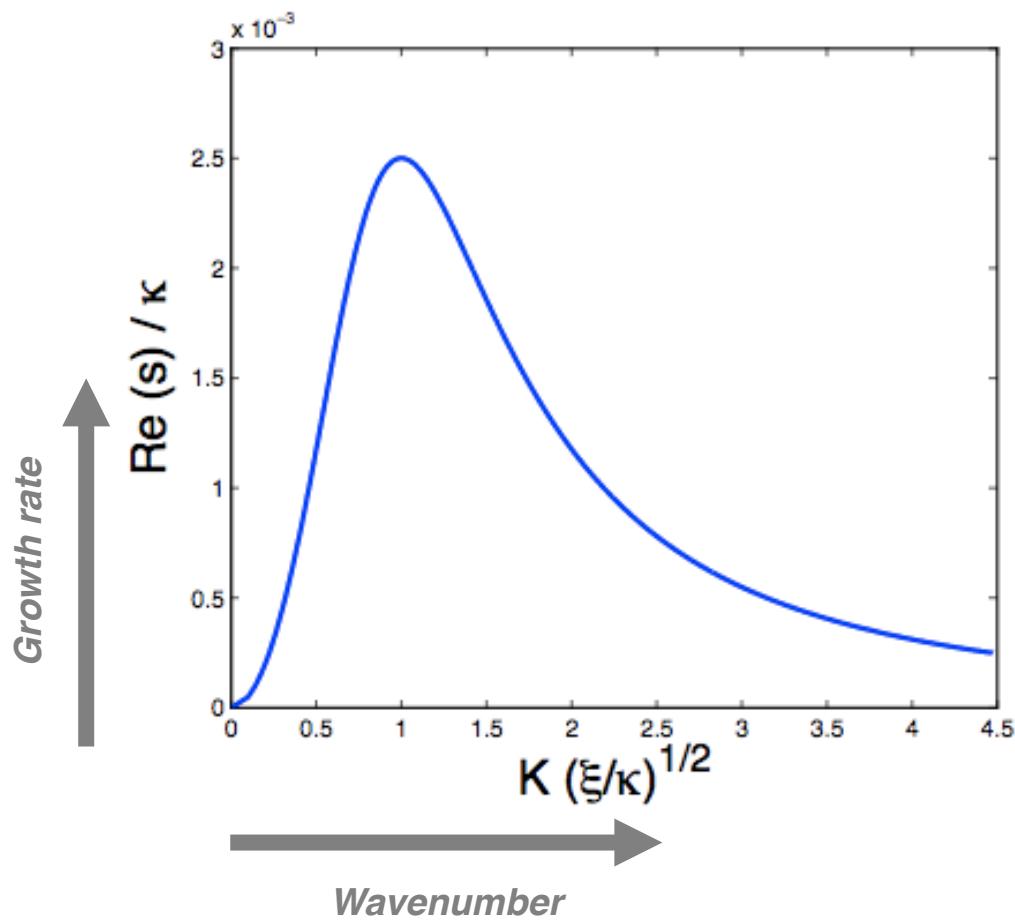
$$\begin{aligned}\frac{\partial \rho}{\partial t} + (\mathbf{u} \cdot \nabla) \rho &= -\rho \nabla \cdot \mathbf{u}, \\ \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} &= -\frac{1}{\rho} \nabla p + \mathbf{g}, \\ \frac{\partial p}{\partial t} + (\mathbf{u} \cdot \nabla) p &= -\gamma p \nabla \cdot \mathbf{u} - \frac{p}{T} \frac{(T - T_0)}{\tau},\end{aligned}$$

$$\tau_{\max} = \frac{1}{\gamma \Omega};$$

$$\sigma_{\max} = -\frac{N^2}{4\Omega}.$$



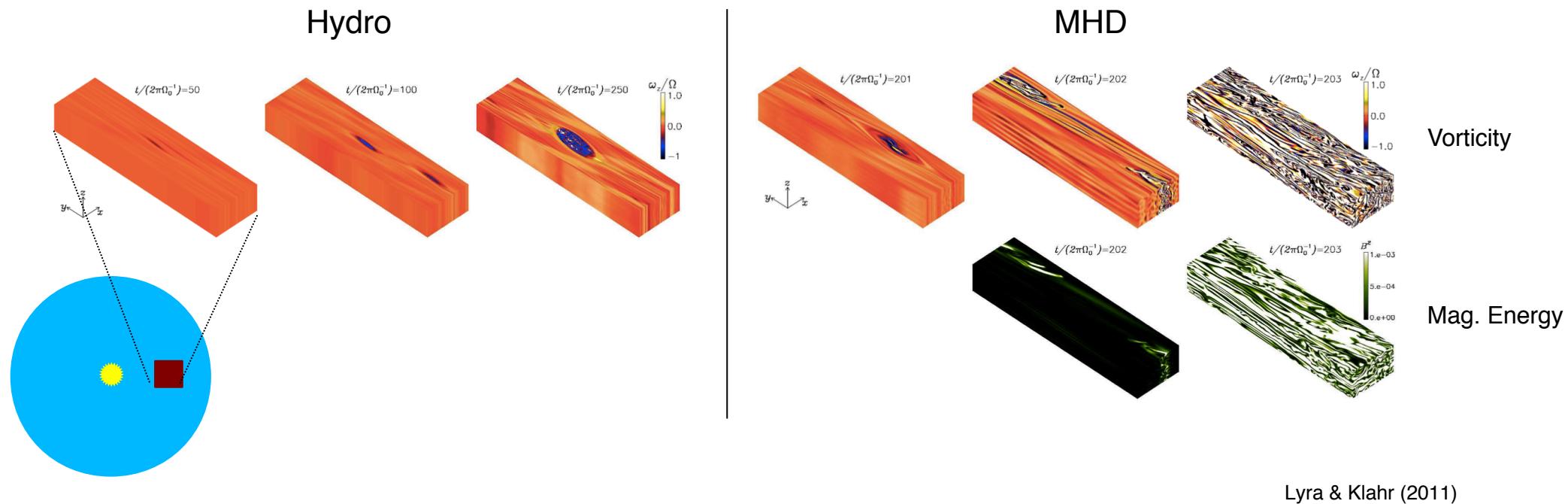
# Convective Overstability



**Figure 2.** Four panels indicating the convective overstability mechanism. In panel (a) a fluid blob is embedded in a radial entropy gradient. In panel (b) it undergoes half an epicycle and returns to its original radius with a smaller entropy than when it began  $S_1 < S_0$ . It hence feels a buoyancy acceleration inwards and the epicycle is amplified. The process occurs in reverse once the epicycle is complete, shown in panel (c), where now  $S_2 > S_0$ . The oscillations hence grow larger and larger.

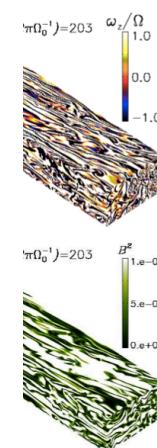
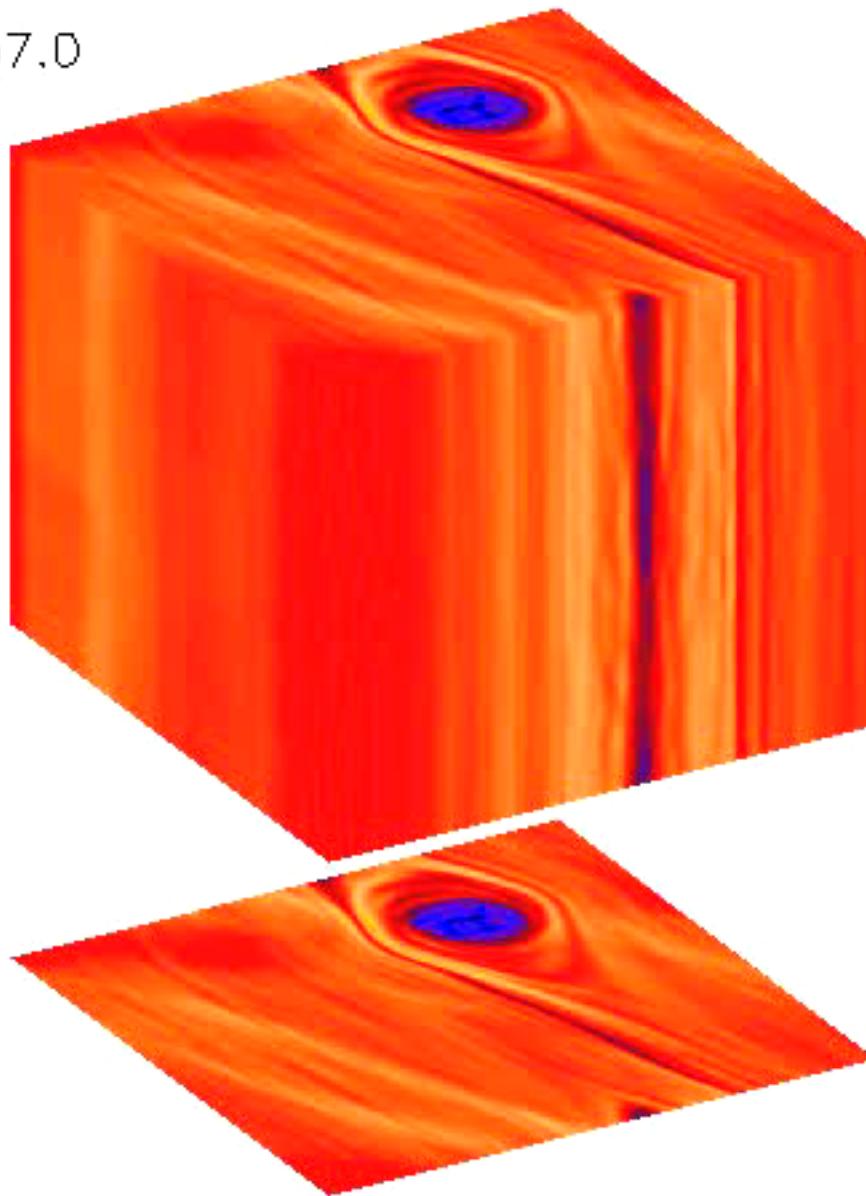
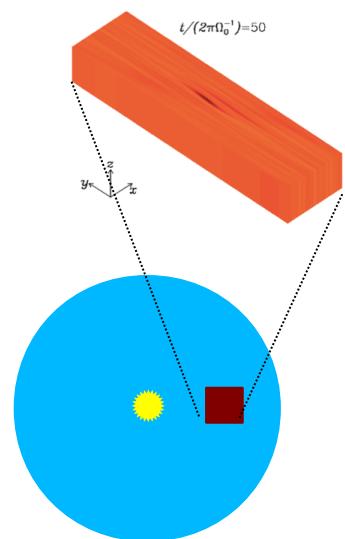
# COV/SBI and MRI

What happens when the disk is magnetized?

