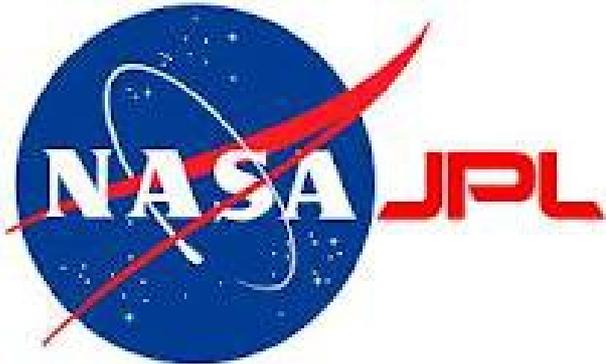


Sharp eccentric rings in debris disks

Who needs a planet?

Wladimir (Wlad) Lyra

Carl Sagan Fellow



NASA-JPL/Caltech

American Museum of Natural History

Max-Planck Institute for Astronomy

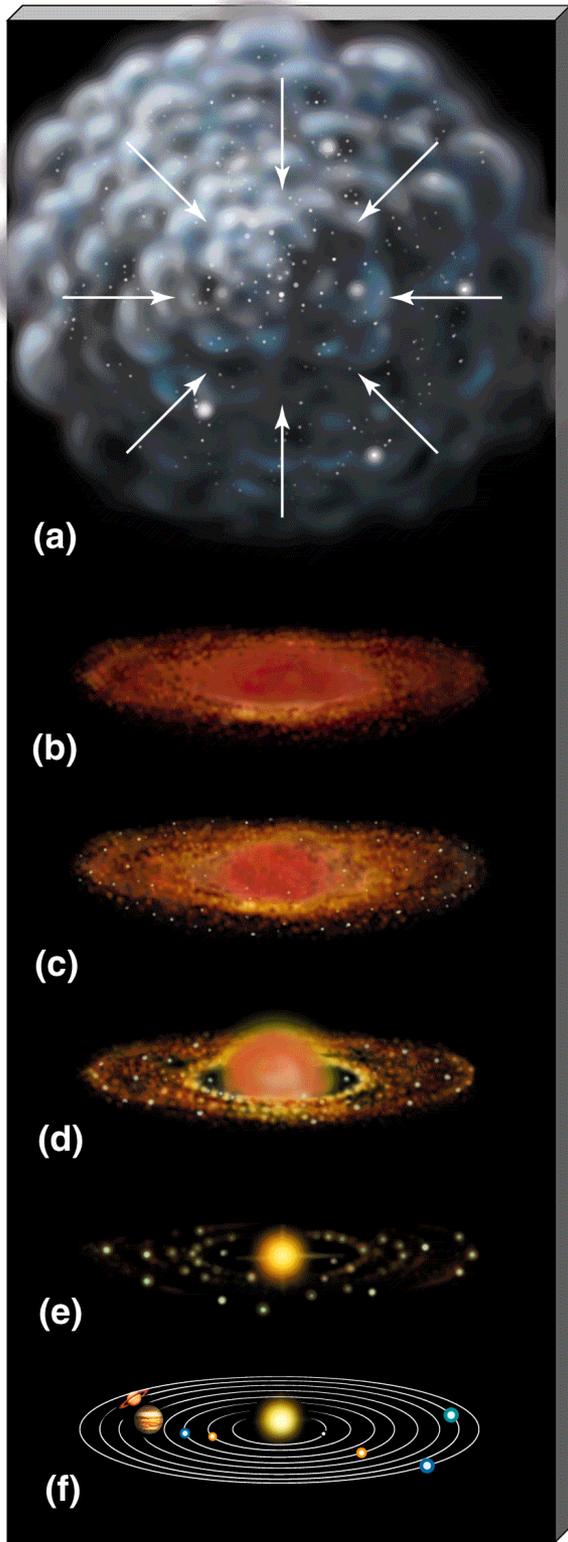
University of Uppsala



Montréal QC, September 2012

Collaborators:

Axel Brandenburg (Stockholm), Kees Dullemond (Heidelberg), Anders Johansen (Lund),
Brandon Horn (Columbia), Hubert Klahr (Heidelberg), Marc Kuchner (Goddard),
Mordecai-Mark Mac Low (AMNH), Sijme-Jan Paardekooper (Cambridge),
Nikolai Piskunov (Uppsala), Natalie Raettig (Heidelberg), Zsolt Sandor (Innsbruck),
Neal Turner (JPL), Andras Zsom (MIT).



A disk life story

Gas-rich phase (< 10 Myr)

T-Tauri Disks

Accretion and Planet Formation

Thinning phase (~10 Myr)

Transitional Disks

Planet retention

Gas-poor phase (>10 Myr)

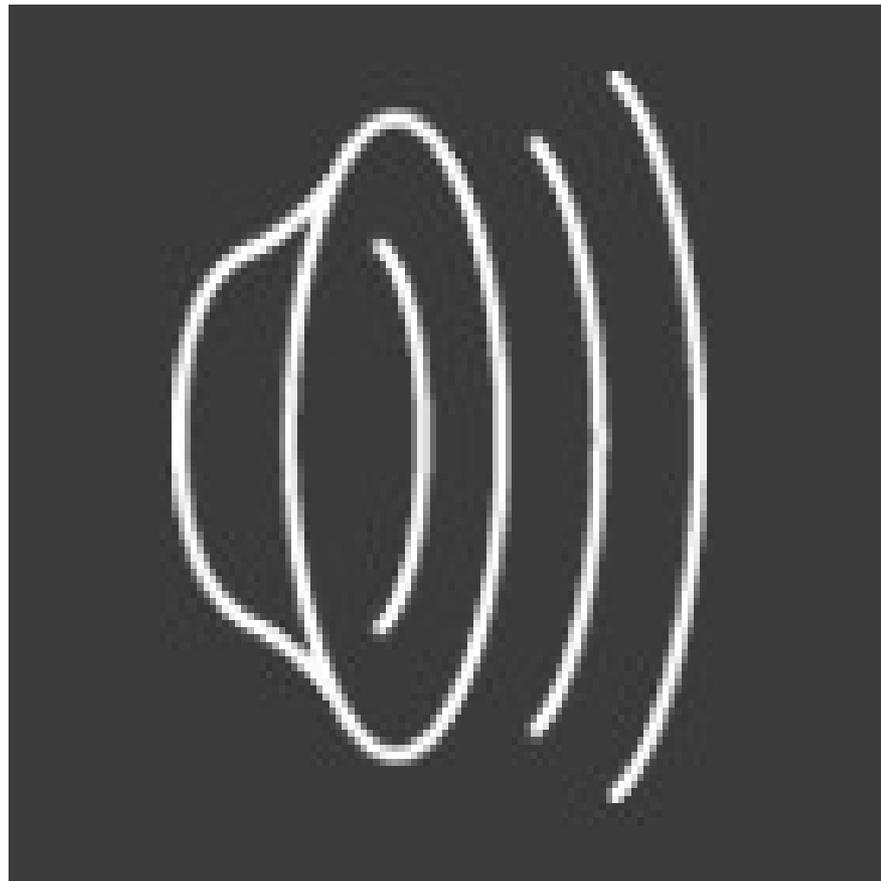
Debris Disks

Stabilization of architecture and Planet Detection

Accretion in disks occurs via turbulent viscosity

Turbulence in disks is enabled by the **Magneto-Rotational Instability**

MRI sketch



Magnetized disk

Lyra et al. (2008a)

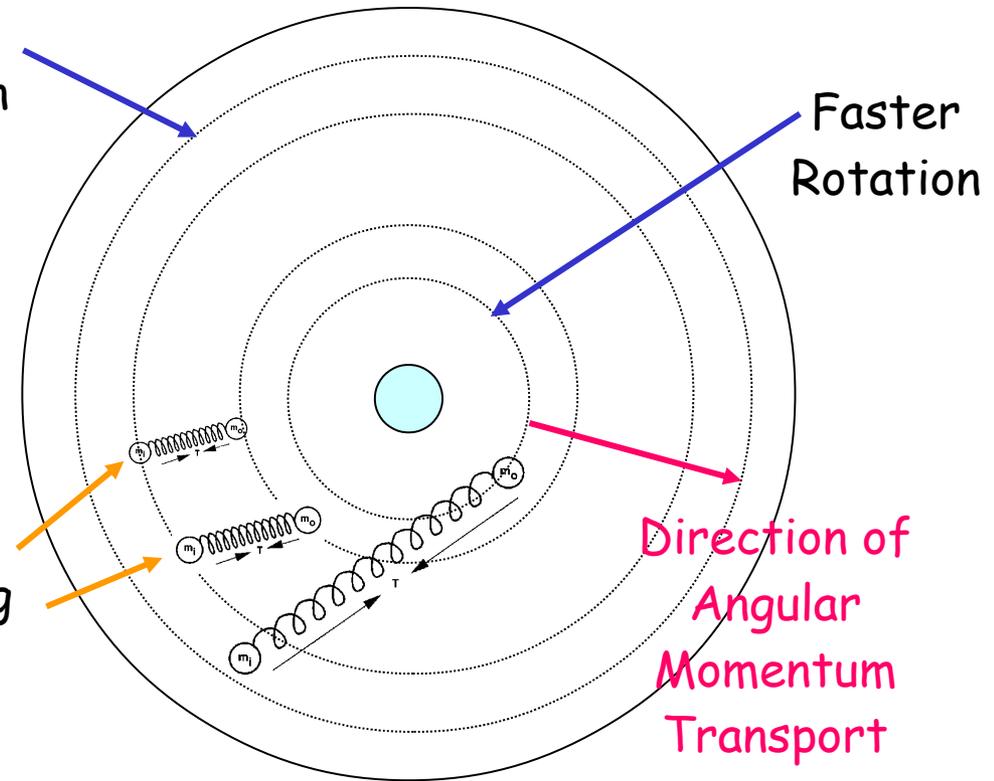
Slower
rotation

Faster
Rotation

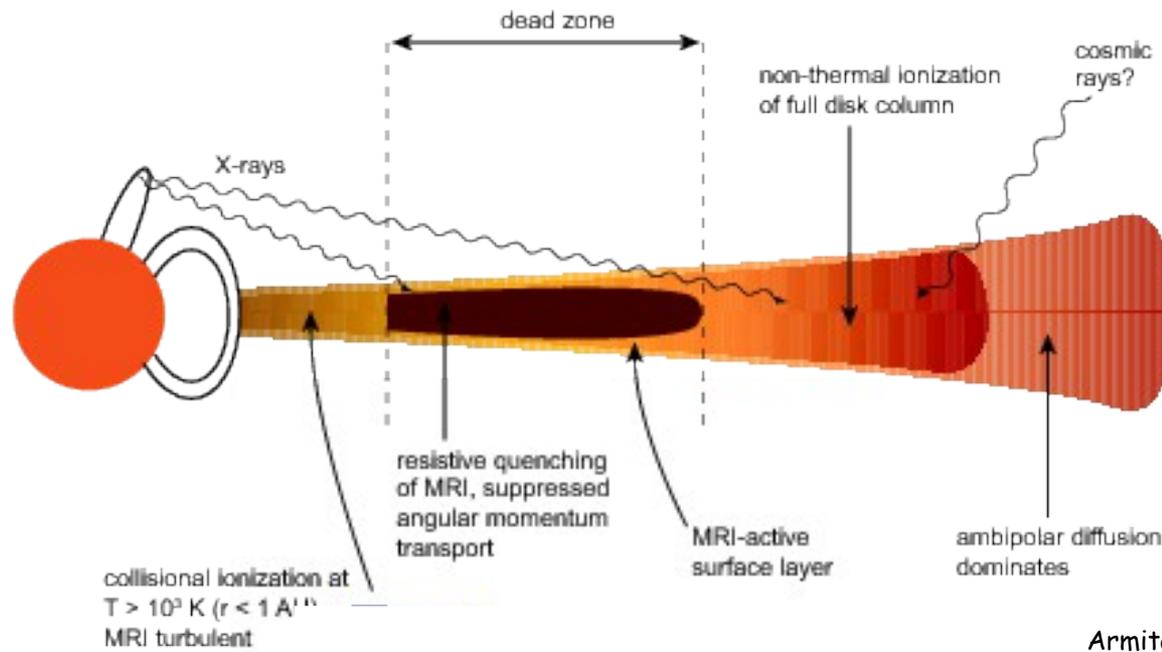
etching
plifies
field

Direction of
Angular
Momentum
Transport

Unstable if angular
velocity decreases
outward



Alas... Dead zones are robust features of accretion disks



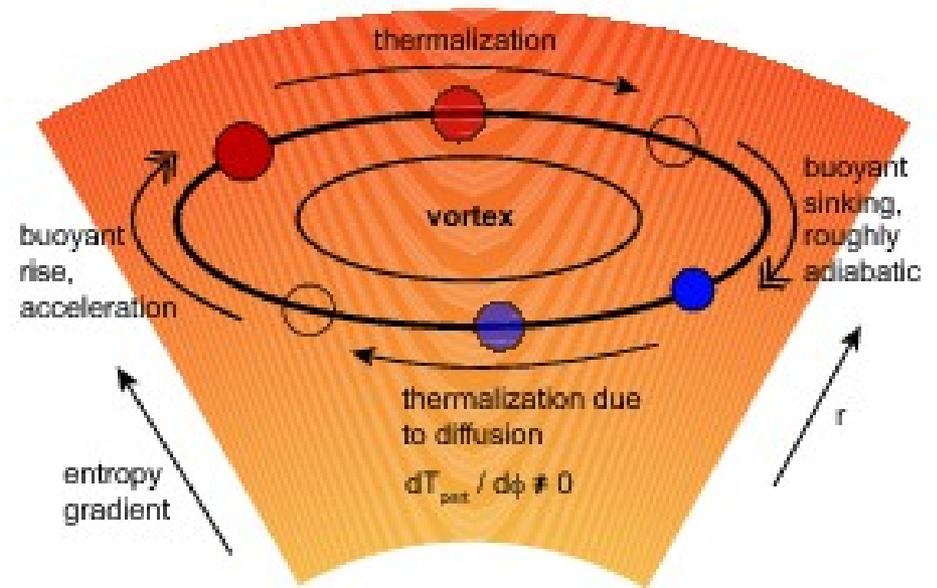
Disks are cold and thus poorly ionized
(Blaes & Balbus 1994)

Therefore, accretion is **layered** (Gammie 1996)

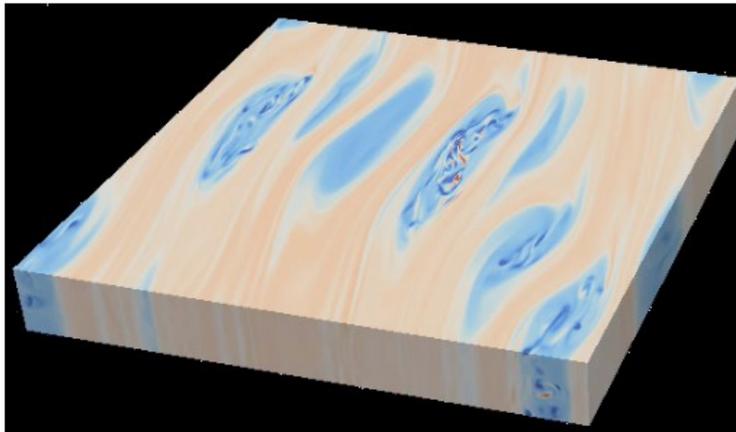
There should be a non-magnetic,
hydrodynamical, source of turbulence in the **dead zone**.