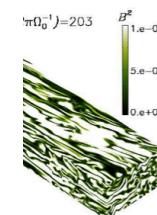


# COV/SBI and MRI

$t=1257.0$

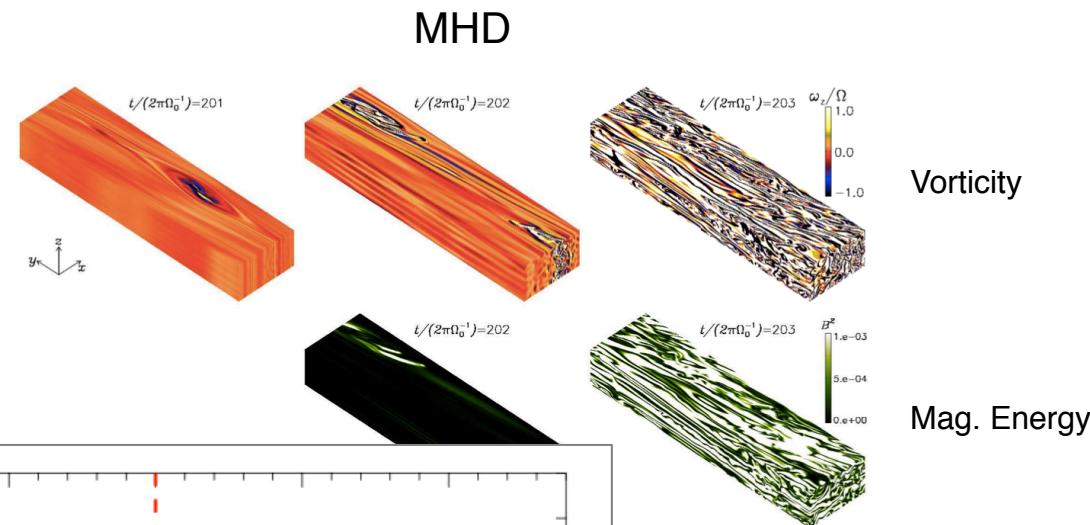
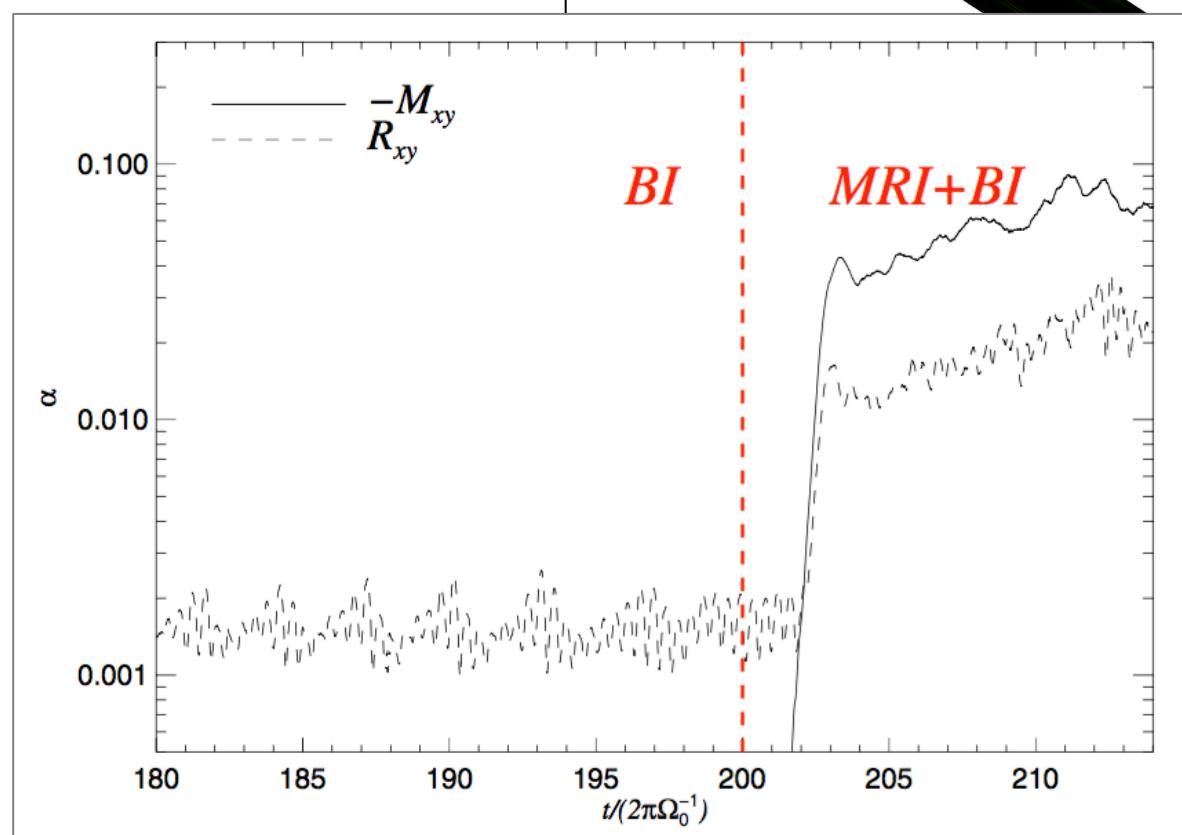
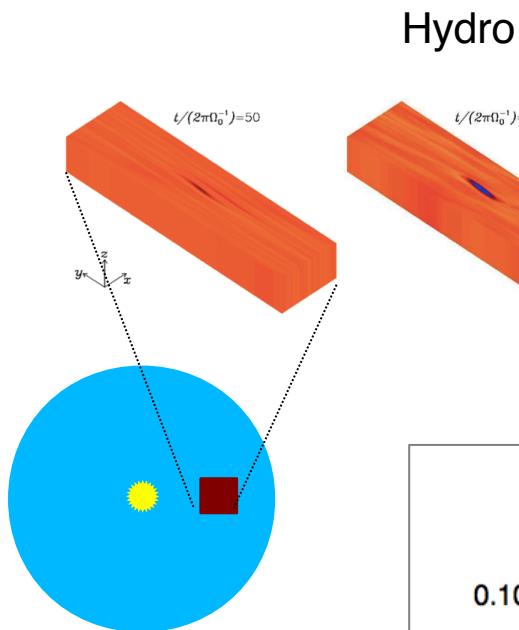


Vorticity



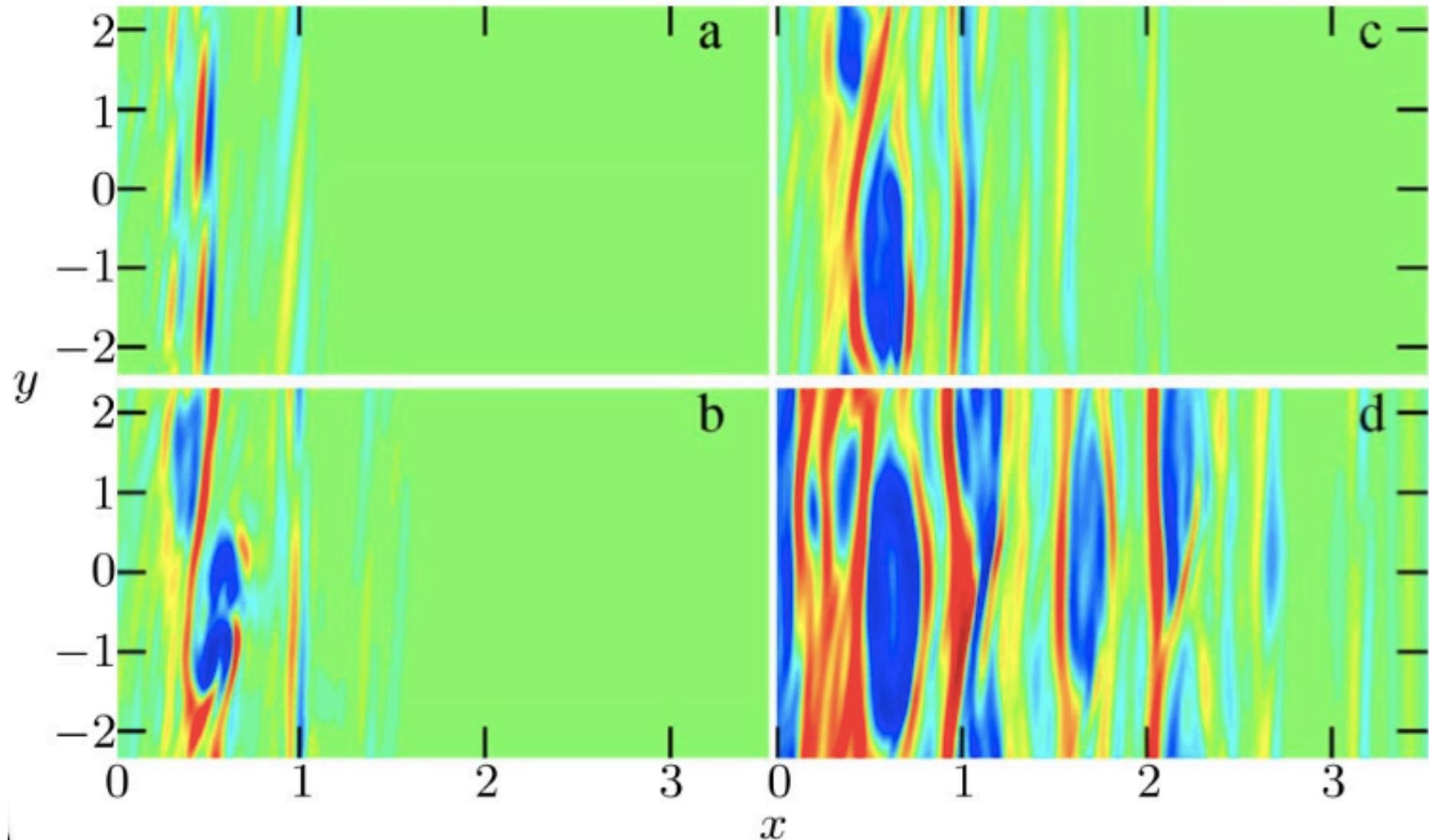
Mag. Energy

# COV/SBI and MRI



Lyra & Klahr (2011)

# Zombie Vortex Instability



Cascade of baroclinic critical layers

"Because the vortices arise from these dead zones, and because new generations of giant vortices march across these dead zones, we affectionately refer to them as 'zombie vortices,'" said Marcus.