Baroclinic Instability - Excitation and self-sustenance of vortices

Sketch of the Baroclinic Instability



Armitage (2010)

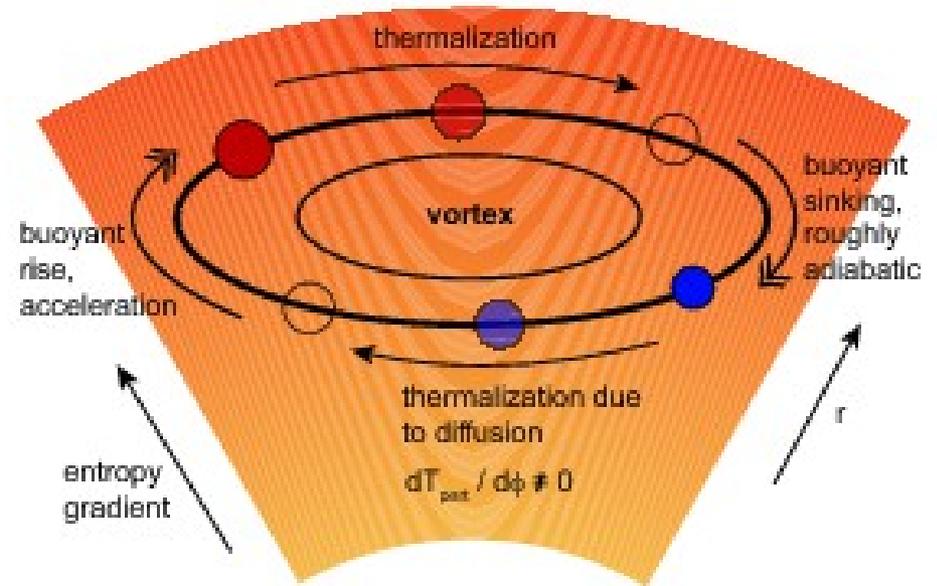


Lesur & Papaloizou (2010)

$$\frac{\partial \omega}{\partial t} = \underbrace{-(\mathbf{u} \cdot \nabla) \omega}_{\text{advection}} - \underbrace{\omega (\nabla \cdot \mathbf{u})}_{\text{compression}} + \underbrace{(\omega \cdot \nabla) \mathbf{u}}_{\text{stretching}} + \frac{1}{\rho^2} \underbrace{\nabla \rho \times \nabla p}_{\text{baroclinicity}} + \nu \underbrace{\nabla^2 \omega}_{\text{dissipation}}$$

Baroclinic Instability - Excitation and self-sustenance of vortices

Sketch of the Baroclinic Instability



Armitage (2010)

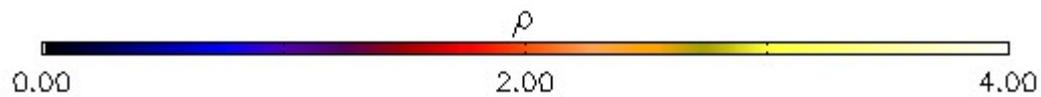
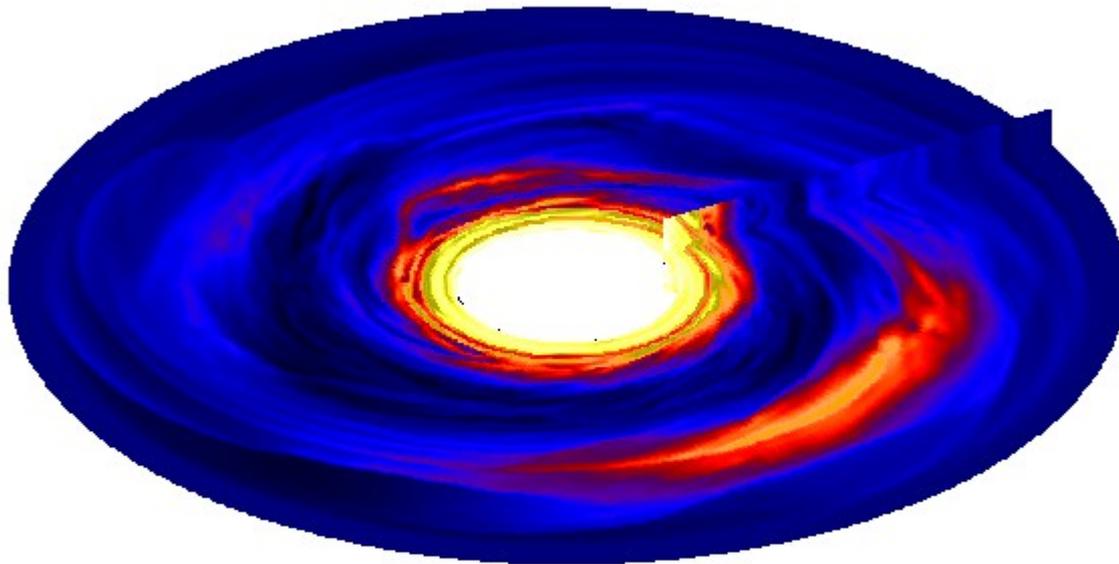


Lyra & Klahr (2011)

$$\frac{\partial \omega}{\partial t} = \underbrace{-(\mathbf{u} \cdot \nabla) \omega}_{\text{advection}} - \underbrace{\omega (\nabla \cdot \mathbf{u})}_{\text{compression}} + \underbrace{(\omega \cdot \nabla) \mathbf{u}}_{\text{stretching}} + \frac{1}{\rho^2} \underbrace{\nabla \rho \times \nabla p}_{\text{baroclinicity}} + \nu \underbrace{\nabla^2 \omega}_{\text{dissipation}}$$

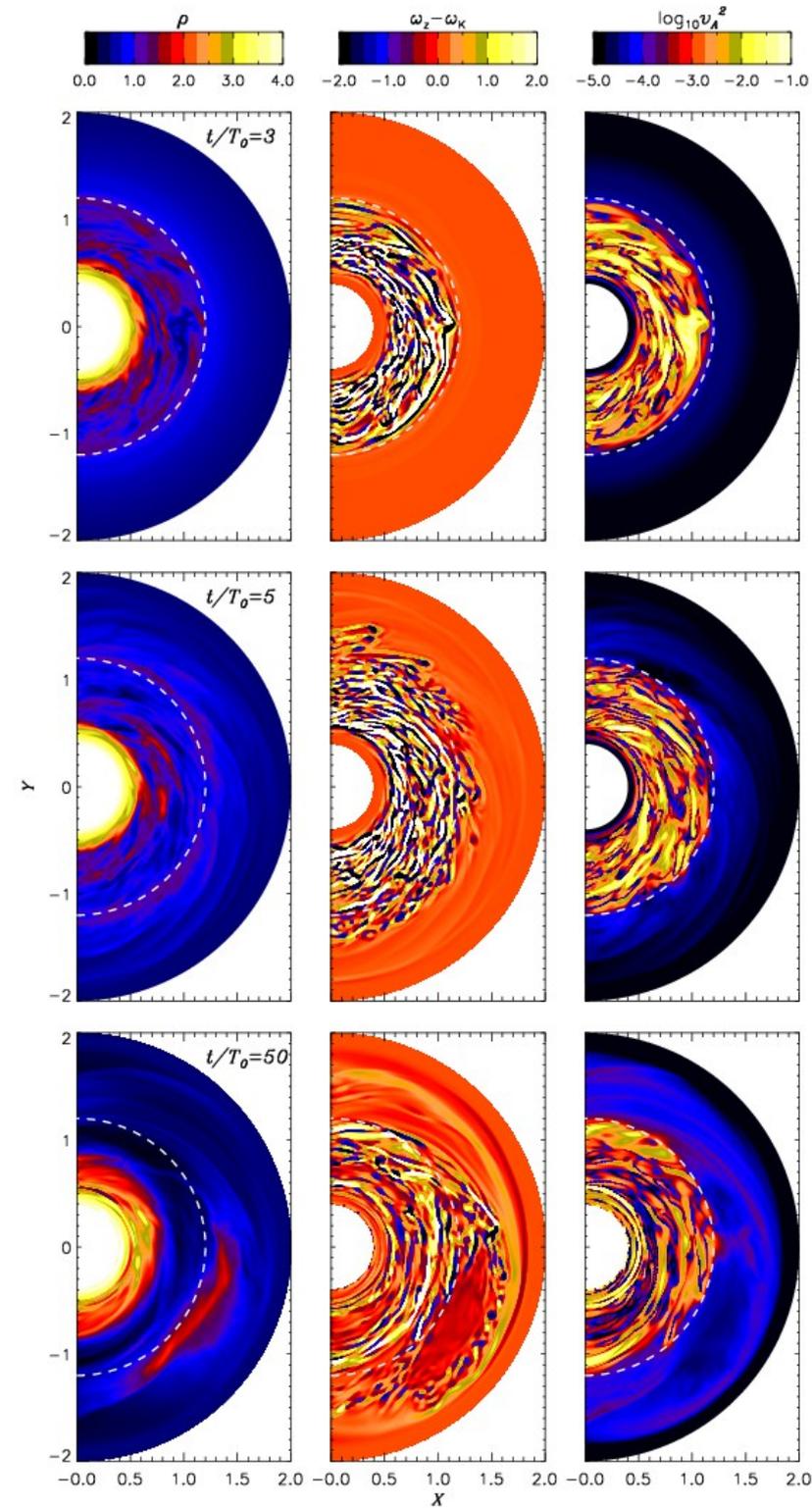
Active/dead zone boundary

$t = 22.28 \tau_0$



Magnetized inner disk + resistive outer disk

Lyra & Mac Low (2012)

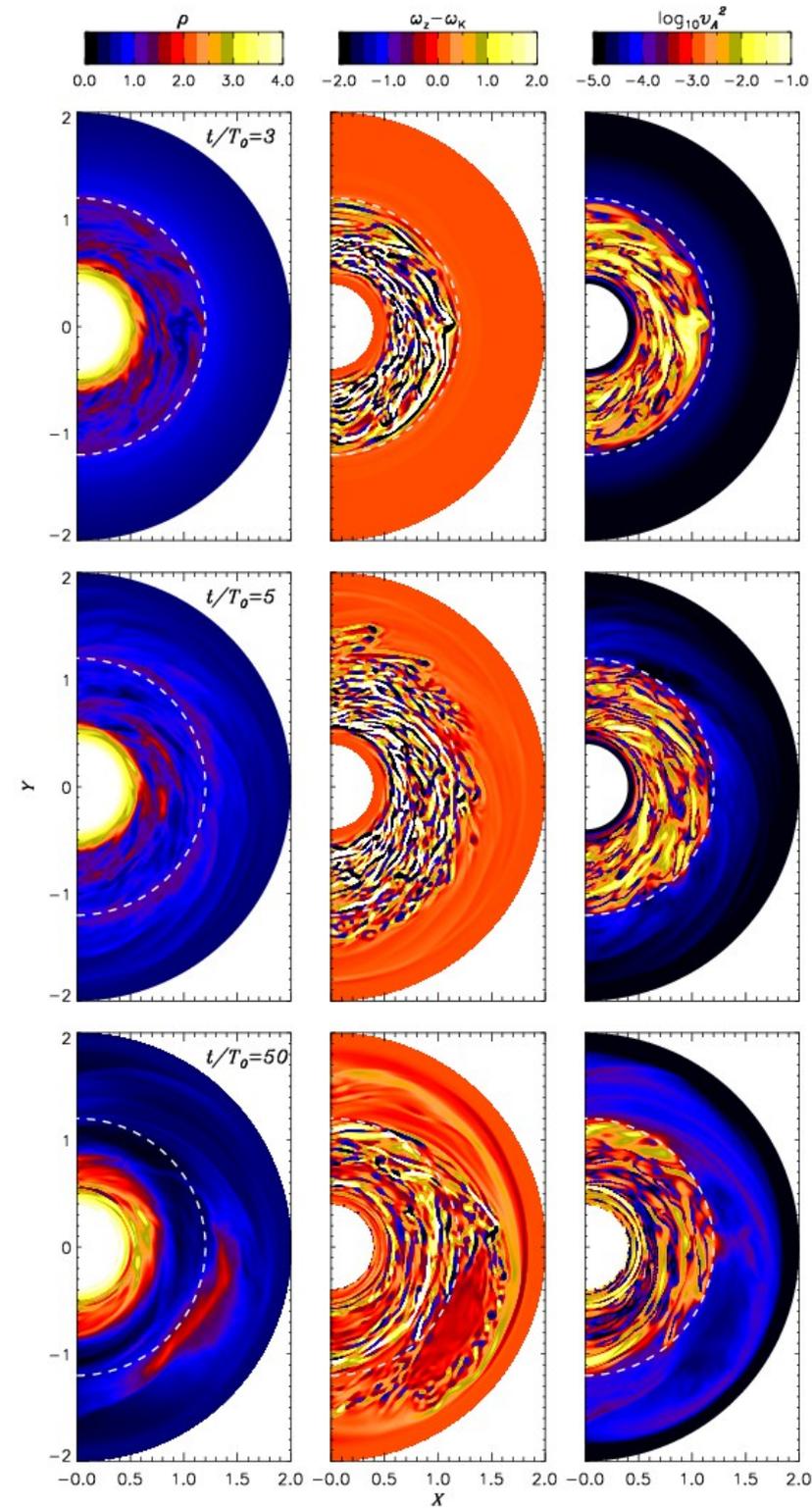


Active/dead zone boundary



Magnetized inner disk + resistive outer disk

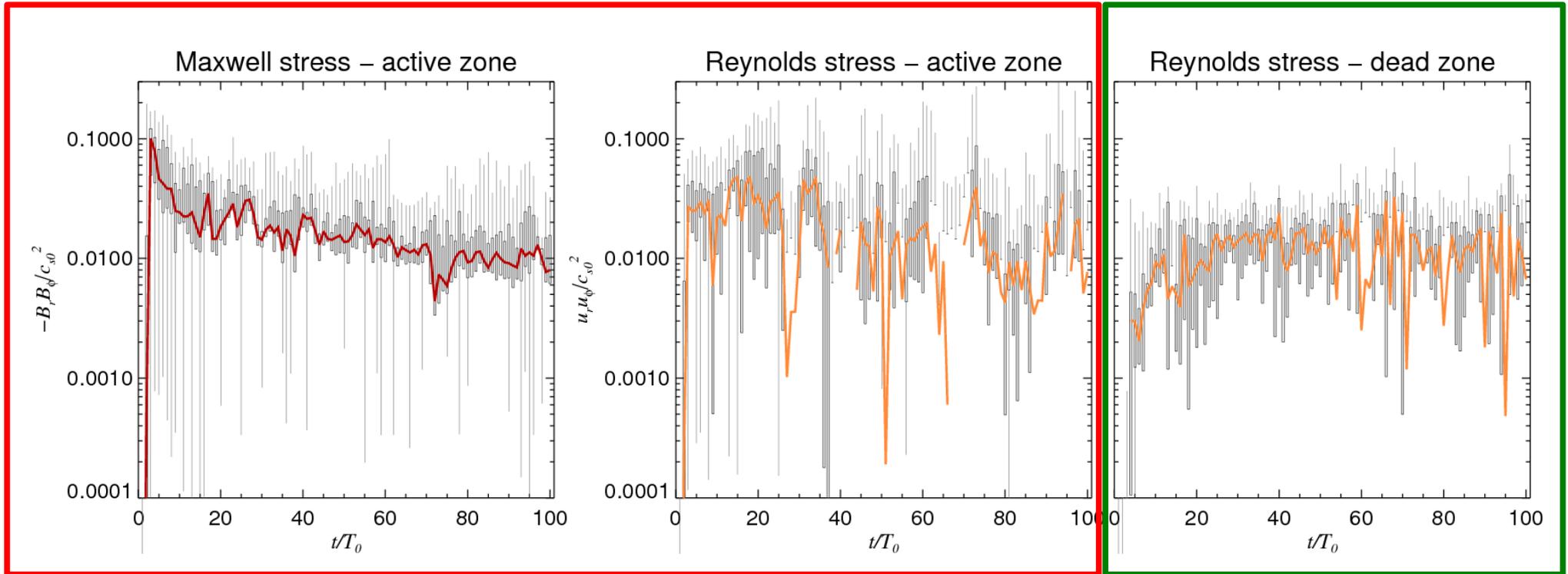
Lyra & Mac Low (2012)



Significant angular momentum transport

Active zone

Dead zone

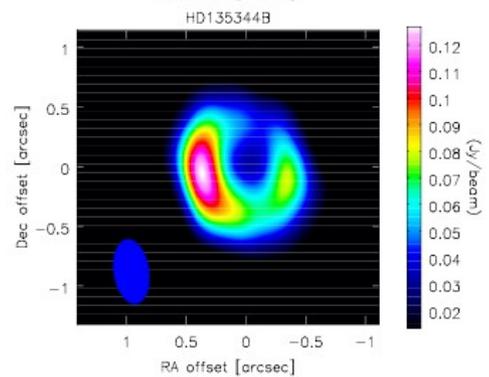
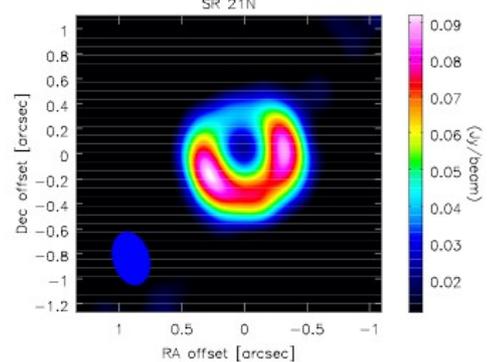
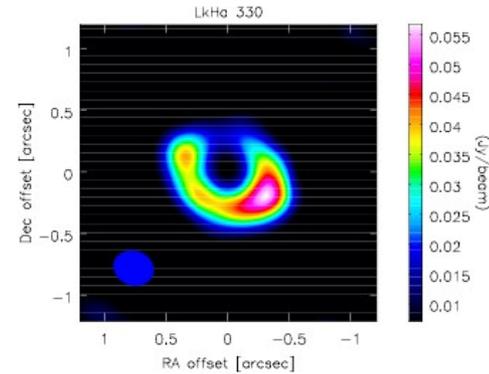
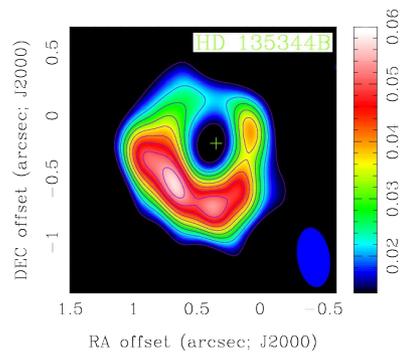
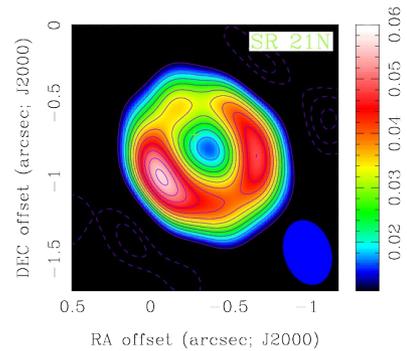
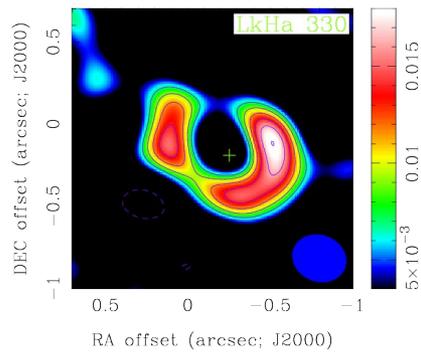


Large mass accretion rates in the **dead zone**,
comparable to the MRI in the **active zone**!

A possible detection of vortices in disks

Observations

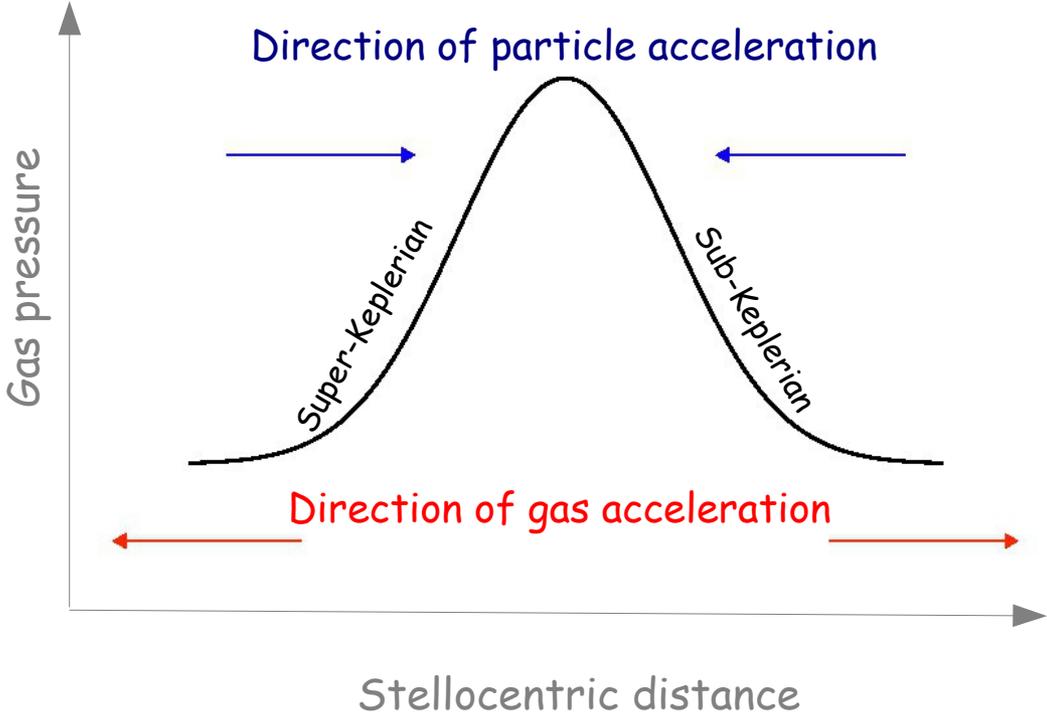
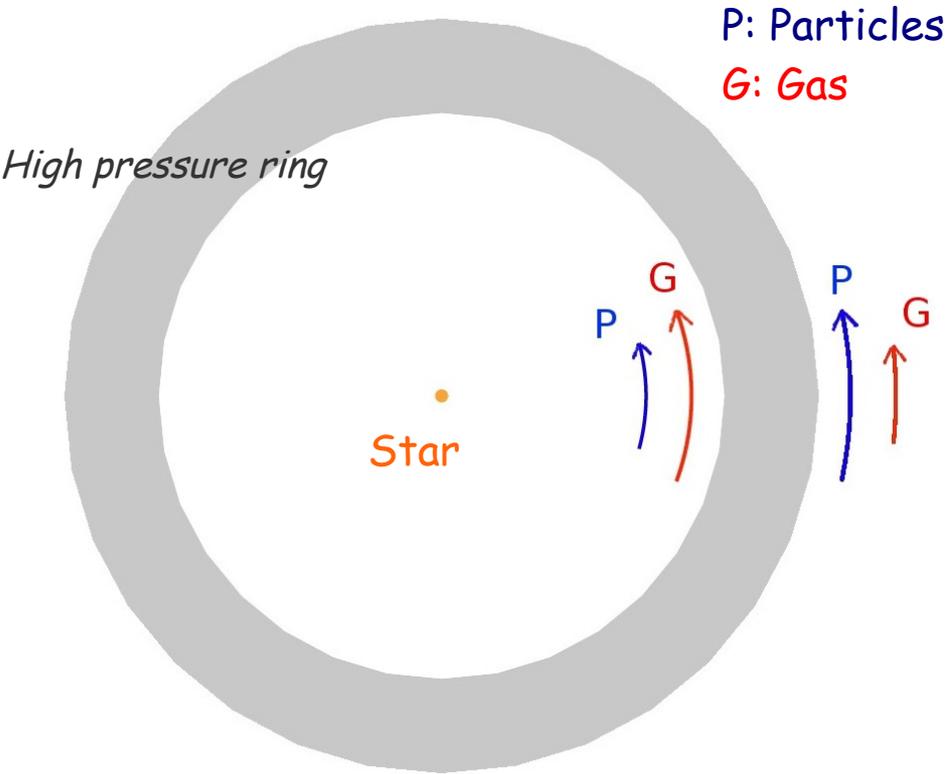
Brown et al. (2009)



Simulated observations
of Rossby vortices

Regaly et al. (2012)

Forming planets in turbulent disks



Adapted from Whipple (1972)

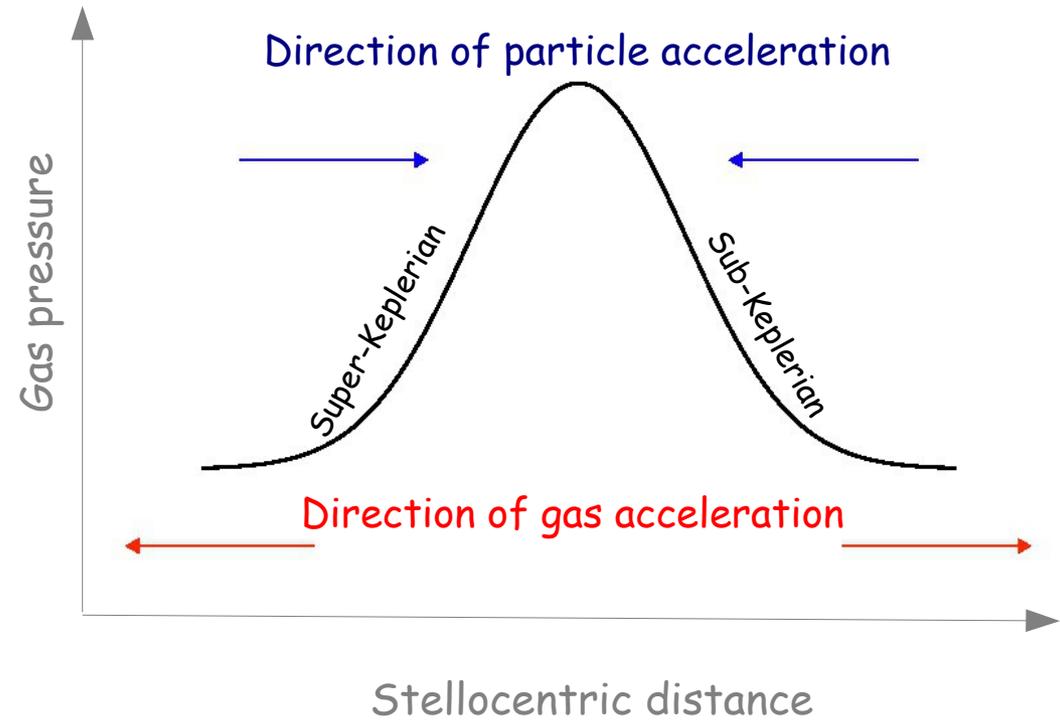
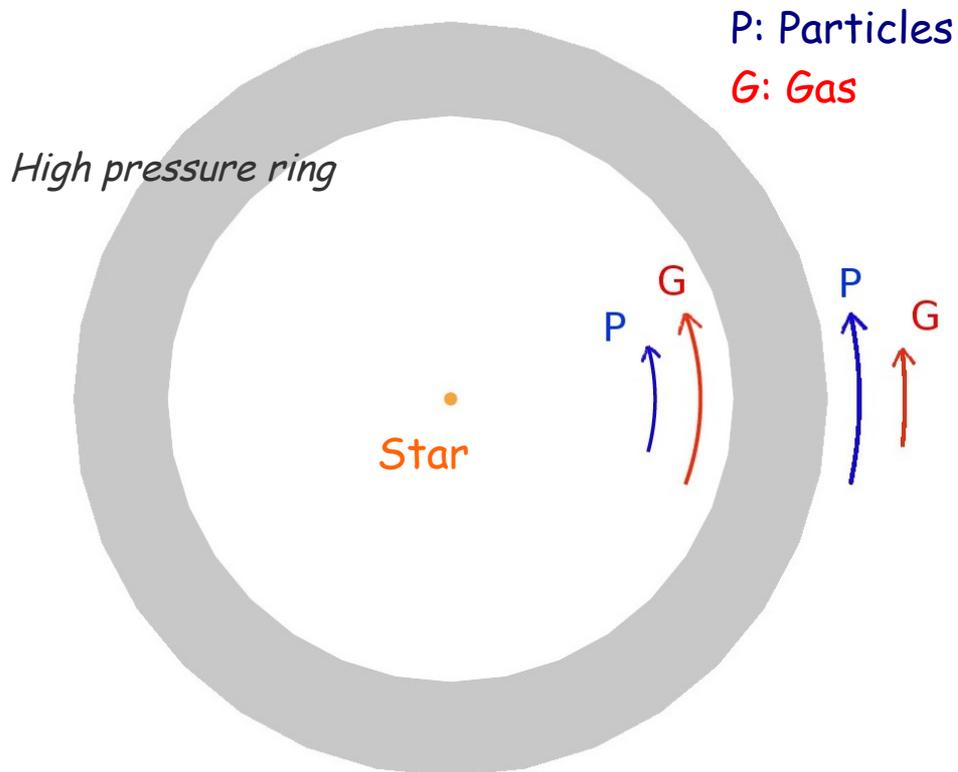
Forming planets in turbulent disks

Gas $\frac{D \mathbf{u}}{Dt} = -\nabla \Phi - \rho^{-1} \nabla p$

$$\mathbf{w} = \mathbf{u} + \tau \rho^{-1} \nabla p$$

The drag force pushes the particles *toward* the pressure gradient

Particles $\frac{d \mathbf{w}}{dt} = -\nabla \Phi - \frac{(\mathbf{w} - \mathbf{u})}{\tau}$



Solid particles

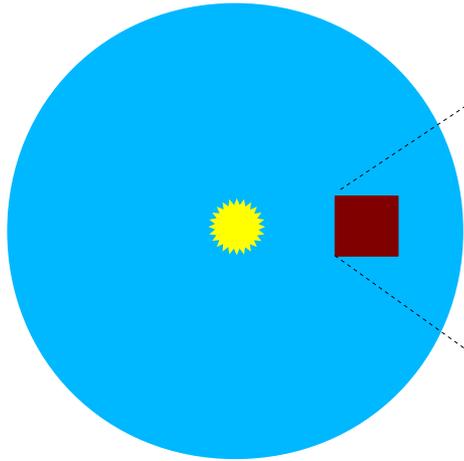
move toward

pressure maxima

Turbulence concentrates solids mechanically in pressure maxima



Gravitational collapse into planetesimals



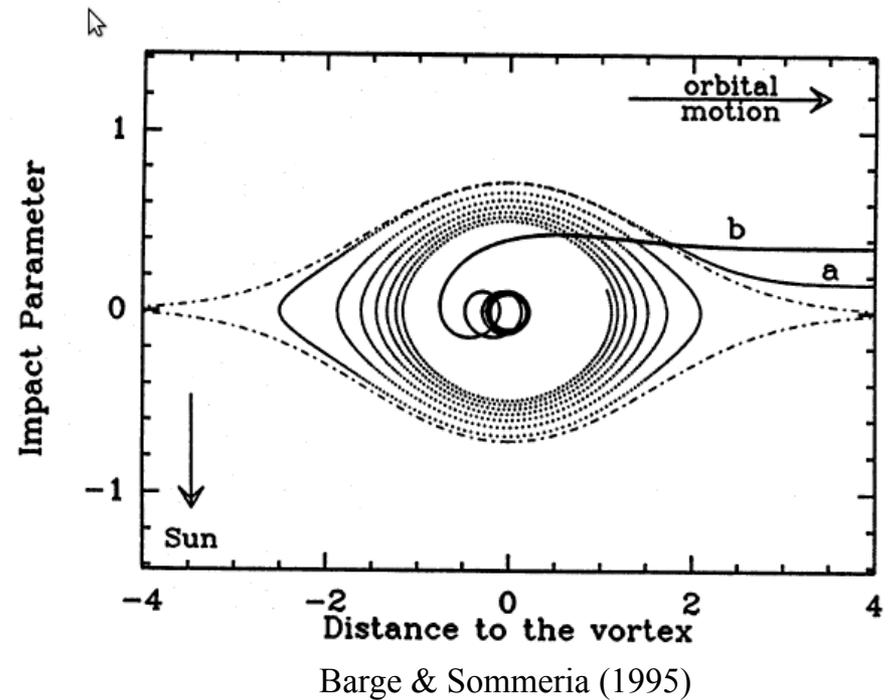
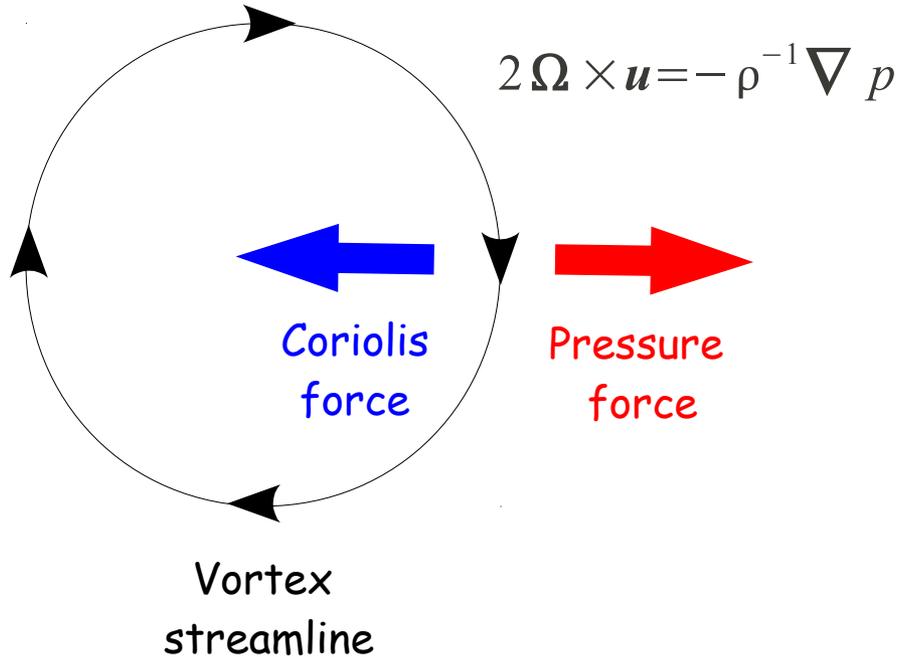
Johansen et al. (2007)

Turbulent eddies concentrate solids,
turning them into planetesimals...

...and vortices are **huge** eddies!