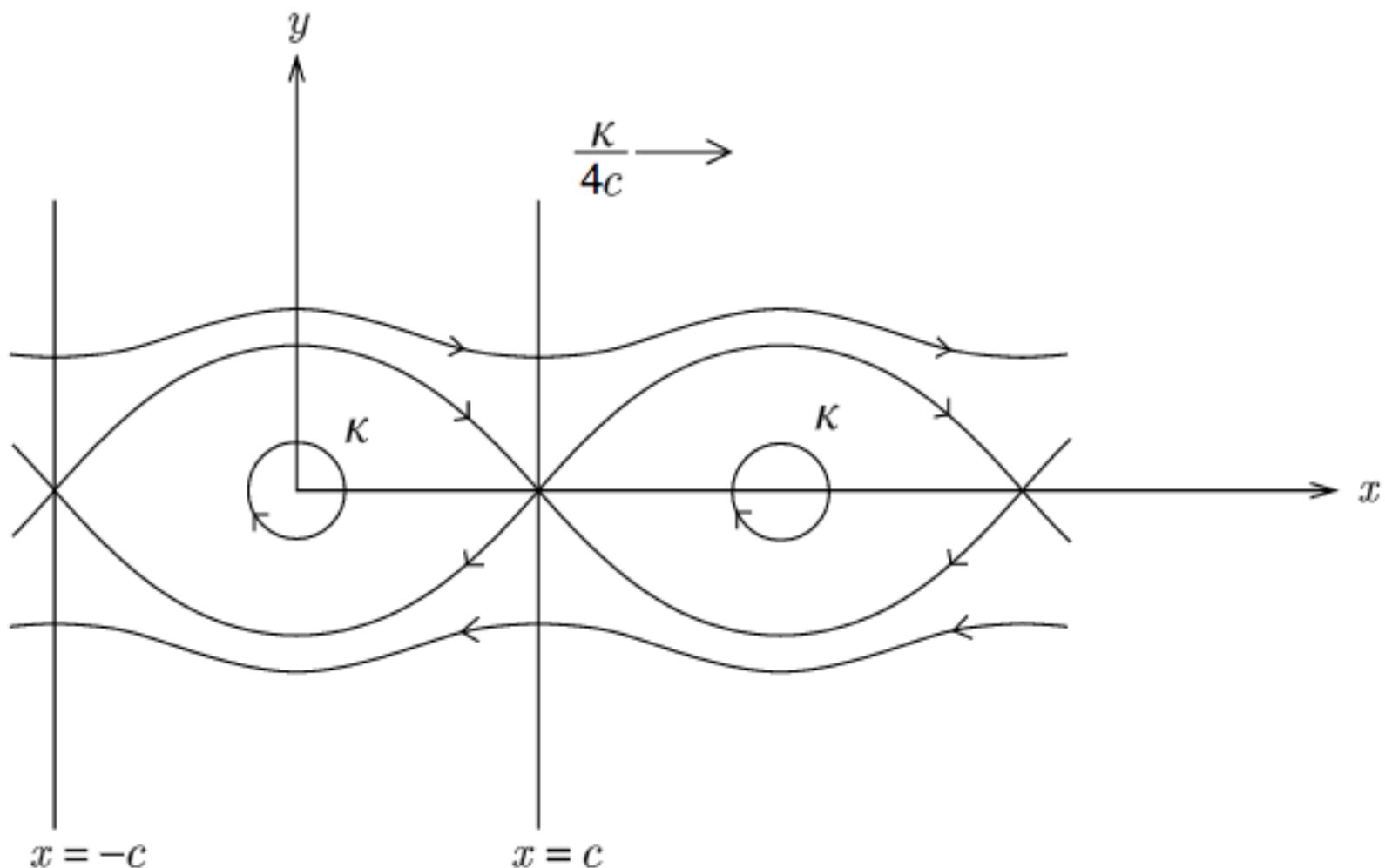


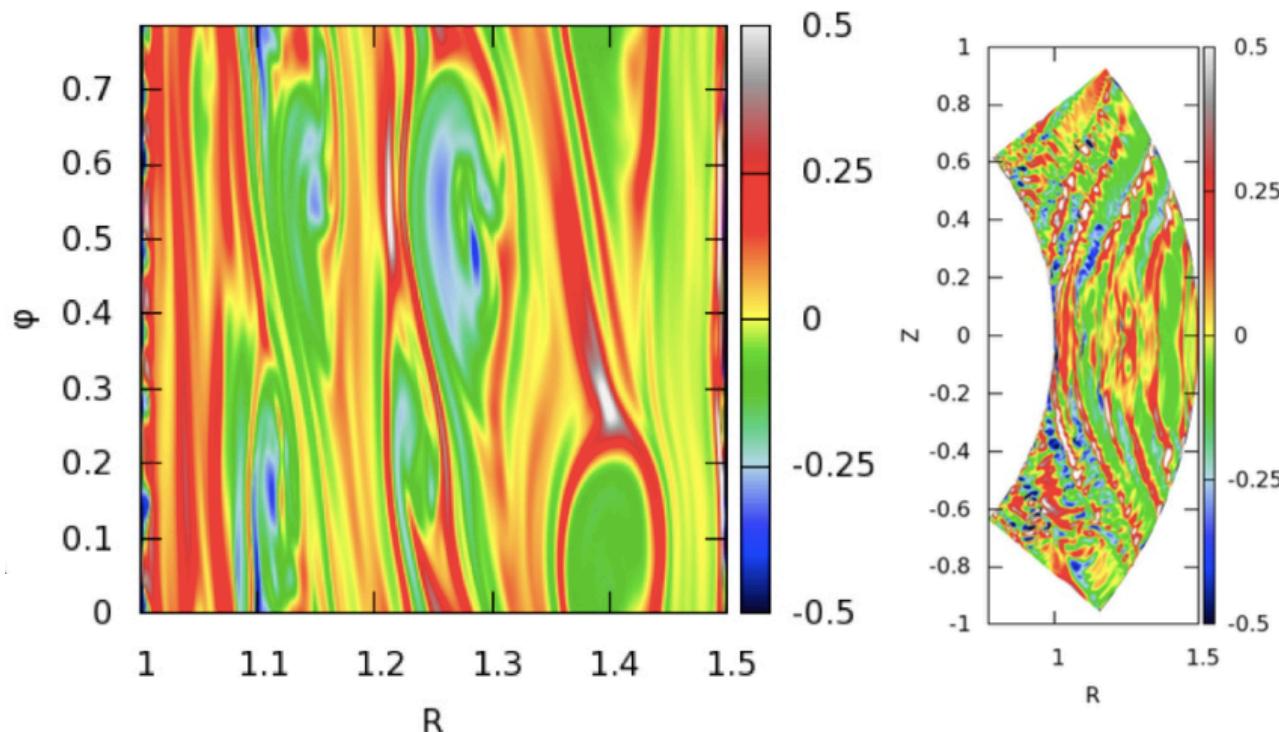
Figure 1. Notation and cat's-eye streamlines.

# Vertical shear instability

$$\rho_{\text{mid}} = \rho_0 \left( \frac{R}{R_0} \right)^p, \quad \Omega = \Omega_K \left[ 1 + \frac{1}{2} \left( \frac{H}{R} \right)^2 \left( p + q + \frac{q}{2} \frac{Z^2}{H^2} \right) \right]$$