Planet Formation via *Vortex Thruway*

Geostrophic balance:



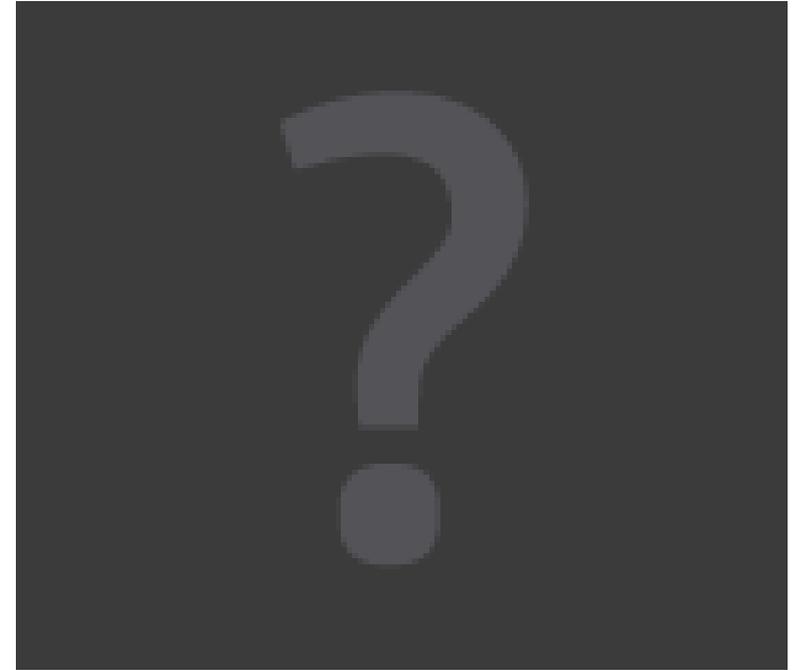
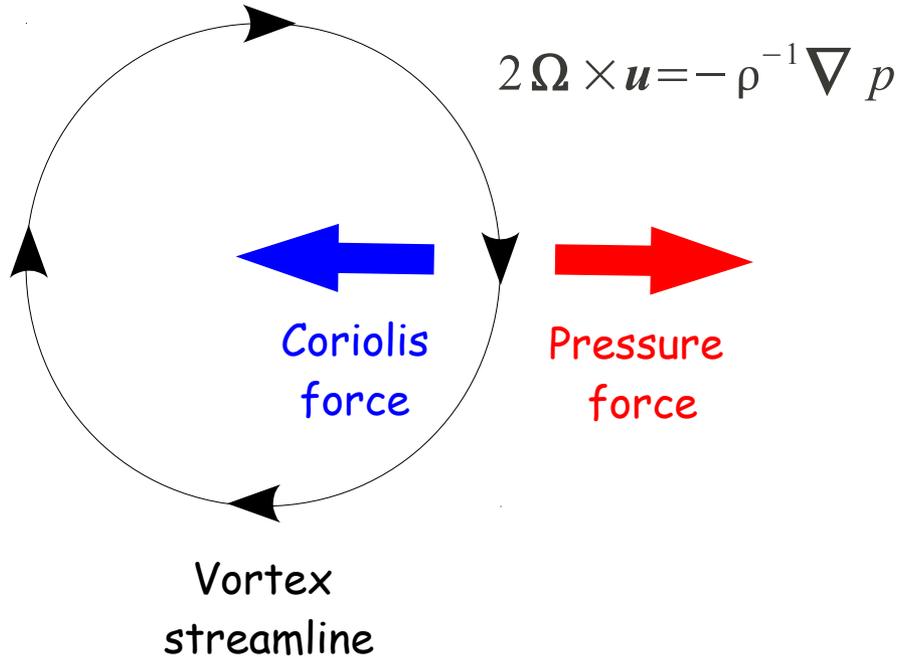
Particles do not feel the pressure gradient.
They sink towards the center, where they accumulate.

Aid to planet formation (Barge & Sommeria 1995)

Speed up planet formation enormously
(Lyra et al. 2008b, 2009a, 2009b, Raettig, Lyra & Klahr 2012)

Planet Formation via *Vortex Thruway*

Geostrophic balance:



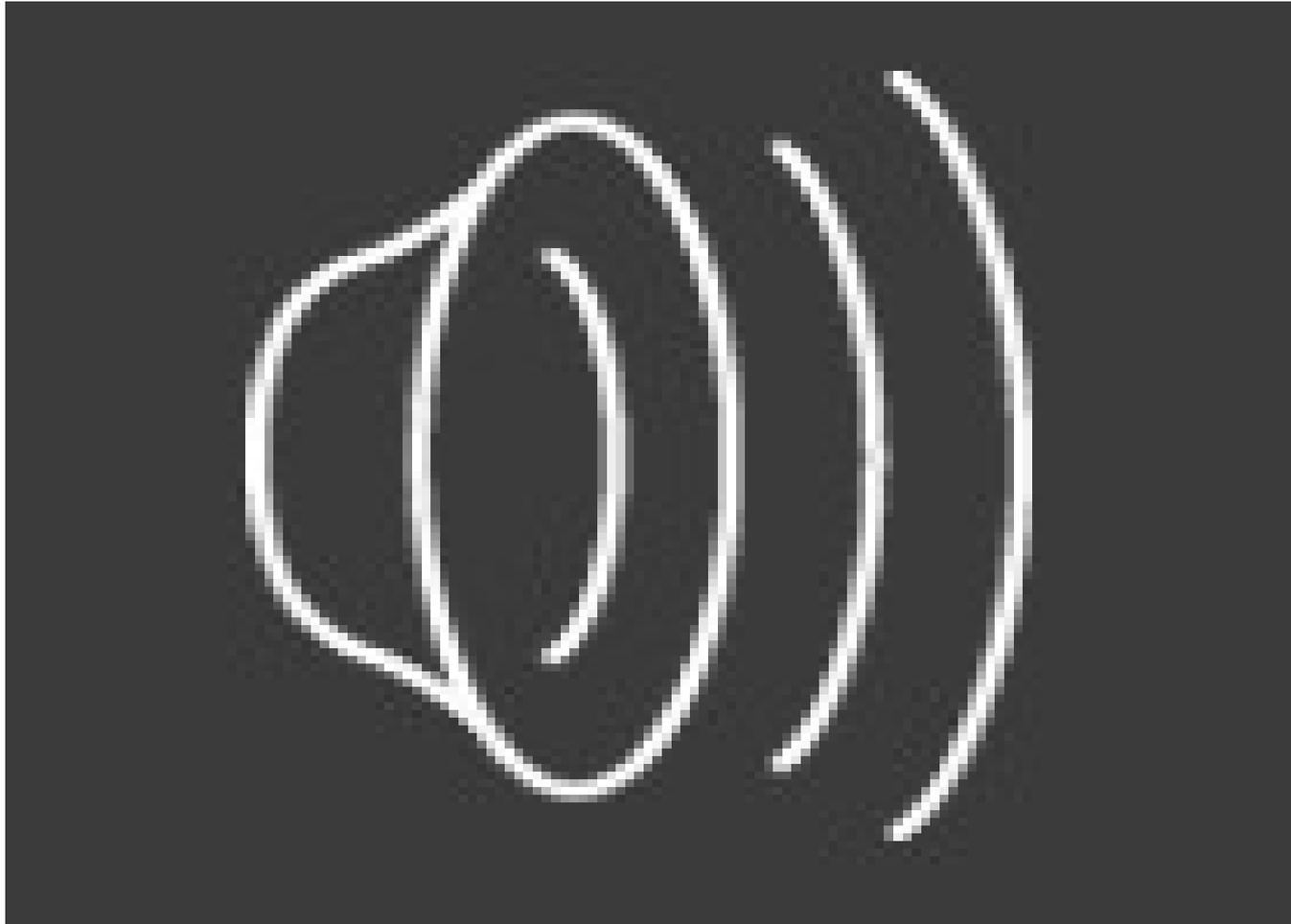
Raettig et al. (2012)

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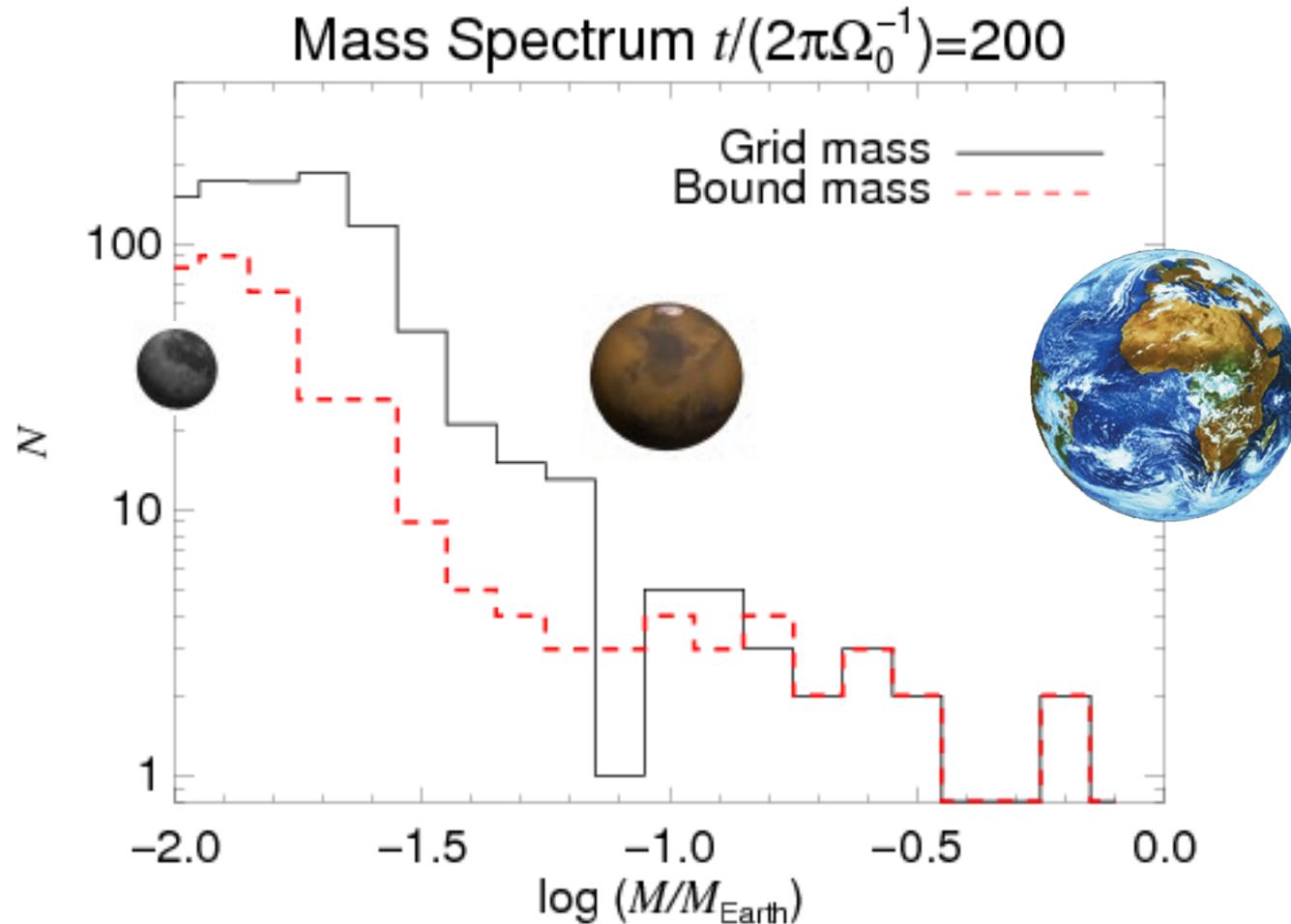
Speed up planet formation enormously
(Lyra et al. 2008b, 2009a, 2009b, Raettig, Lyra & Klahr 2012)

The Initial Mass Function of planets



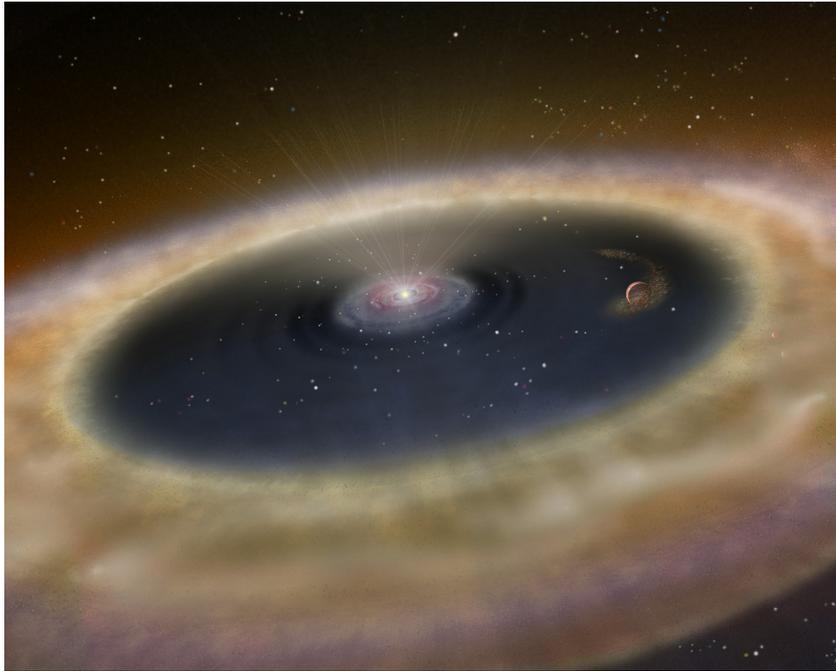
- Mass spectrum by the end of the simulation
 - 300 bound clumps were formed
- Power law $d(\log N)/d(\log M) = -2.3 \pm 0.2$
- 20 of these are more massive than Mars

The Initial Mass Function of planets

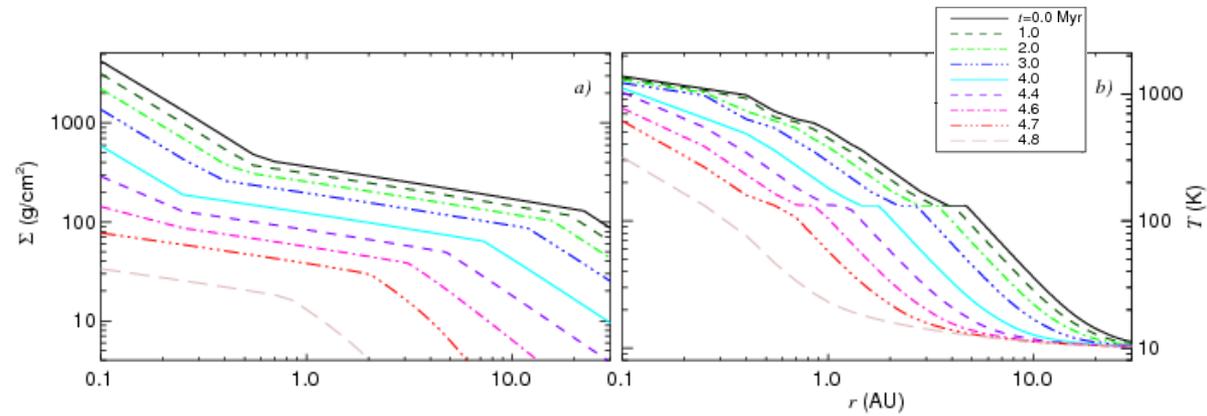


- Mass spectrum by the end of the simulation
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- 20 of these are more massive than Mars

Transitional disks - The thinning phase



Disks evolve in time, due to photoevaporative winds and viscous evolution

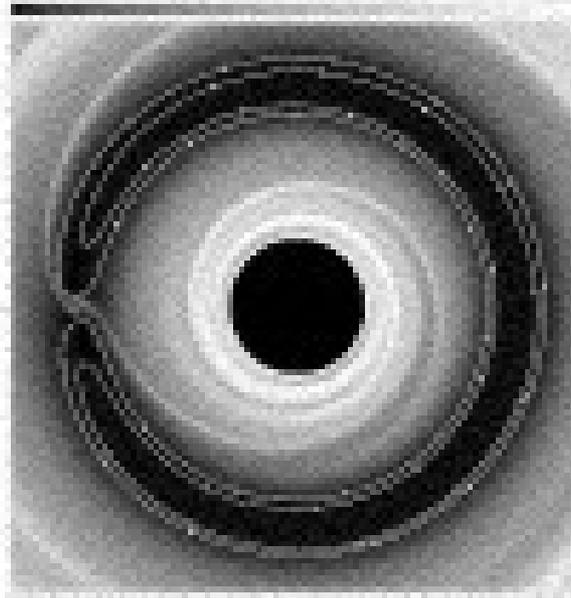


Lyra, Paardekooper, & Mac Low (2010)

Planets form and start to migrate

Planet-disk interaction leads to **angular momentum exchange**

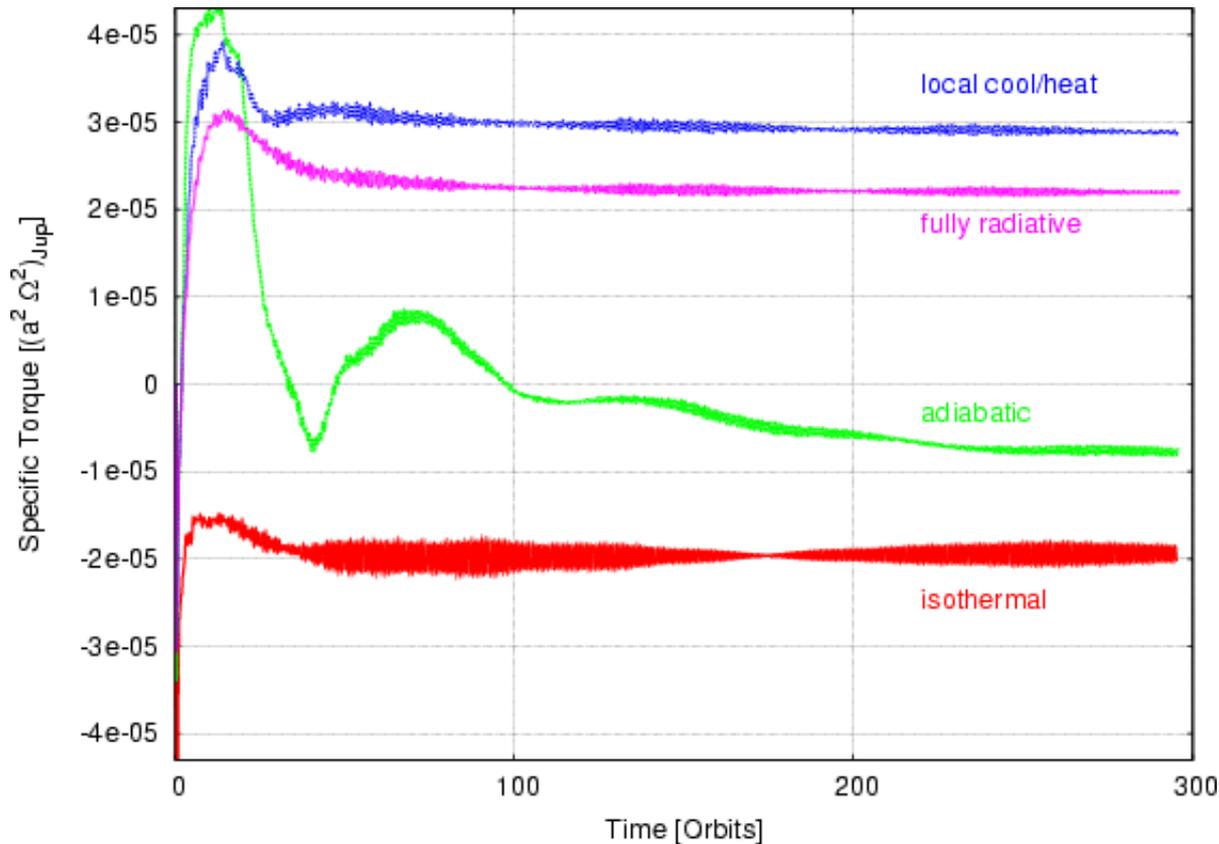
One armed spiral: Lindblad resonance
Horseshoe libration: Co-rotational torques



Lubow et al. (1999)

In isothermal disks,
the result is *inward migration*.

Planets form and start to migrate



Kley & Crida (2008)

Rule of thumb: Migration is
outwards in
steep temperature gradients,
inwards in
isothermal regions.

Paardekooper & Mellema (2006)

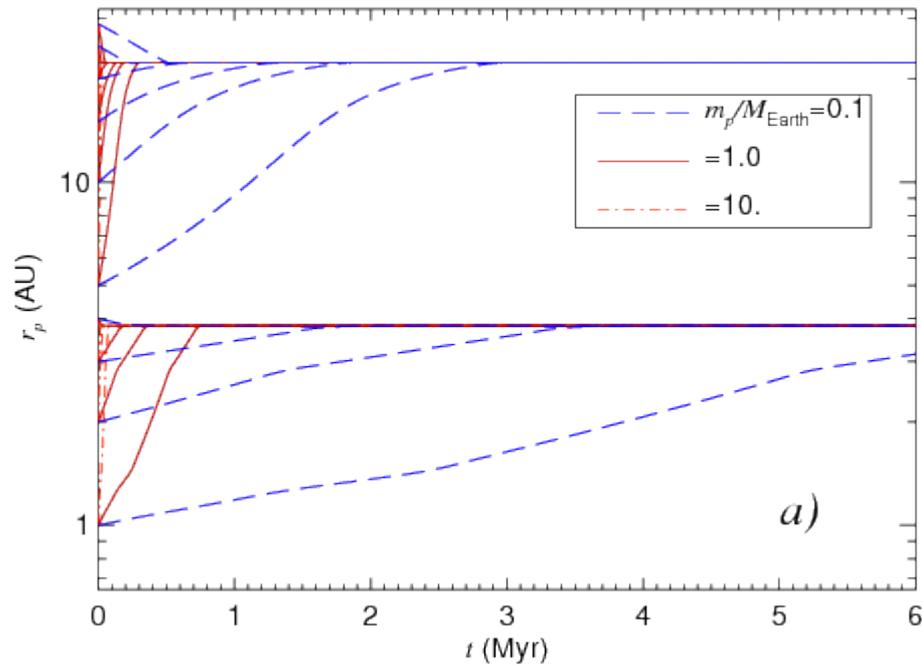
Non-isothermal
co-rotational torque may lead
to outward migration

Hot topic!

- Paardekooper & Mellema 2008
- Baruteau & Masset 2008
- Paardekooper & Papaloizou 2008
- Kley & Crida 2009
- Kley et al 2009
- Paardekooper et al. 2010
- Bitsch & Kley 2010
- Lyra et al. 2010
- Paardekooper et al. 2011
- Ayliffe & Bate 2011
- Yamada & Inaba 2011
- Kley 2011

Planets form and start to migrate

Planet-disk interaction leads to **angular momentum exchange**

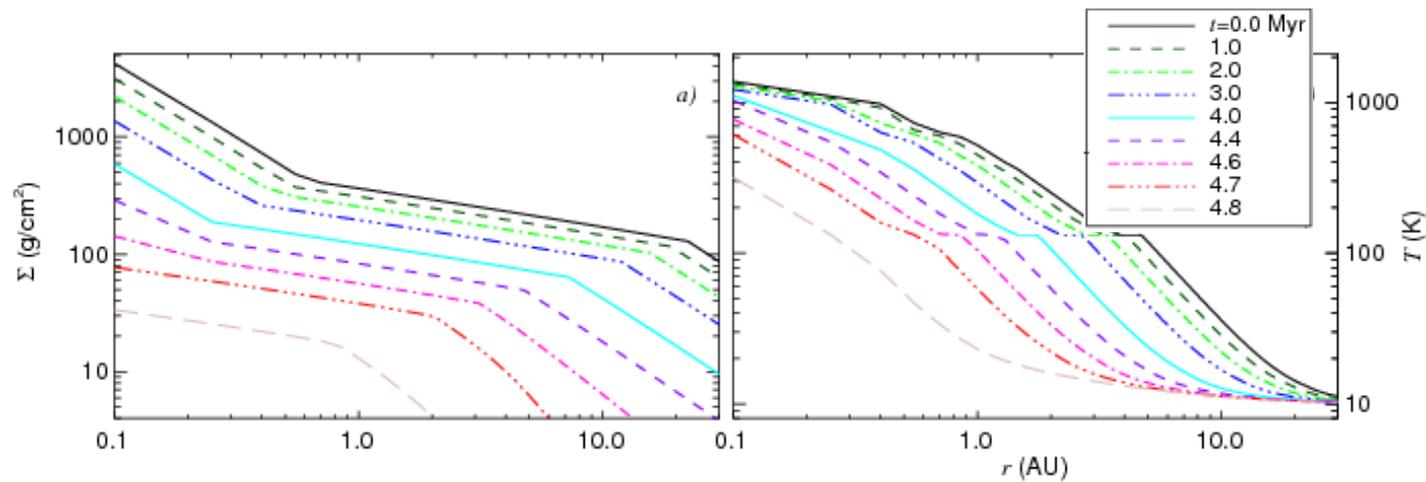


Lyra, Paardekooper, & Mac Low (2010)

Planet traps where migration
is **convergent**
($\tau \neq 0$, $d\tau/dr < 0$).

Migration in Evolutionary Models

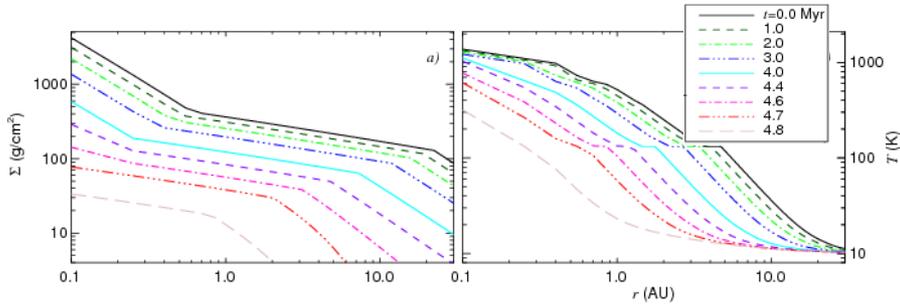
Disks evolve in time, due to
photoevaporative winds and viscous evolution



Lyra, Paardekooper, & Mac Low (2010)

Migration in Evolutionary Models

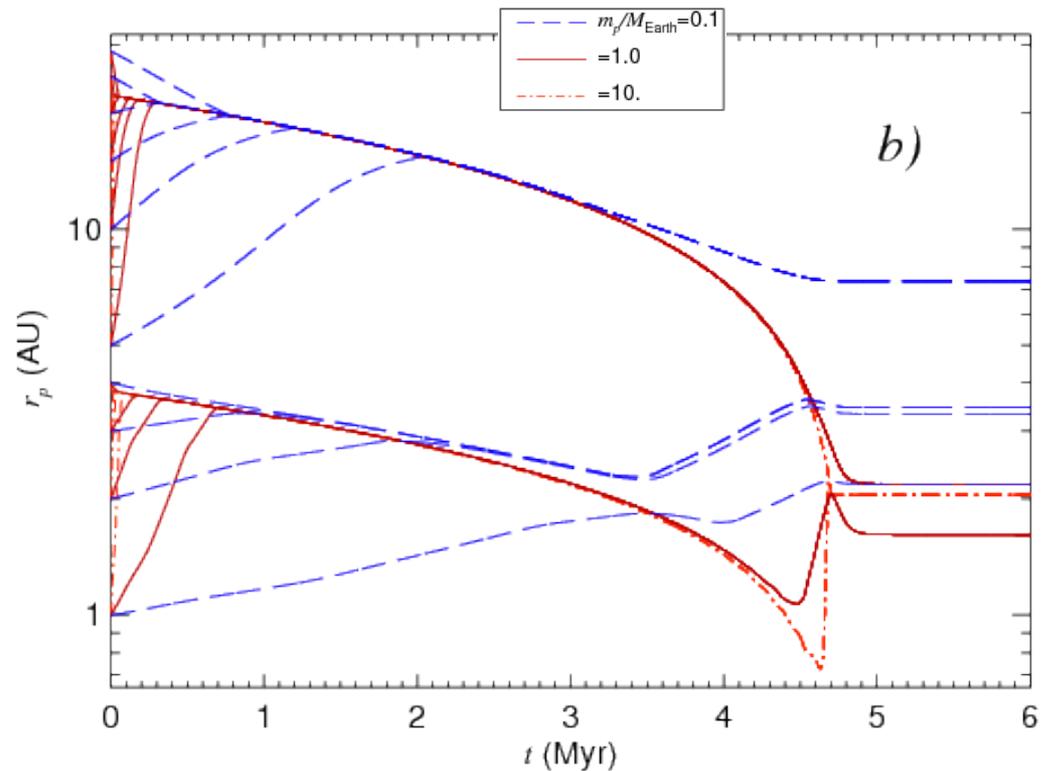
Disks evolve in time, due to photoevaporative winds and viscous evolution



Single planets in a planetary trap evolve in **lockstep with the gas** at the accretion timescale.

At some point, the **disk becomes too thin** to drive accretion. The **planet decouples** and is **released** in a safe orbit.

Rule of thumb: **Migration is outwards in steep temperature gradients, inwards in isothermal regions.**



Lyra, Paardekooper, & Mac Low (2010)

Migration in Evolutionary Models

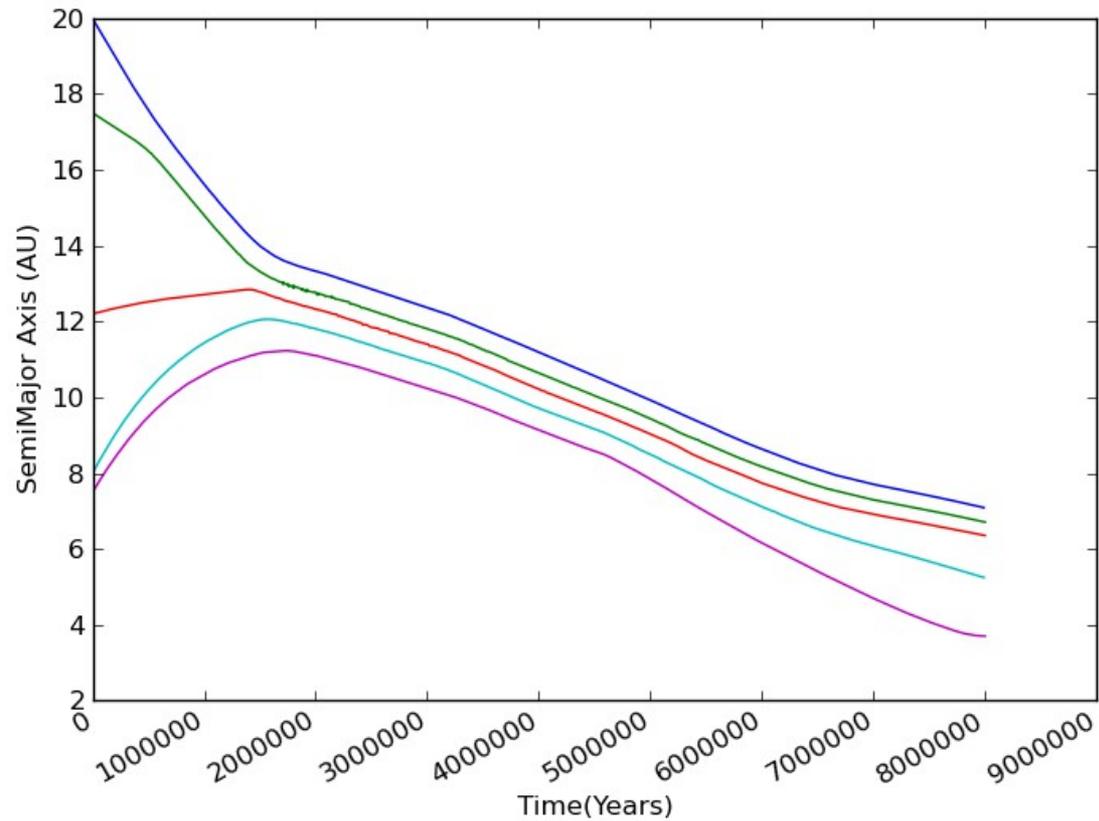
Single planets in a planetary trap evolve in **lockstep with the gas** at the accretion timescale.

At some point, the **disk becomes too thin** to drive accretion.

The planet **decouples** and is **released** in a safe orbit.



Migration + N-Body in Evolutionary Models



Migration in resonance!

see also

Sandor, Lyra & Dullemond (2011)

Hellary & Nelson (2012)

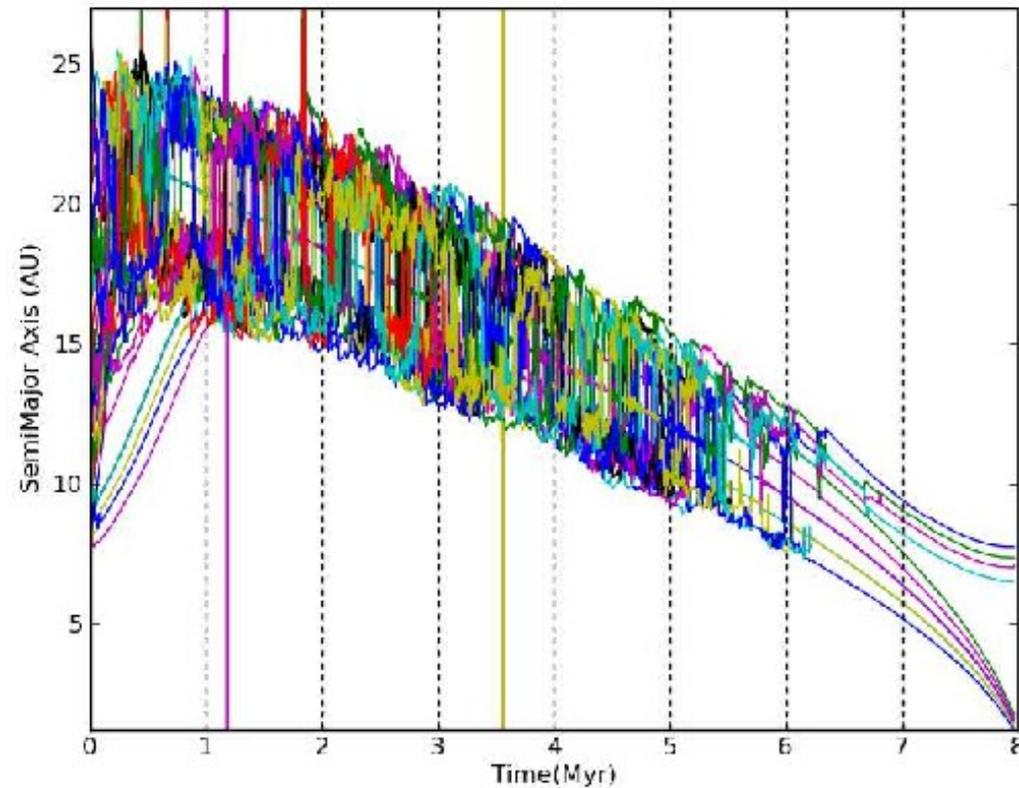
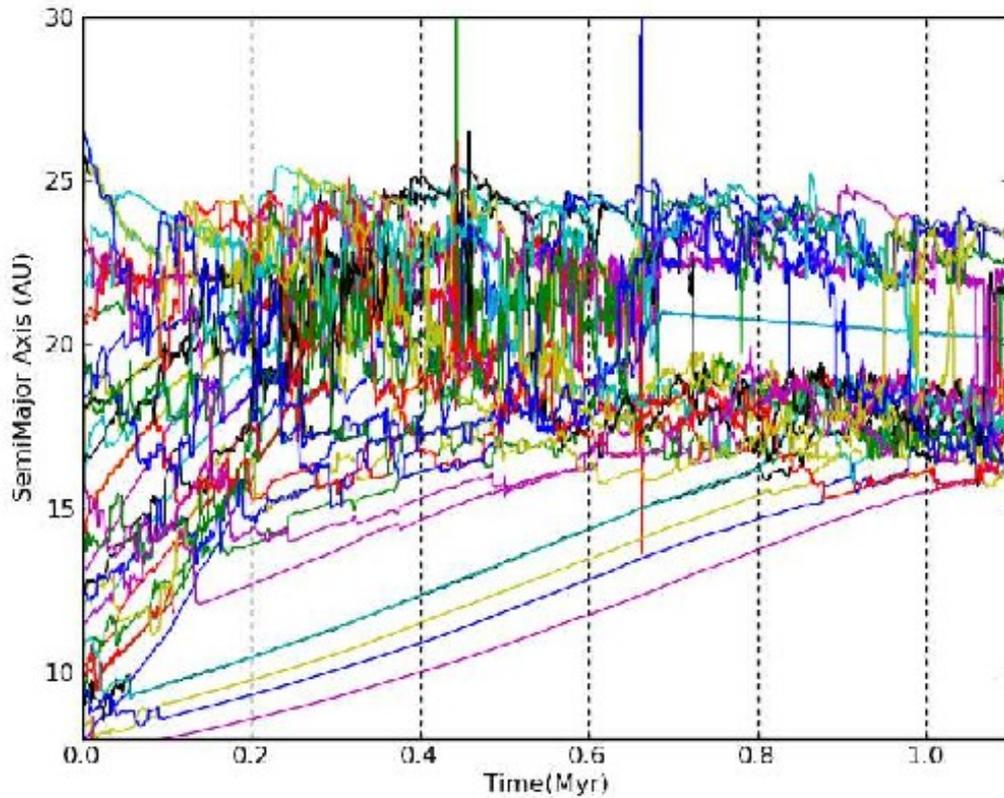
Orbital migration of interacting planets in a radiative evolutionary model