$$c_s^2 = c_0^2 \left( \frac{R}{R_0} \right)^q,$$

$d\Omega / dz \neq 0 ; \kappa_z^2 < 0 \Rightarrow \text{Rayleigh unstable}$



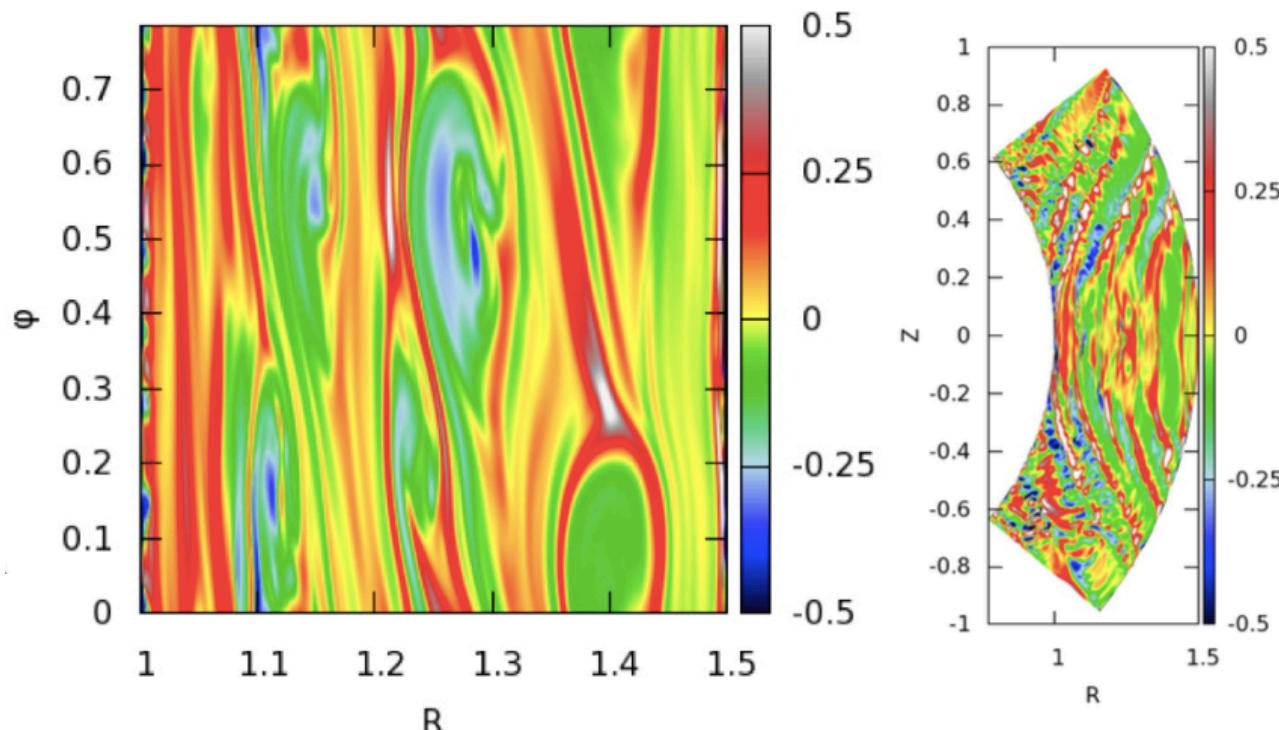
# Vertical shear instability

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$d\Omega / dz \neq 0 ; \kappa_z^2 < 0 \Rightarrow$  Rayleigh unstable

Solberg-Hoiland stability criterion  
 $\kappa^2 + N^2 > 0$



# Vertical shear instability

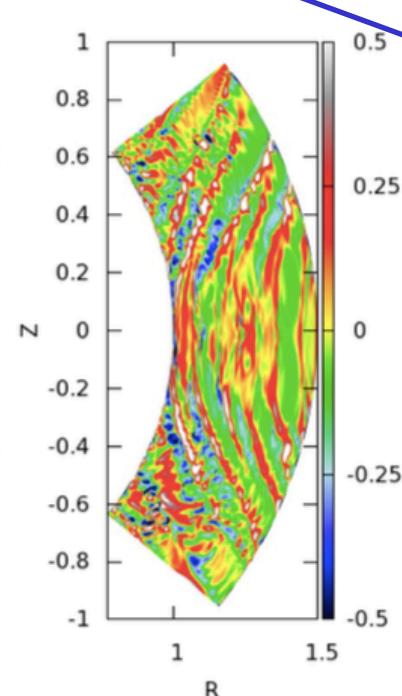
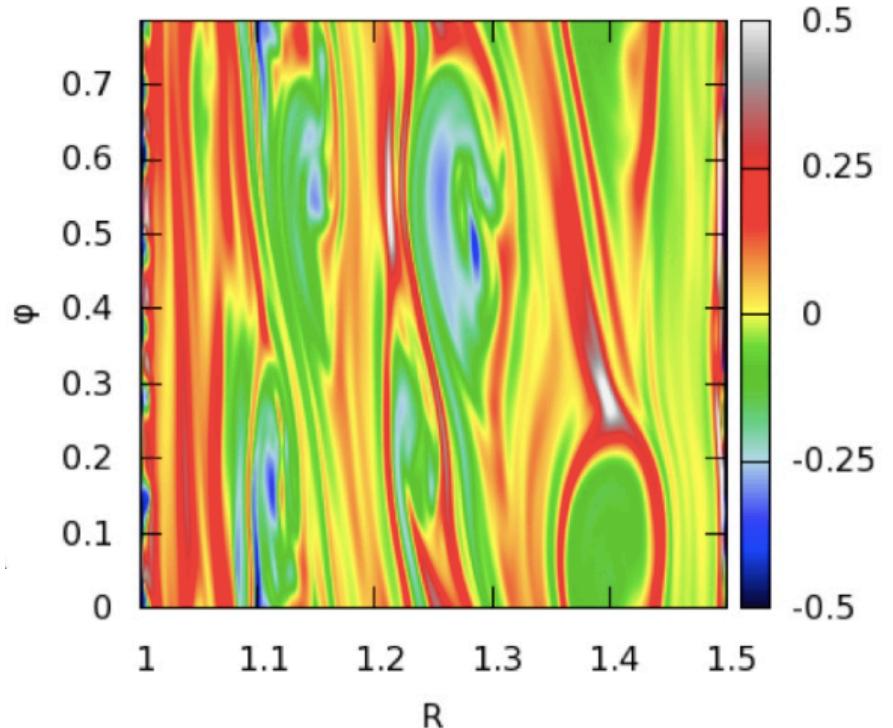
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$$c_s^2 = c_0^2 \left( \frac{R}{R_0} \right)^q,$$

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Solberg-Hoiland stability criterion

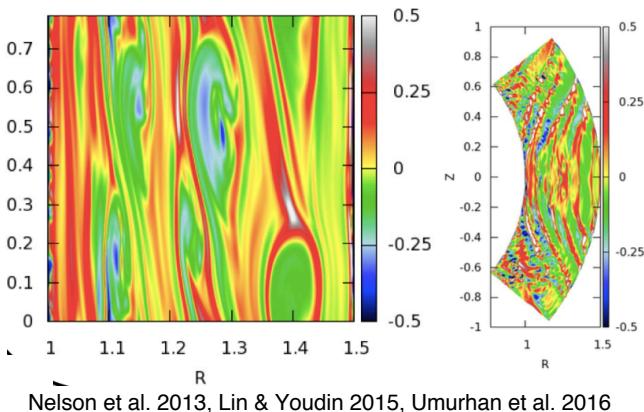
$$\kappa^2 + N^2 > 0$$



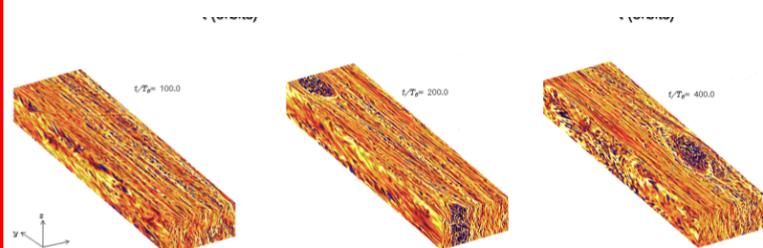
Buoyancy stabilizes!

# Thermal Instabilities

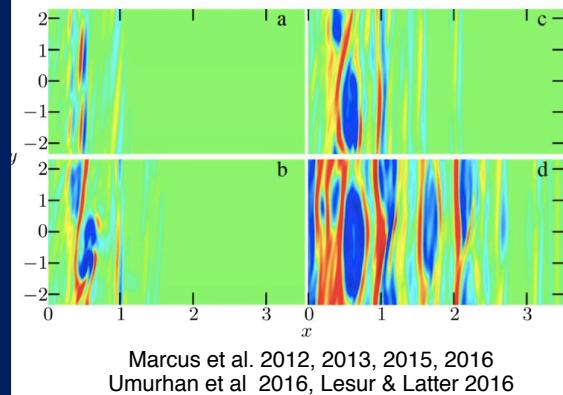
## Vertical Shear Instability



## Convective Overstability



## Zombie Vortex Instability



$$\alpha \sim 10^{-4} - 10^{-3}$$

$\Omega\tau \ll 1$   
( $\kappa < 1 \text{ cm}^2/\text{g}$ )

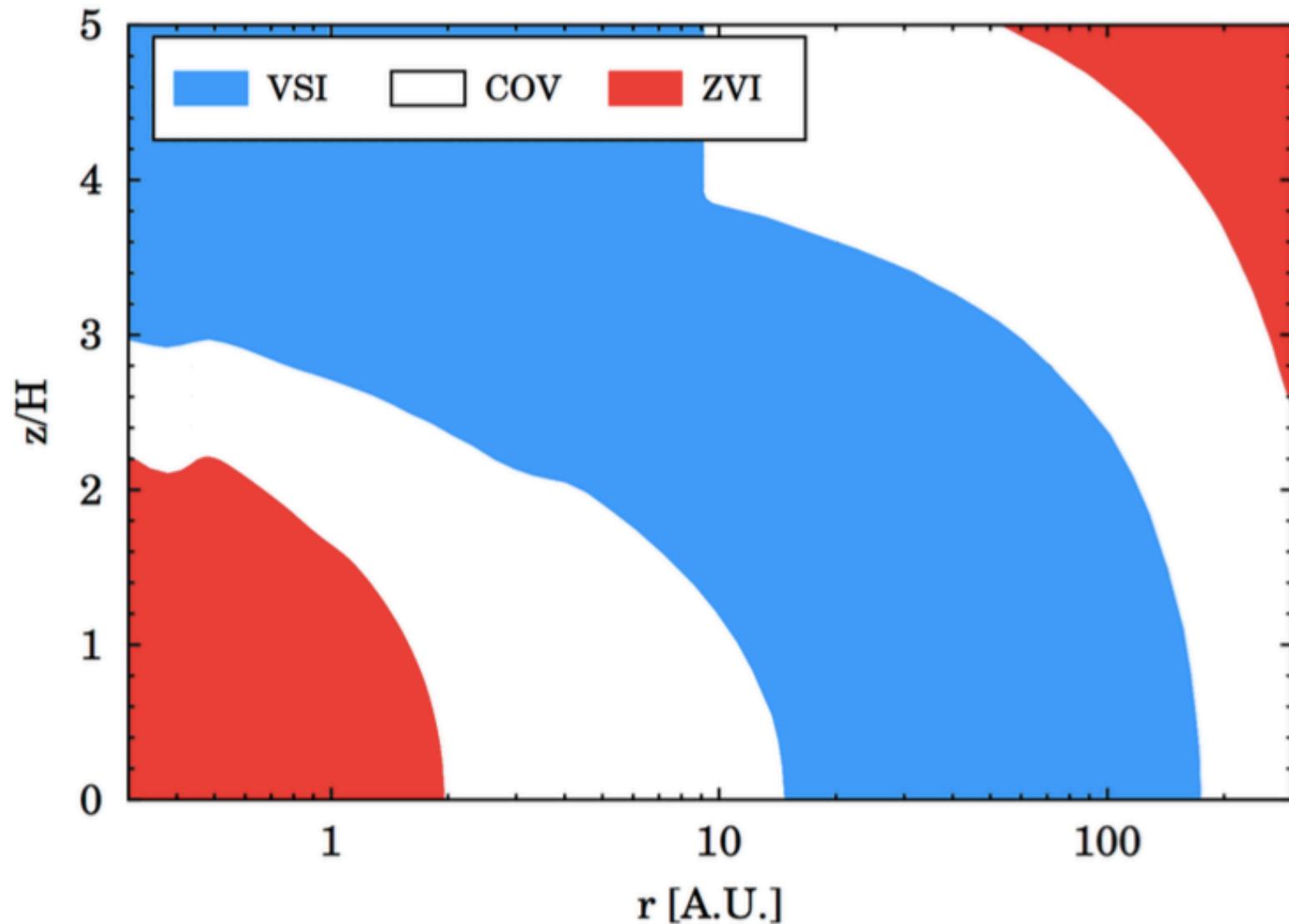
$$\alpha \sim 10^{-4} - 10^{-3}$$

$\Omega\tau \sim 1$   
( $\kappa \sim 1-50 \text{ cm}^2/\text{g}$ )

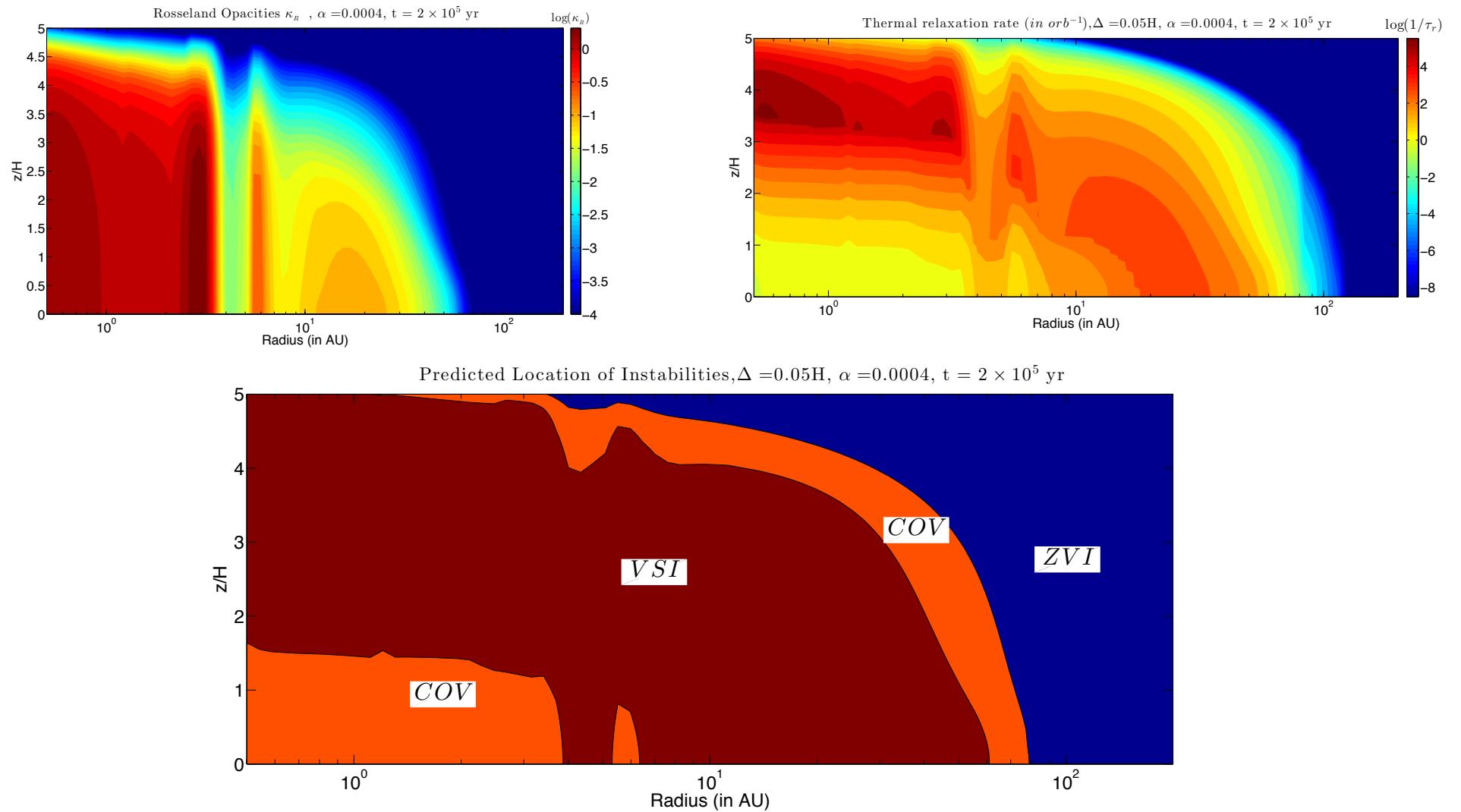
$$\alpha < \sim 10^{-3}$$

$\Omega\tau \gg 1$   
( $\kappa > 50 \text{ cm}^2/\text{g}$ )