Combines

migration + N-body + photoevaporation + turbulence

modelled as stochastic forcing

(Laughlin et al. 2004, Ogiwara et al. 2007)

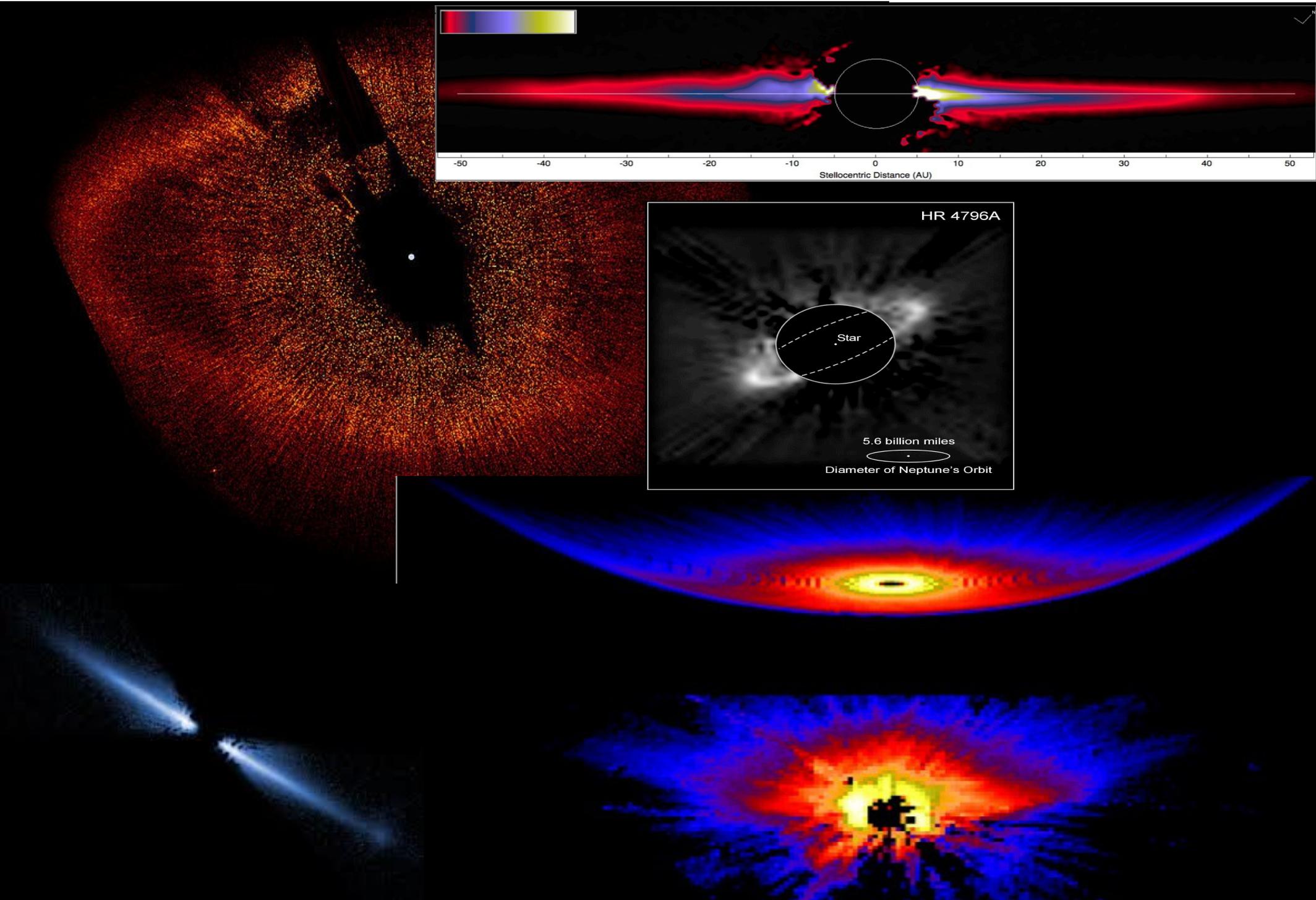


Orbital migration of interacting planets in a radiative evolutionary model

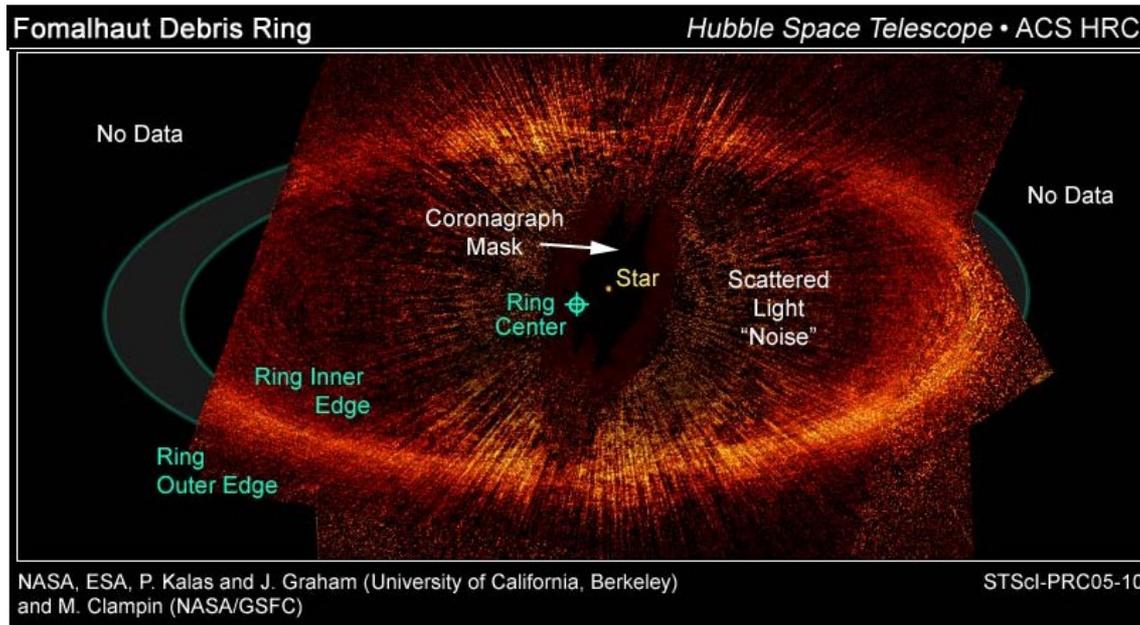


Horn et al. (2012)

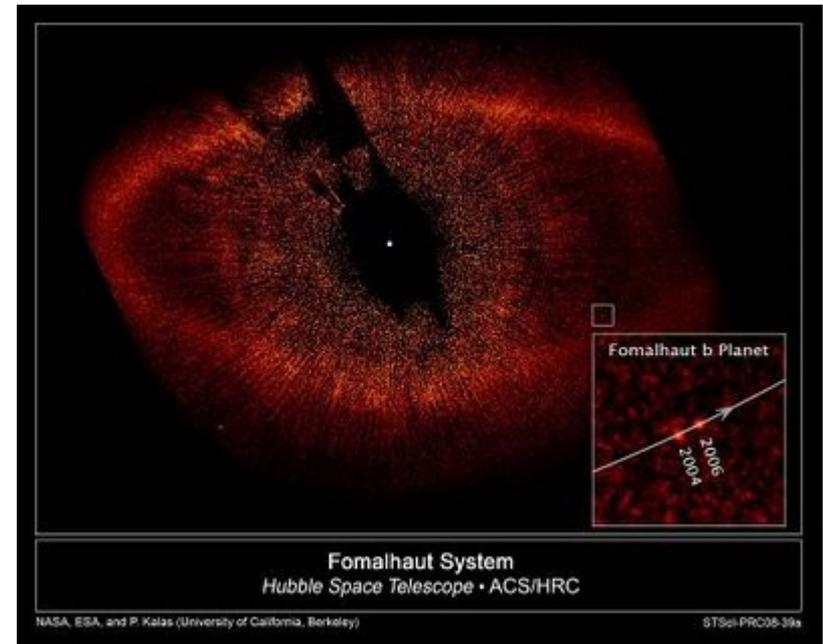
Debris disks - The gas-poor phase



Sharp and eccentric rings in debris disks: Signposts of planets



Narrow sharp eccentric ring



Detection of a source
quickly heralded as a planet
Fomalhaut b

Sharp and eccentric rings in debris disks: Signposts of planets ?

However.....

THE ASTROPHYSICAL JOURNAL, 747:116 (7pp), 2012 March 10
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doi:10.1088/0004-637X/747/2/116

INFRARED NON-DETECTION OF FOMALHAUT b: IMPLICATIONS FOR THE PLANET INTERPRETATION

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¹Department of Astrophysical Sciences, Princeton University, Princeton, NJ, USA; janson@astro.princeton.edu

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Received 2011 December 16; accepted 2012 January 12; published 2012 February 23

ABSTRACT

The nearby A4-type star Fomalhaut hosts a debris belt in the form of an eccentric ring, which is thought to be caused by dynamical influence from a giant planet companion. In 2008, a detection of a point source inside the inner edge of the ring was reported and was interpreted as a direct image of the planet, named Fomalhaut b. The detection was made at $\sim 600\text{--}800$ nm, but no corresponding signatures were found in the near-infrared range, where the bulk emission of such a planet should be expected. Here, we present deep observations of Fomalhaut with *Spitzer*/IRAC at $4.5\ \mu\text{m}$, using a novel point-spread function subtraction technique based on angular differential imaging and Locally Optimized Combination of Images, in order to substantially improve the *Spitzer* contrast at small separations. The results provide more than an order of magnitude improvement in the upper flux limit of Fomalhaut b and exclude the possibility that any flux from a giant planet surface contributes to the observed flux at visible wavelengths. This renders any direct connection between the observed light source and the dynamically inferred giant planet highly unlikely. We discuss several possible interpretations of the total body of observations of the Fomalhaut system and find that the interpretation that best matches the available data for the observed source is scattered light from a transient or semi-transient dust cloud.

Key words: circumstellar matter – planetary systems – stars: early-type

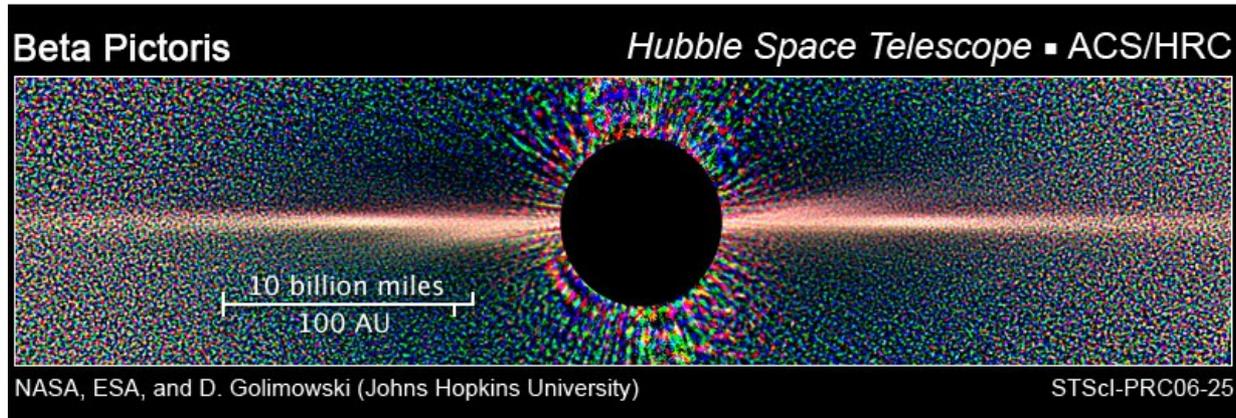
Online-only material: color figures

Planet not detected in infrared

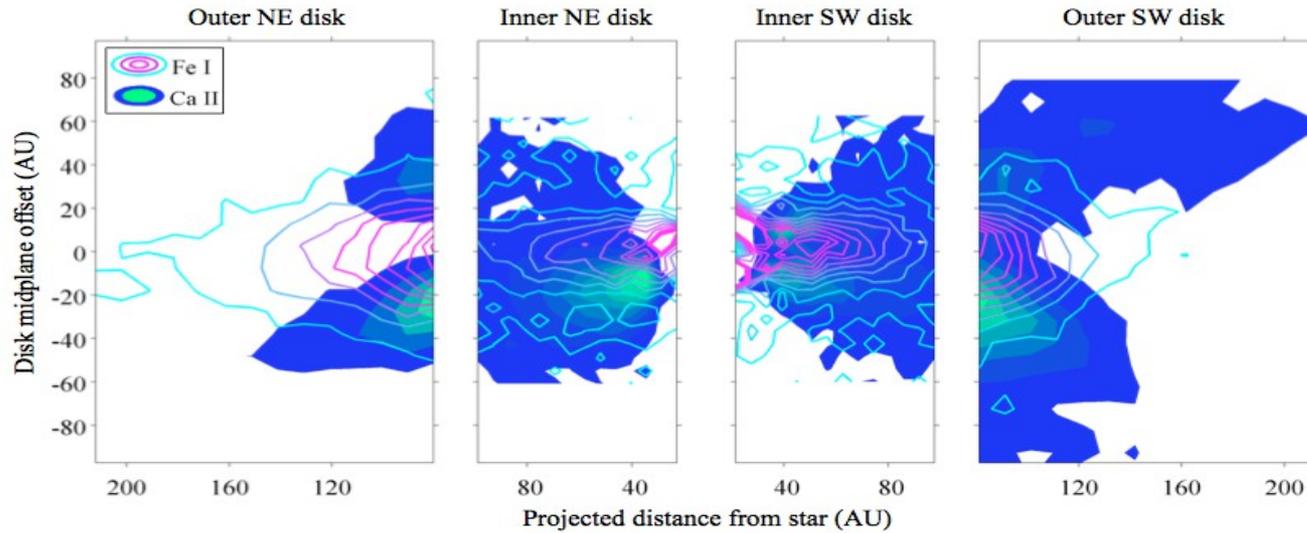
Are there
alternative explanations?

Debris disks are not completely gas-free

Dust



Gas



VLT imaging by
Nilsson et al. (2012)

Gas in debris disks

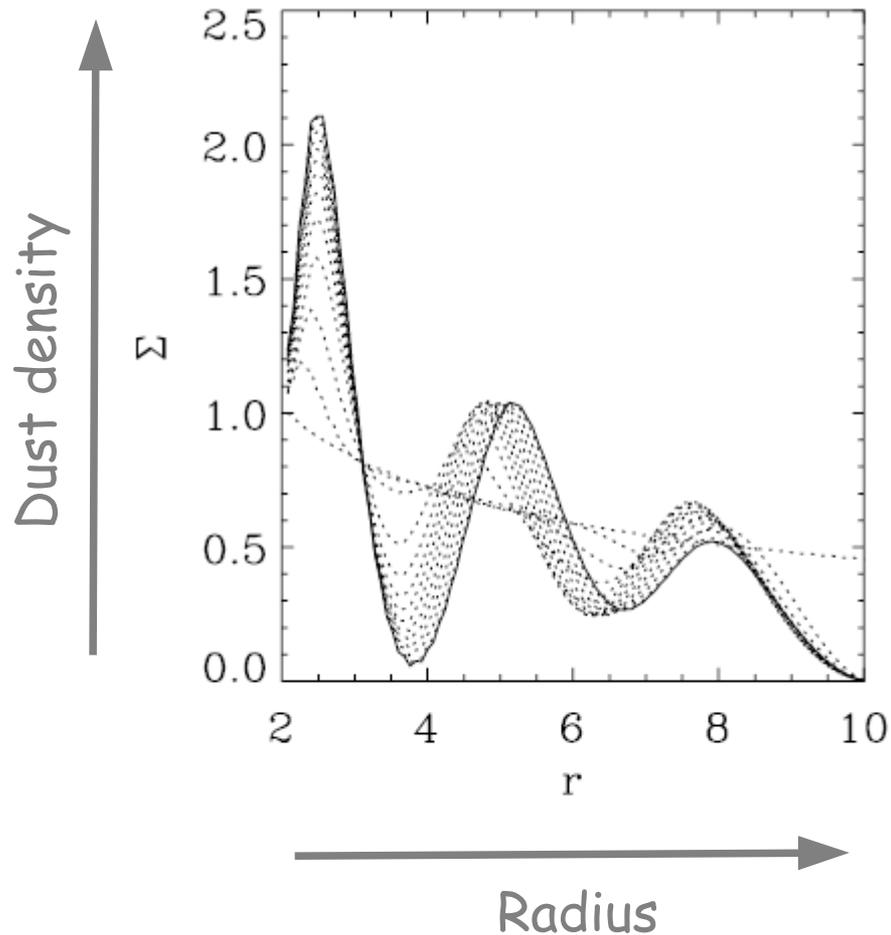
Detections

β Pictoris	many species	Lagrange et al. (1998), ...
51 Ophiuchi	many species	Roberge et al. (2002)
σ Herculis	C II, N II	Chen & Jura (2003)
HD 32297	Na I, CII	Redfield (2007), Donaldson et al. (2012)
HD 135344	H ₂ , CO	Thi et al. (2001), Pontoppidan et al. (2008)
49 Ceti	H ₂ , CO	Dent et al. (2005), Roberge et al. (2012)
AU Mic	H ₂	France et al. (2007)
HD172555	SiO	Lisse et al. (2009)

Source of gas: Outgassing processes

Infalling comets	Beust & Valiron (2007)
Grain sublimation	e.g. Rafikov (2012)
Grain-Grain collisions	Czechowski & Mann (2007)
Photo-stimulated desorption	Chen et al. (2007)
Planet-Planet collisions	Van den Ancker (2001), Lisse (2008)
Primordial?	

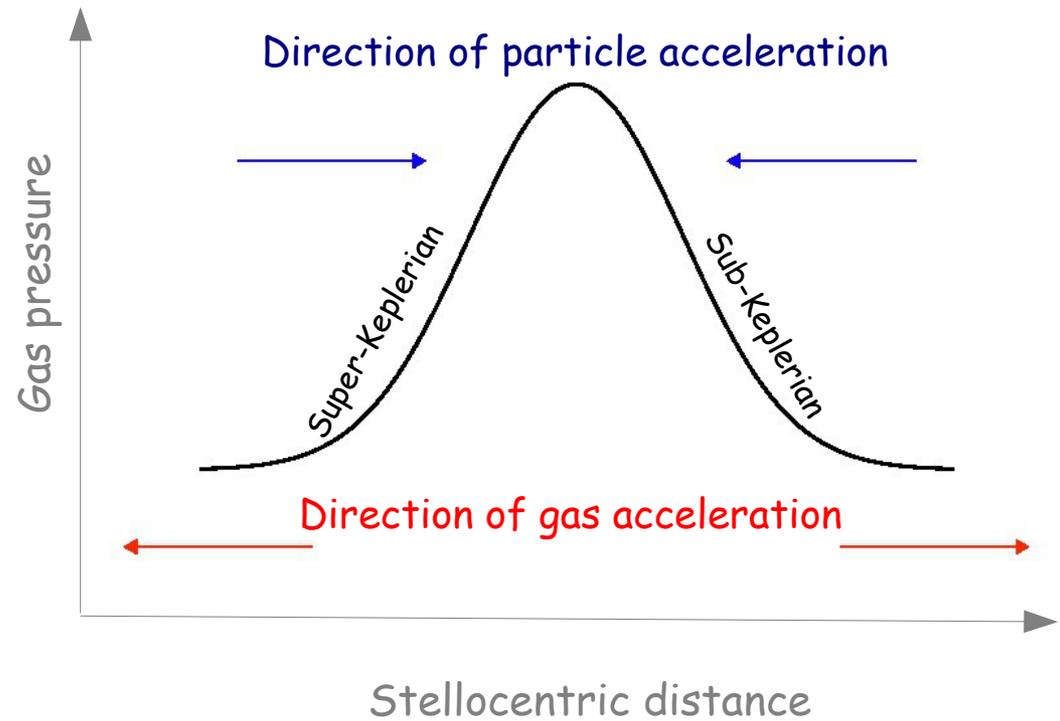
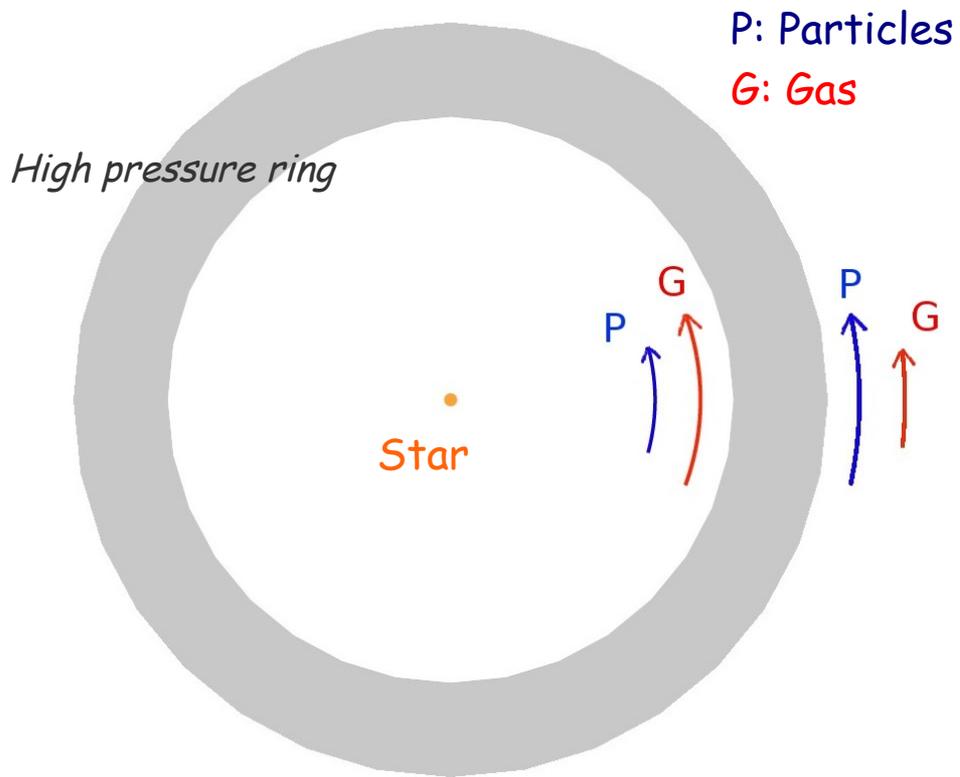
Dust and gas together leads to instability...



Klahr & Lin (2005)

Suggested that an instability might cause dust in debris disks to clump together.

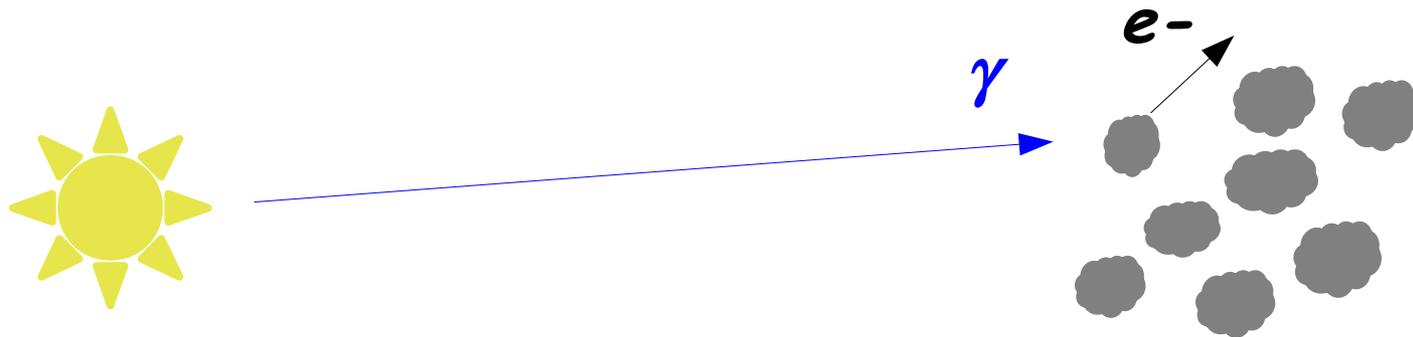
Particles move toward pressure maxima



Adapted from Whipple (1972)

Photoelectric heating

In optically thin debris disks,
the **dust** is the **main heating agent** for the gas.



Dust intercepts starlight directly,
emits electron, that heats the gas.

Gas is photoelectrically heated by the dust itself

Runaway process: instability

Dust heats gas

Heated gas = high pressure region

High pressure concentrates dust

Runaway process: instability



Dust heats gas

Heated gas = high pressure region

High pressure concentrates dust



Model equations

Klahr & Lin (2005) used a simplified, 1-D model.

$$\frac{\partial}{\partial t} \Sigma_d + \frac{1}{r} \frac{\partial}{\partial r} r \Sigma_d v_r = 0.$$

Continuity equation

$$V_\phi = \Omega r + \frac{1}{2\Omega \Sigma_g} \frac{\partial}{\partial r} P$$

Terminal velocity

$$T_g = T_0 \left(\frac{\Sigma_d}{\Sigma_0} \right)^\beta,$$

Equation of state

Model equations

Our simulation adds much more physics, and works in 2D.

Klahr & Lin (2005)

1D

$$\frac{\partial}{\partial t} \Sigma_d + \frac{1}{r} \frac{\partial}{\partial r} r \Sigma_d v_r = 0.$$

$$V_\phi = \Omega r + \frac{1}{2\Omega \Sigma_g} \frac{\partial}{\partial r} P$$

$$T_g = T_0 \left(\frac{\Sigma_d}{\Sigma_0} \right)^\beta,$$

Inertia for both gas and dust

Energy equation

*Drag force and
drag force backreaction*

Lyra & Kuchner (2012)

2D

$$\frac{\partial \Sigma_g}{\partial t} = -(\mathbf{u} \cdot \nabla) \Sigma_g - \Sigma_g \nabla \cdot \mathbf{u}$$

$$\frac{\partial \mathbf{u}}{\partial t} = -(\mathbf{u} \cdot \nabla) \mathbf{u} - \frac{1}{\Sigma_g} \nabla P - \nabla \Phi - \frac{\Sigma_d}{\Sigma_g} \mathbf{f}_d$$

$$\frac{\partial S}{\partial t} = -(\mathbf{u} \cdot \nabla) S - \frac{c_v (T - T_p)}{T \tau_T}.$$

$$\frac{d\mathbf{x}}{dt} = \mathbf{v}$$

$$\frac{d\mathbf{v}}{dt} = -\nabla \Phi + \mathbf{f}_d$$

$$\mathbf{f}_d = -\frac{(\mathbf{v} - \mathbf{u})}{\tau_f}$$

$$T_p = T_0 \frac{\Sigma_d}{\Sigma_0}.$$

Linear Analysis

$$D_w \Sigma_d = -\Sigma_d \nabla \cdot \mathbf{w}$$

Dust

$$D_w w_x = 2\Omega w_y - \frac{1}{\tau_f}(w_x - u_x)$$

$$D_w w_y = -\frac{1}{2}\Omega w_x - \frac{1}{\tau_f}(v_y - u_y)$$

$$D_u \Sigma_g = -\Sigma_g \nabla \cdot \mathbf{u}$$

Gas

$$D_u u_x = 2\Omega u_y - \frac{1}{\Sigma_g} \frac{\partial P}{\partial x} - \frac{\epsilon}{\tau_f}(u_x - w_x)$$

$$D_u u_y = -\frac{1}{2}\Omega u_x - \frac{1}{\Sigma_g} \frac{\partial P}{\partial y} - \frac{\epsilon}{\tau_f}(u_y - w_y)$$

$$\psi = \psi_0 + \psi'$$

$$\psi' = \hat{\psi} \exp(ikx + st)$$



Dispersion relation

$$A\omega^5 + B\omega^4 + C\omega^3 + D\omega^2 + E\omega + F = 0$$

$$A = 1$$

$$B = 2\epsilon + 2$$

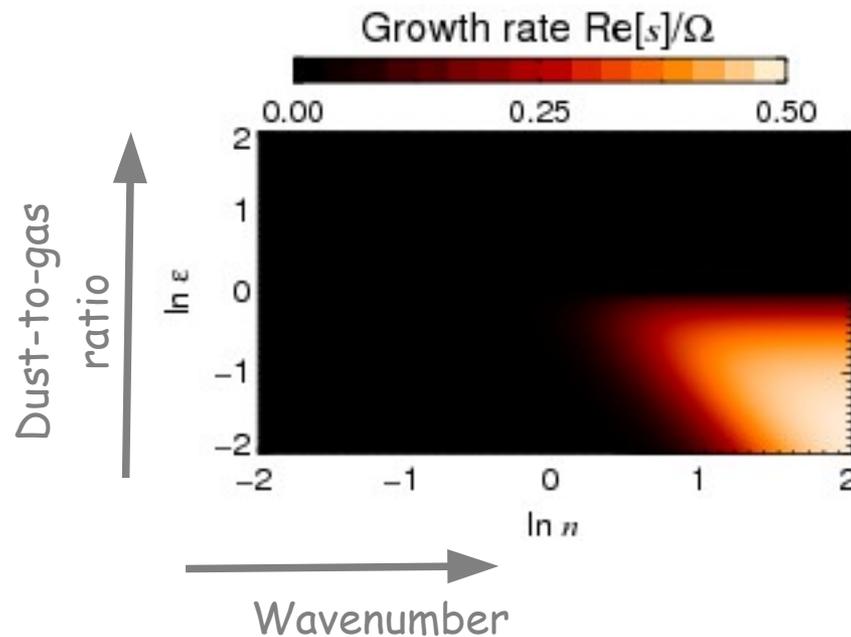
$$C = \epsilon^2 + \epsilon(n^2 + 2) + 3$$

$$D = \epsilon^2 n^2 + \epsilon(3n^2 + 2) + 2$$

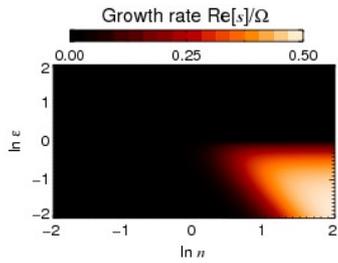
$$E = \epsilon^2(2n^2 + 1) + \epsilon(3n^2 + 2) + 2$$

$$F = \epsilon^2 n^2 - \epsilon n^2$$

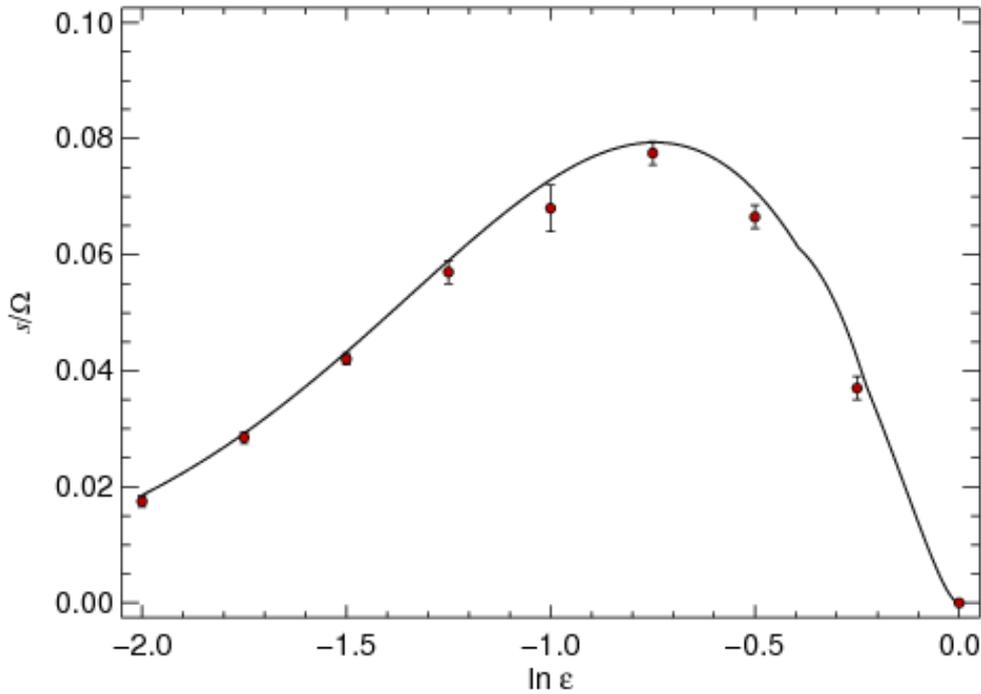
$$\epsilon = \Sigma_d / \Sigma_g \quad n = kH \quad \omega = s/\Omega$$



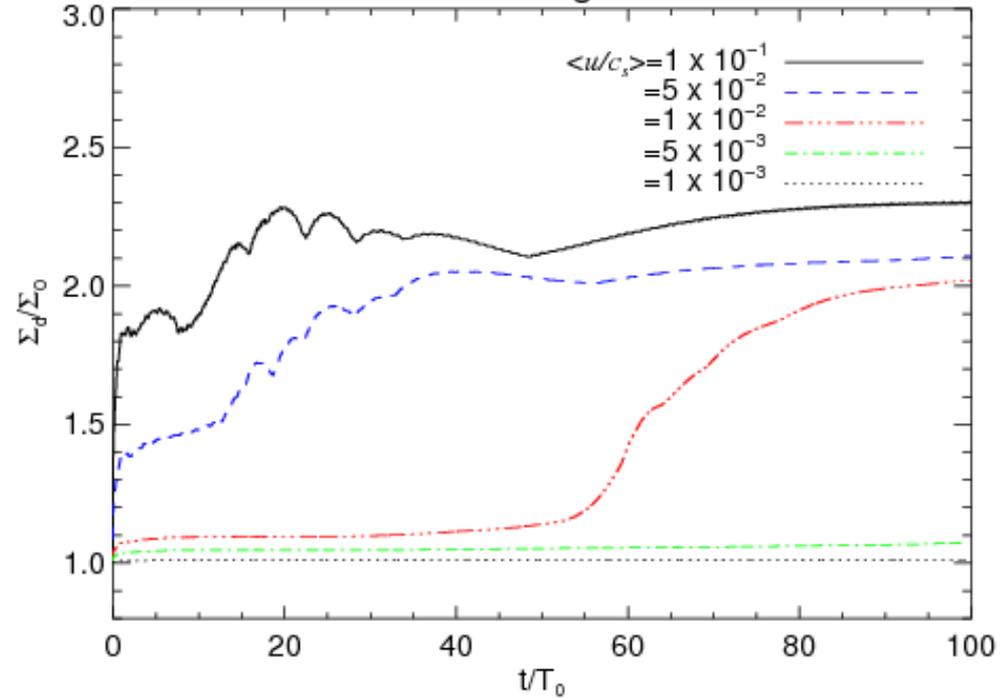
Linear and nonlinear growth



Growth rates $\alpha=10^{-2}$



Nonlinear growth



Linear growth only exists for $\varepsilon < 1$

But there is
nonlinear growth
 beyond !