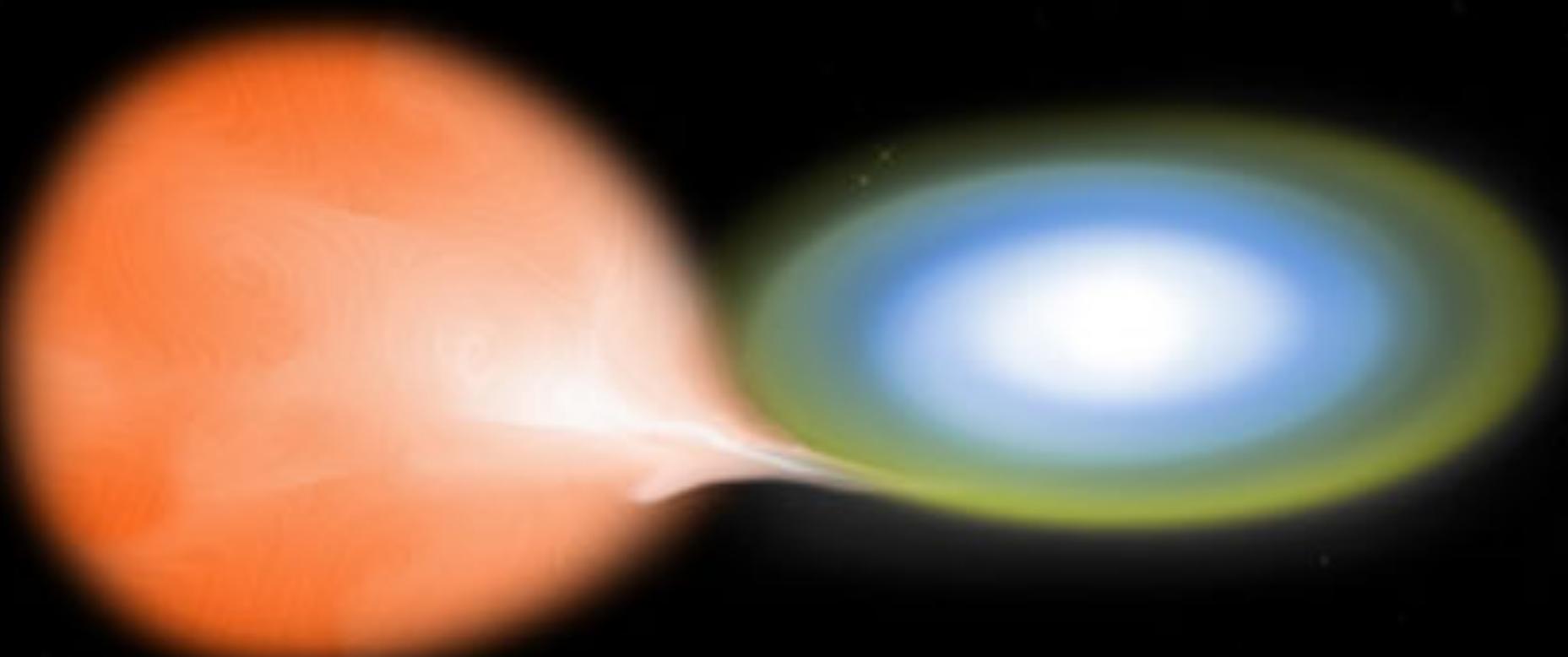
# Synthesis



# Synthesis

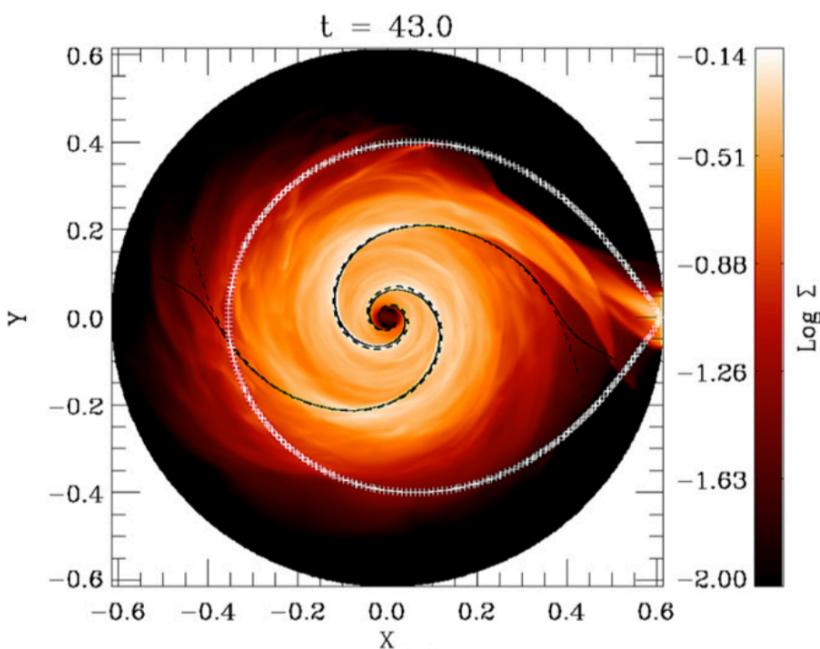


Cuzzi, Estrada, & Umurhan, in prep



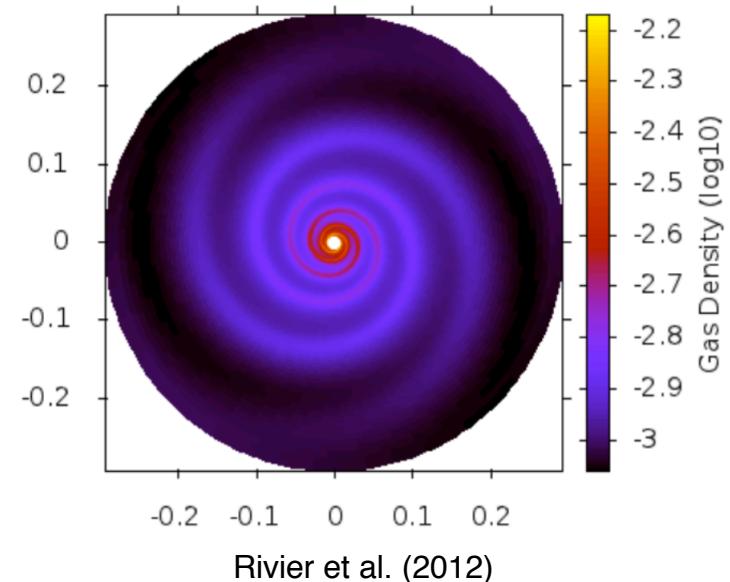
# Tidal deformation

Dwarf novae

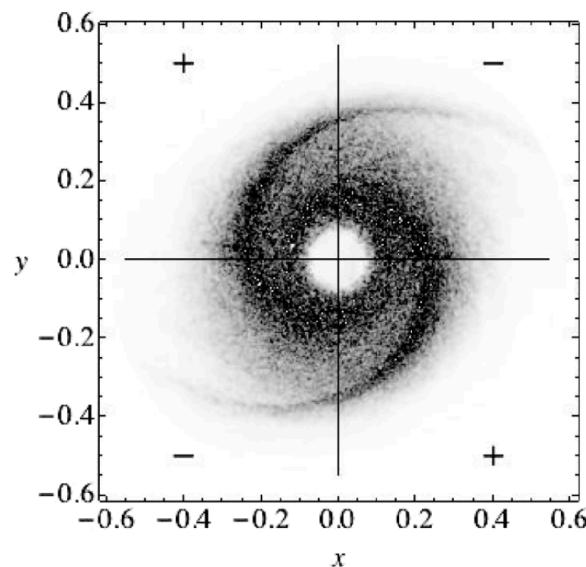


Ju, Stone and Zhu (2016)

Circumplanetary

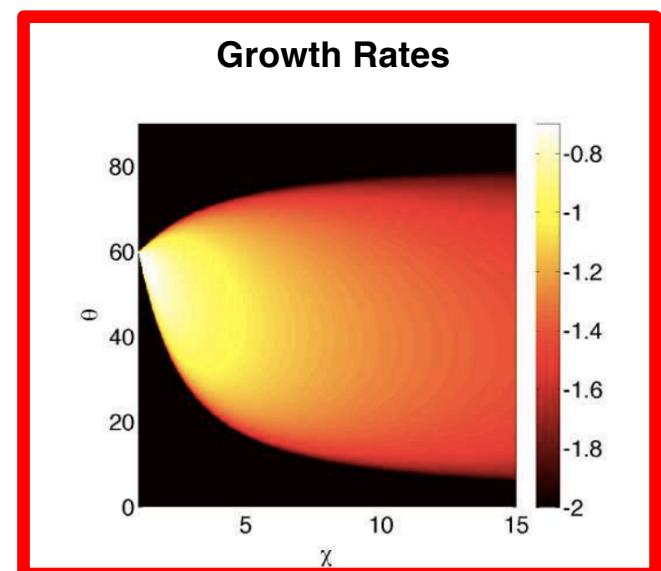
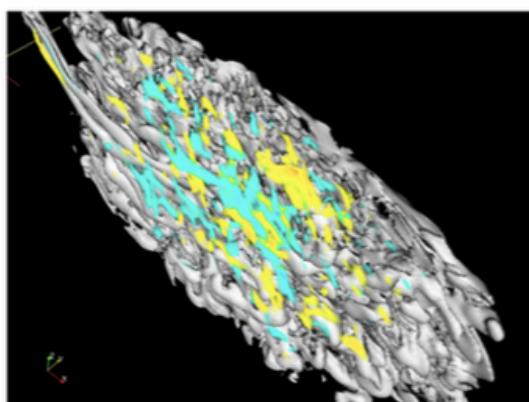
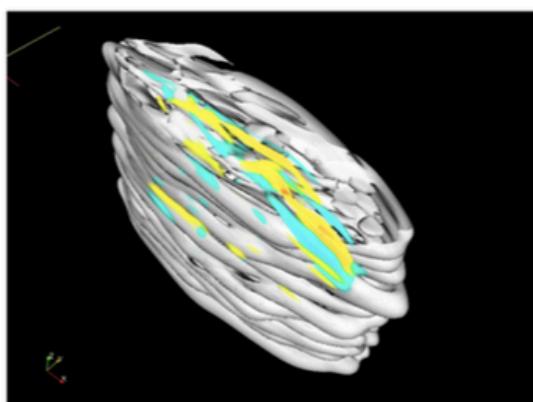
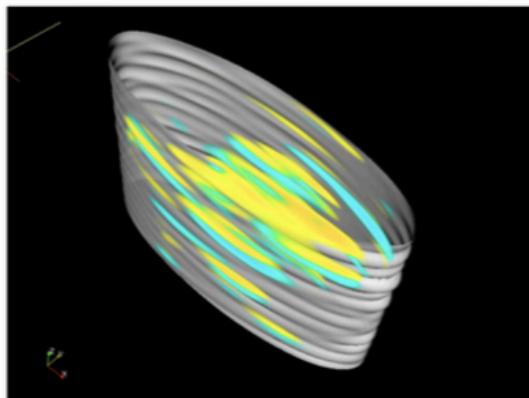
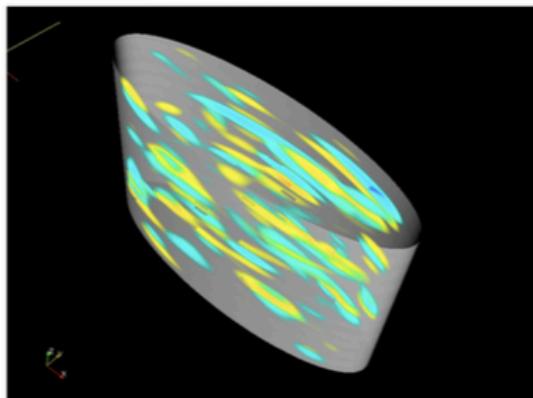


Rivier et al. (2012)



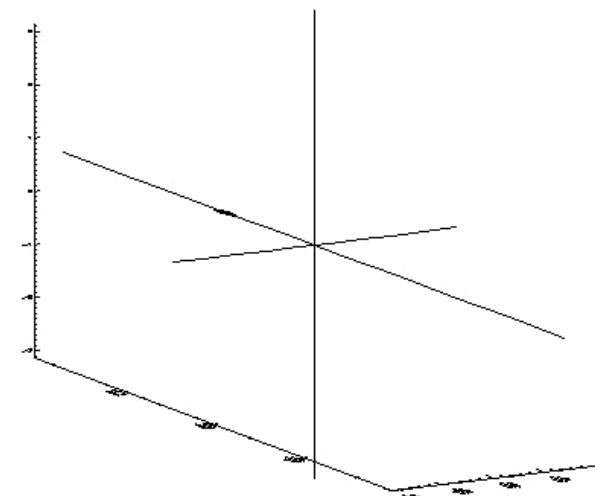
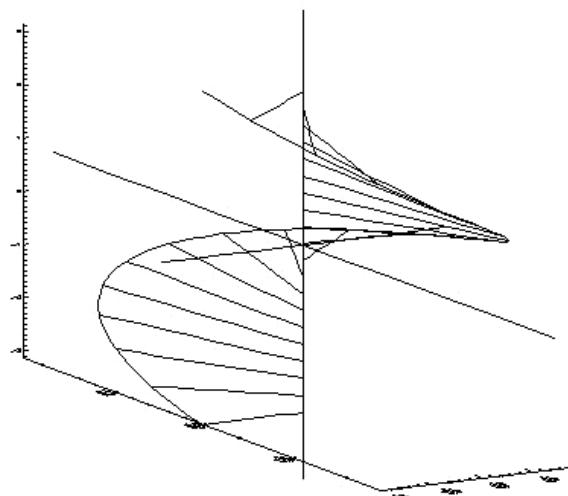
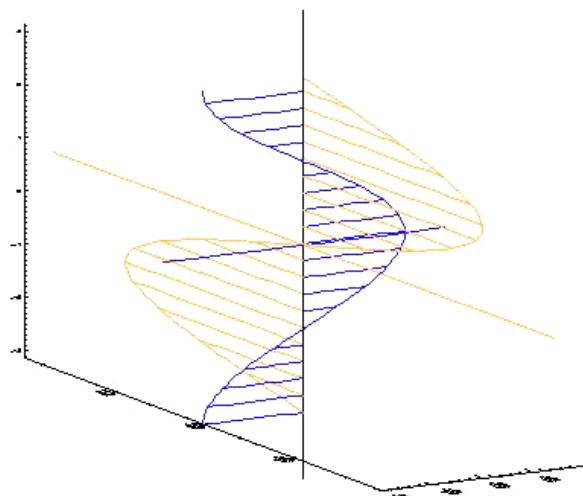
Martin & Lubow (2011)

## Elliptic Instability

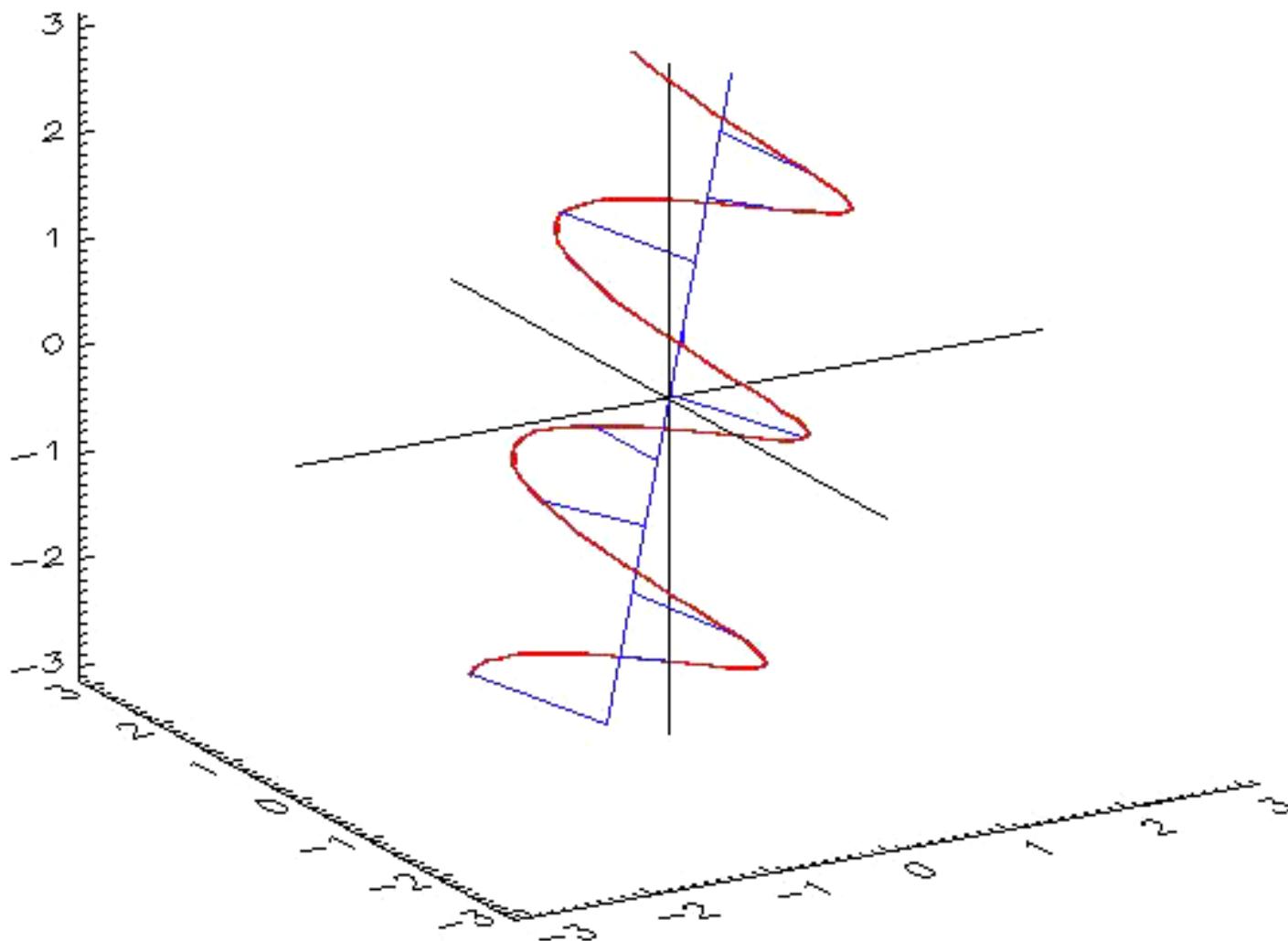


Bayly 1984, Lesur & Papaloizou 2009

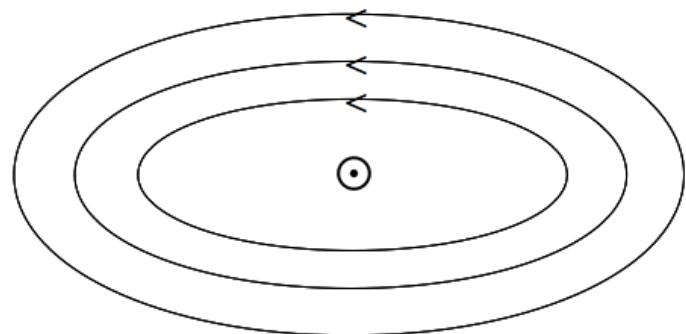
# Fluid in rigid rotation supports a spectrum of oscillations



## Fluid in rigid rotation supports a spectrum of oscillations

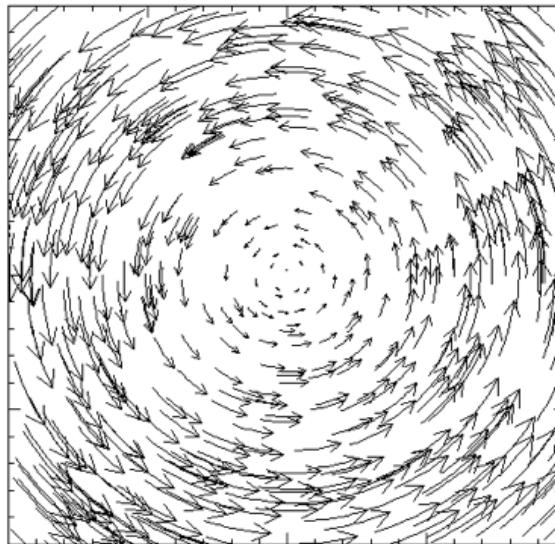


# Introducing ellipticity: Strain



$$U = [-(1-\epsilon)y, (1-\epsilon)x]$$

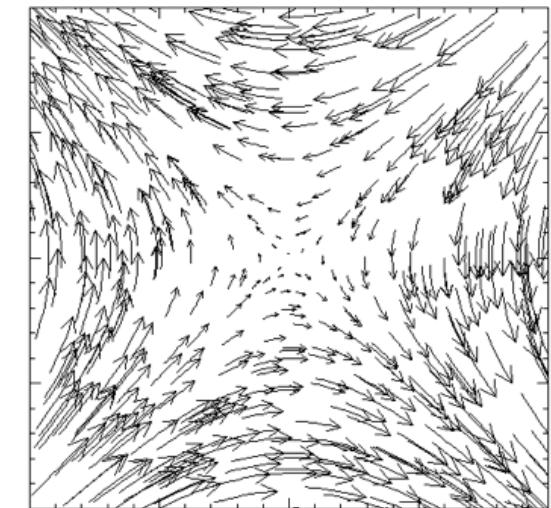
=



$$[-y, x]$$

Rigid rotation

+

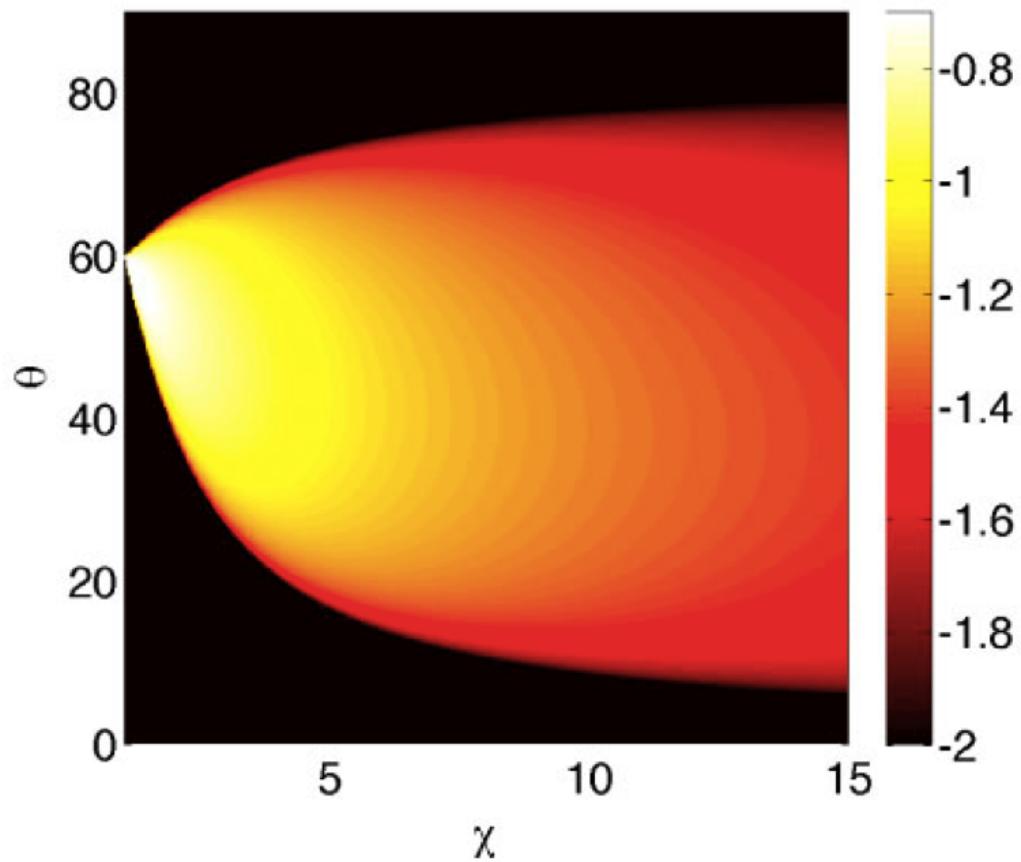


$$-\epsilon [y, x]$$

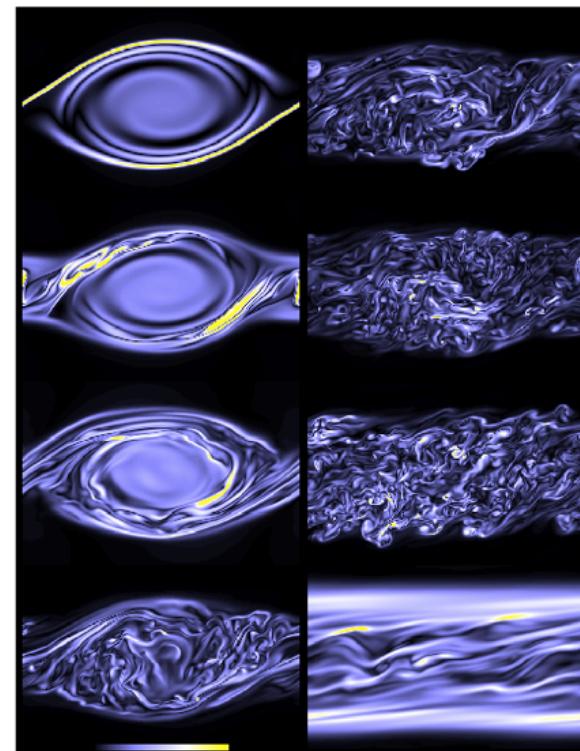
Strain field

Rigid rotation is stable.  
Strain is **not** necessarily so.

# Elliptic Instability



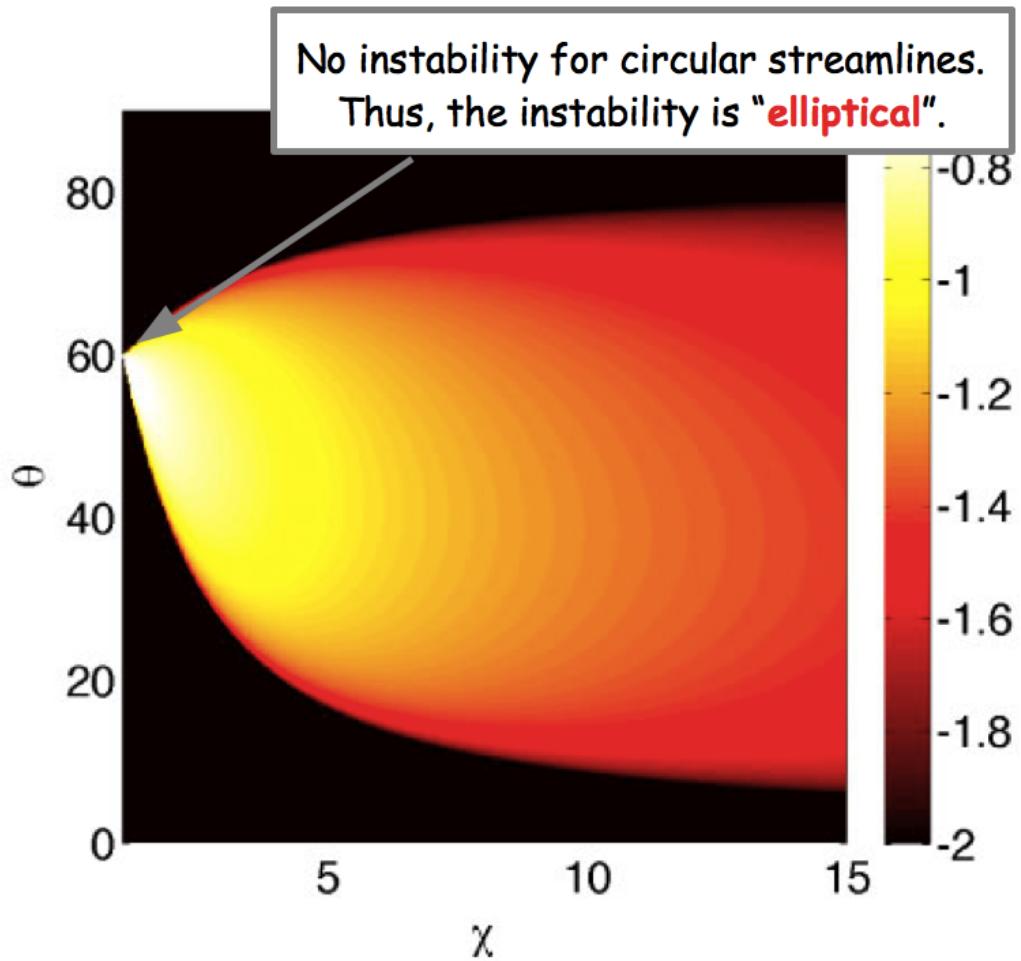
Lesur & Papaloizou (2009)  
After Bayly (1986)



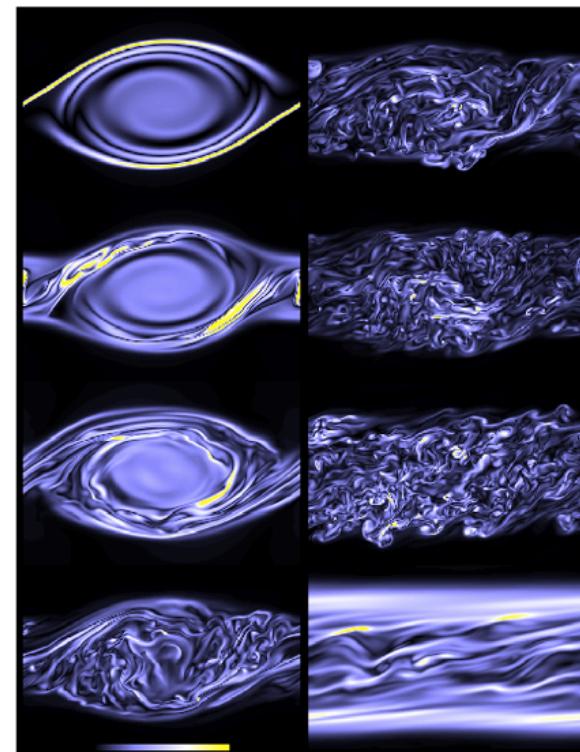
Vortex coherence is destroyed.  
Energy cascades forward and dissipates.  
The flow relaminarizes.

McWilliams (2010)

# Elliptic Instability

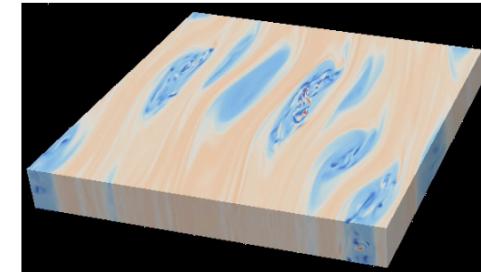


Lesur & Papaloizou (2009)  
After Bayly (1986)



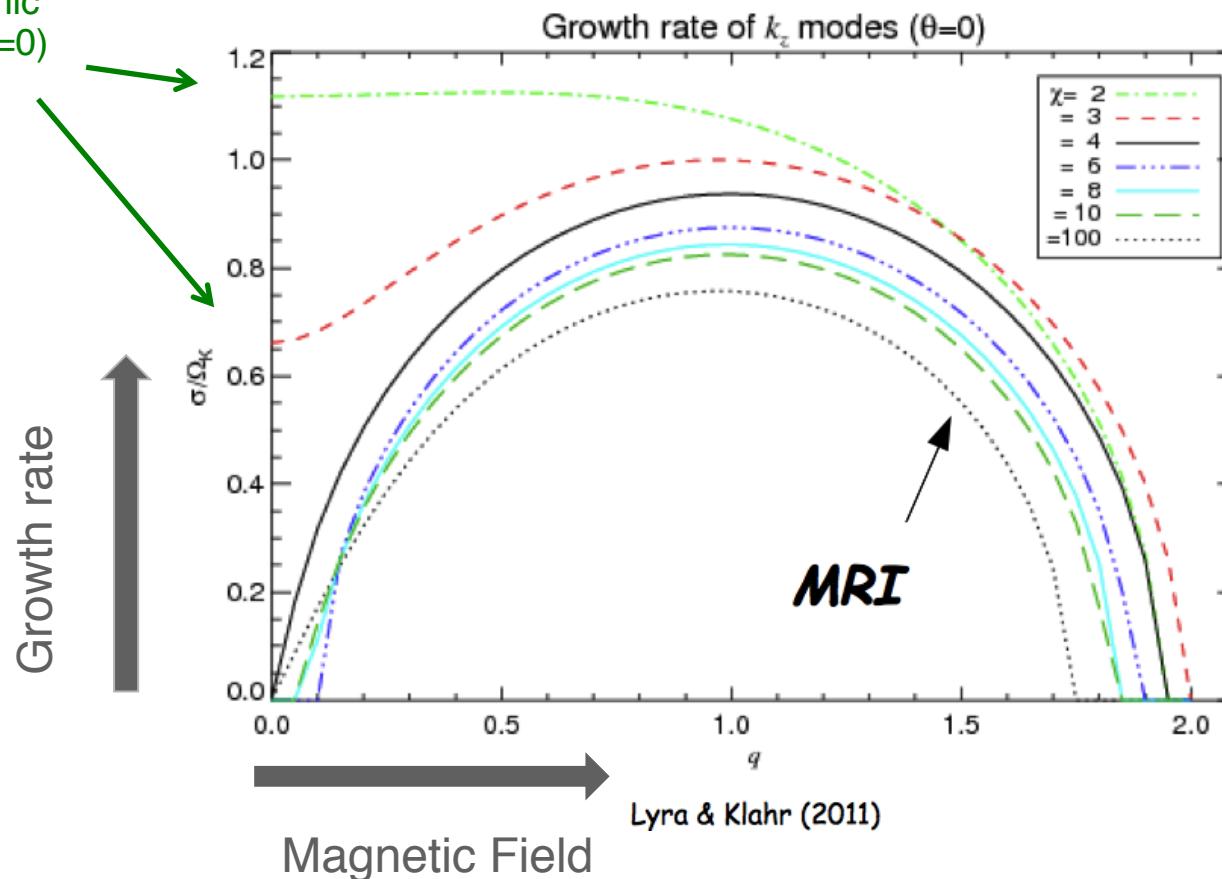
McWilliams (2010)

# Magneto-Elliptic Instability



Lesur & Papaloizou (2010)

Hydrodynamic  
instability ( $B=0$ )



See also

Pierrehumbert 1986

Bayly 1986

Kerswell 2002

Lesur & Papaloizou 2009

Lesur & Papaloizou 2010

Lyra & Klahr 2011

Lyra 2013

**Infinitely elongated vortices** are equivalent to **shear flows**.

They are subject to an MRI-like instability when magnetized.

# Conclusion