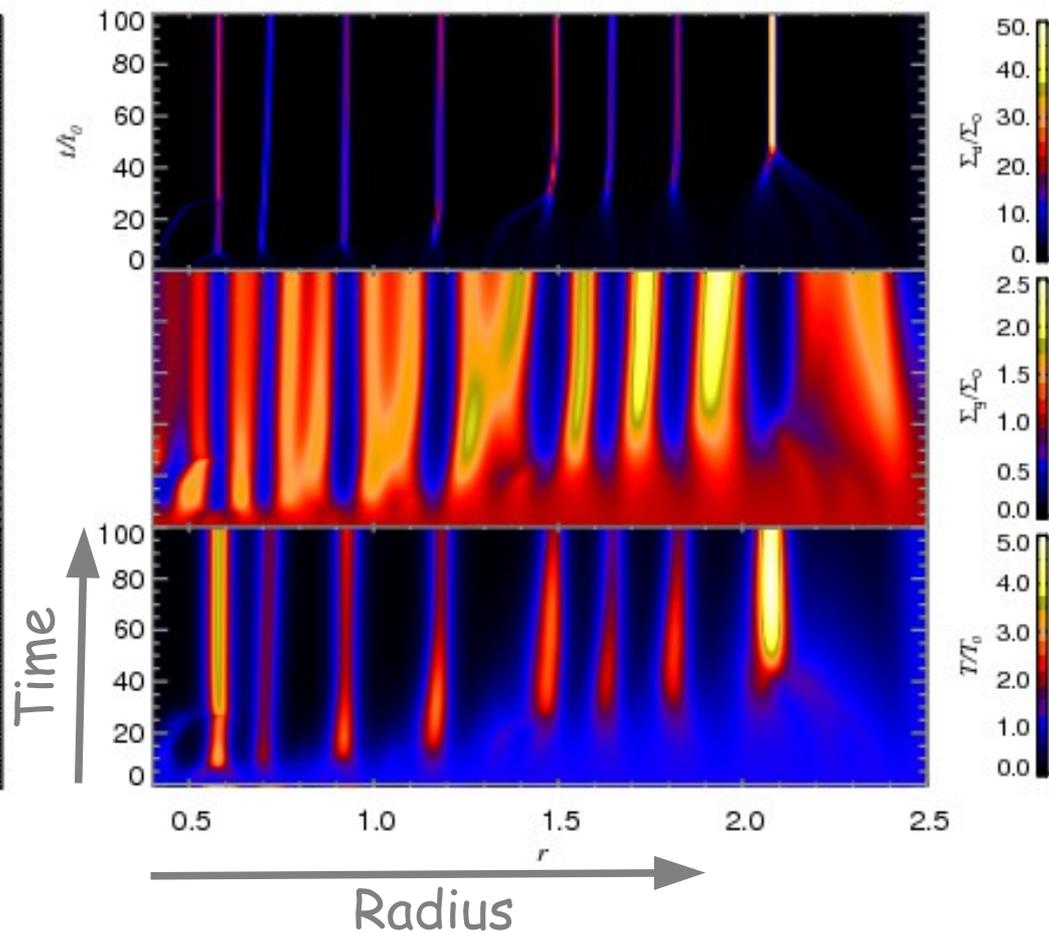
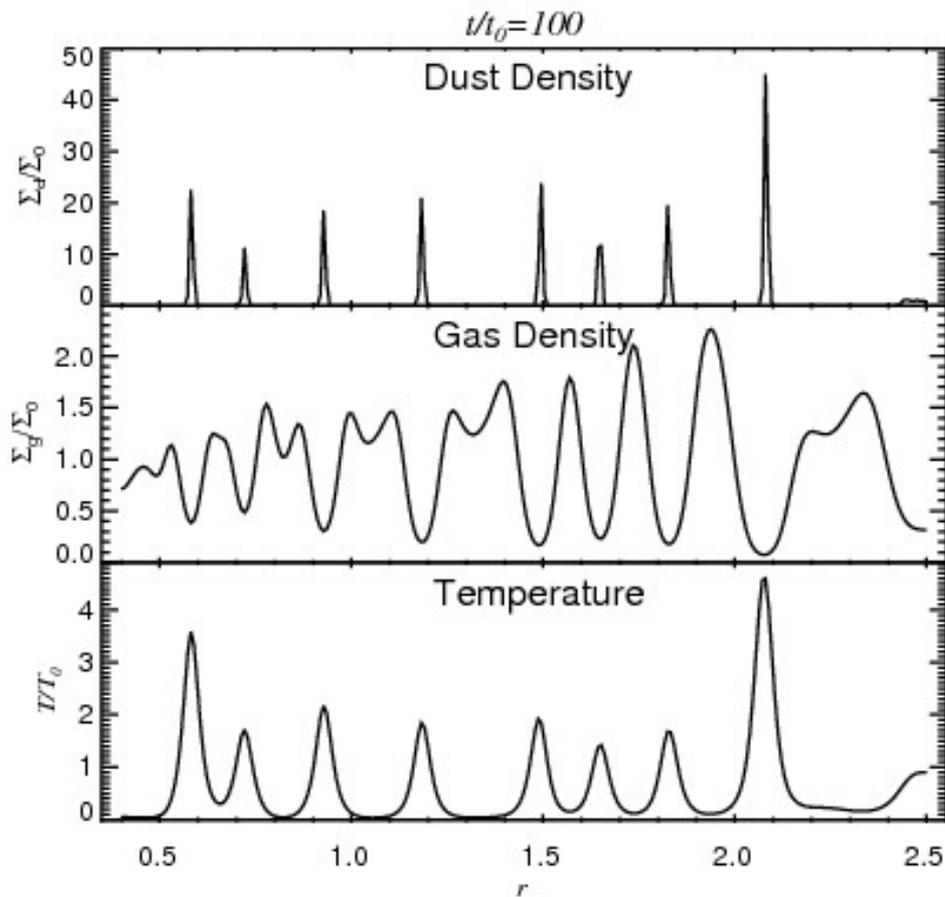
Instability



Dust heats gas

Heated gas = high pressure region

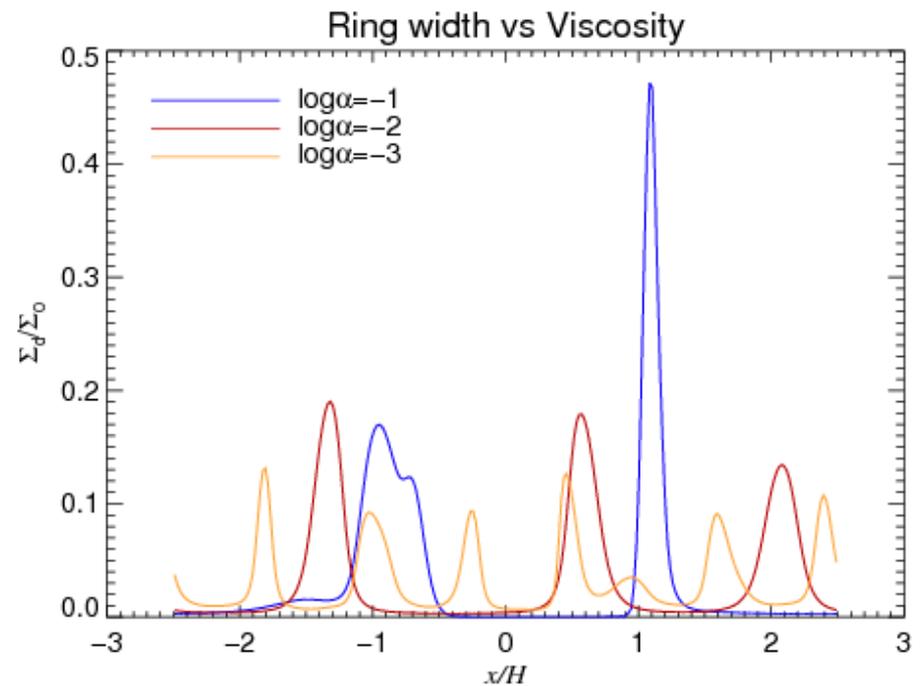
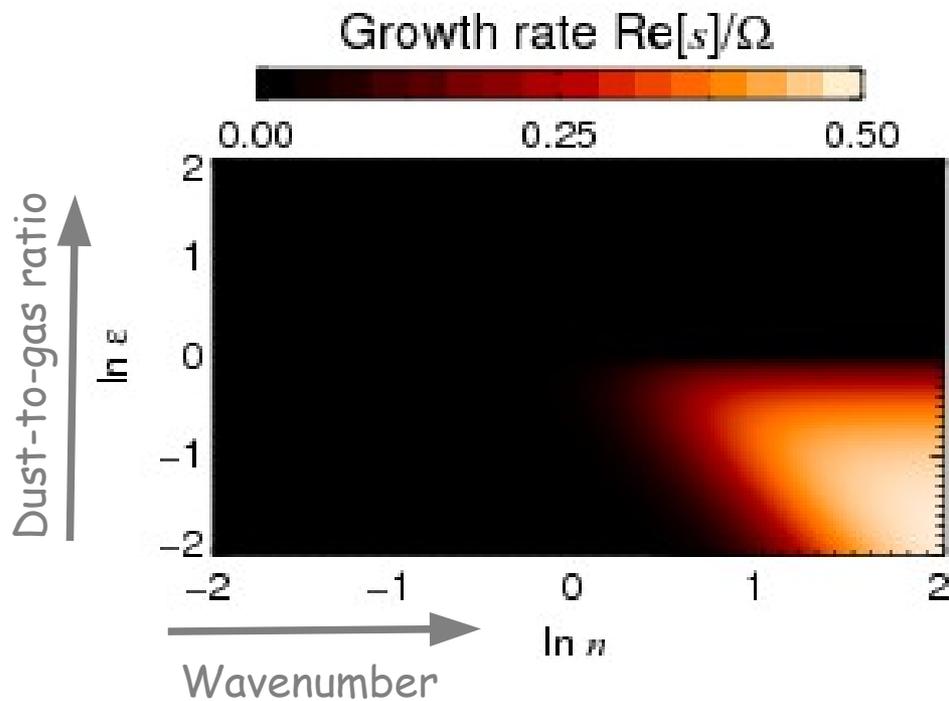
High pressure concentrates dust



Narrow hot dust rings
Cold gas collects between rings

Ring width

Ring spacing and width is determined by the wavelength of maximum growth.

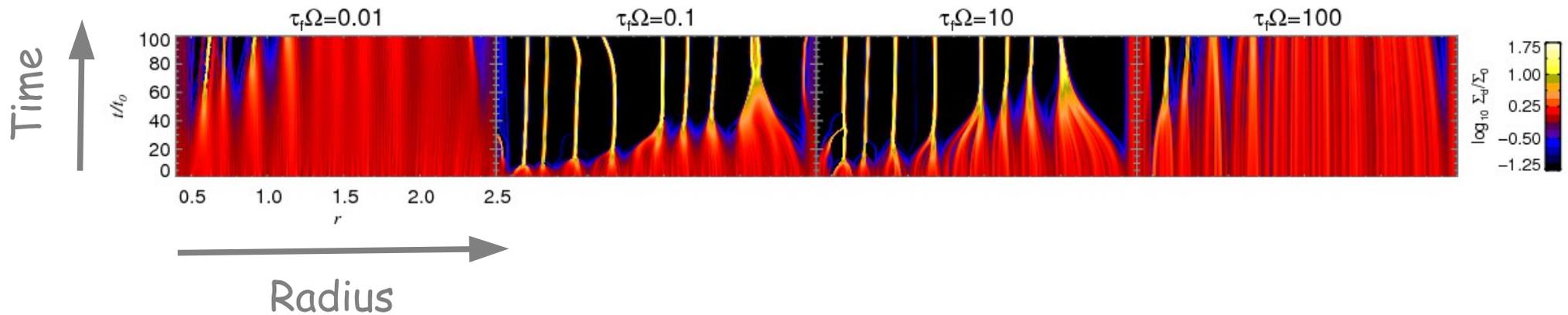


Which in turn is determined by viscosity

Ring width ~ 10 Kolmogorov lengths

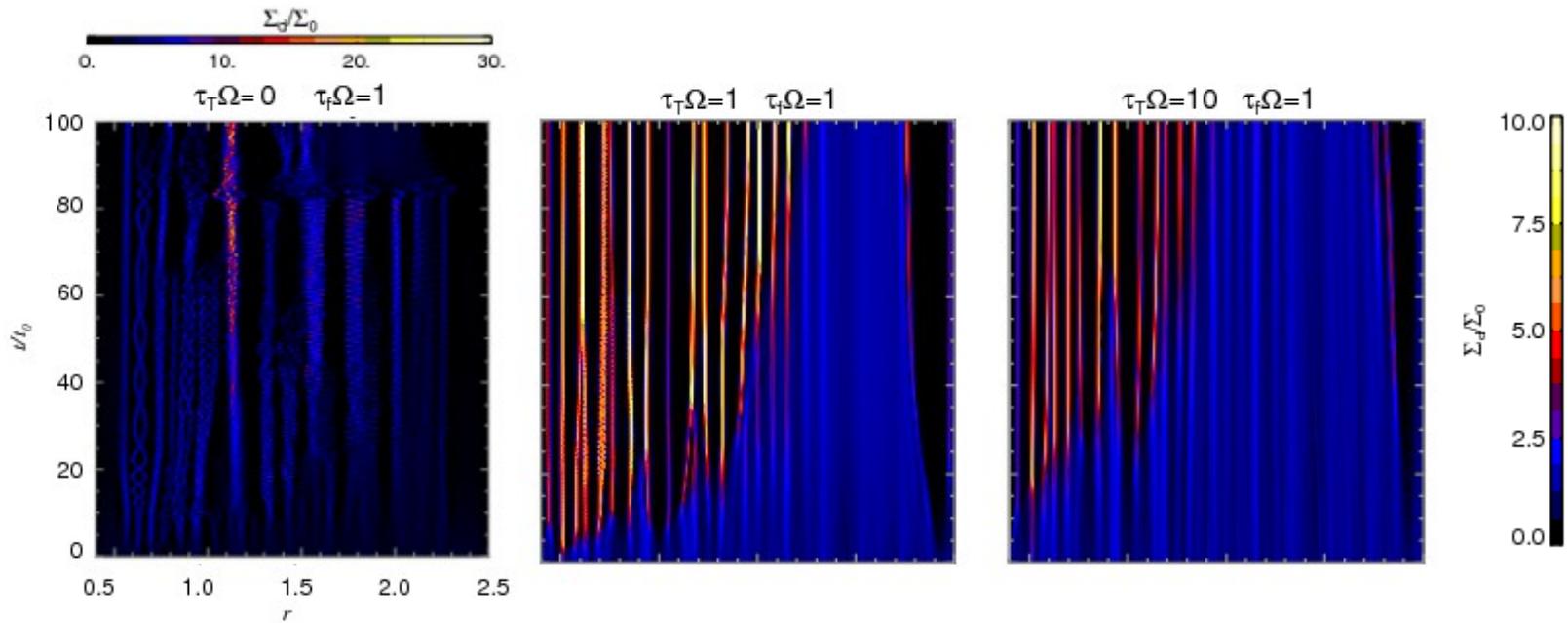
Robustness

Growth over 4 orders of magnitude in dust-gas coupling time (friction time)



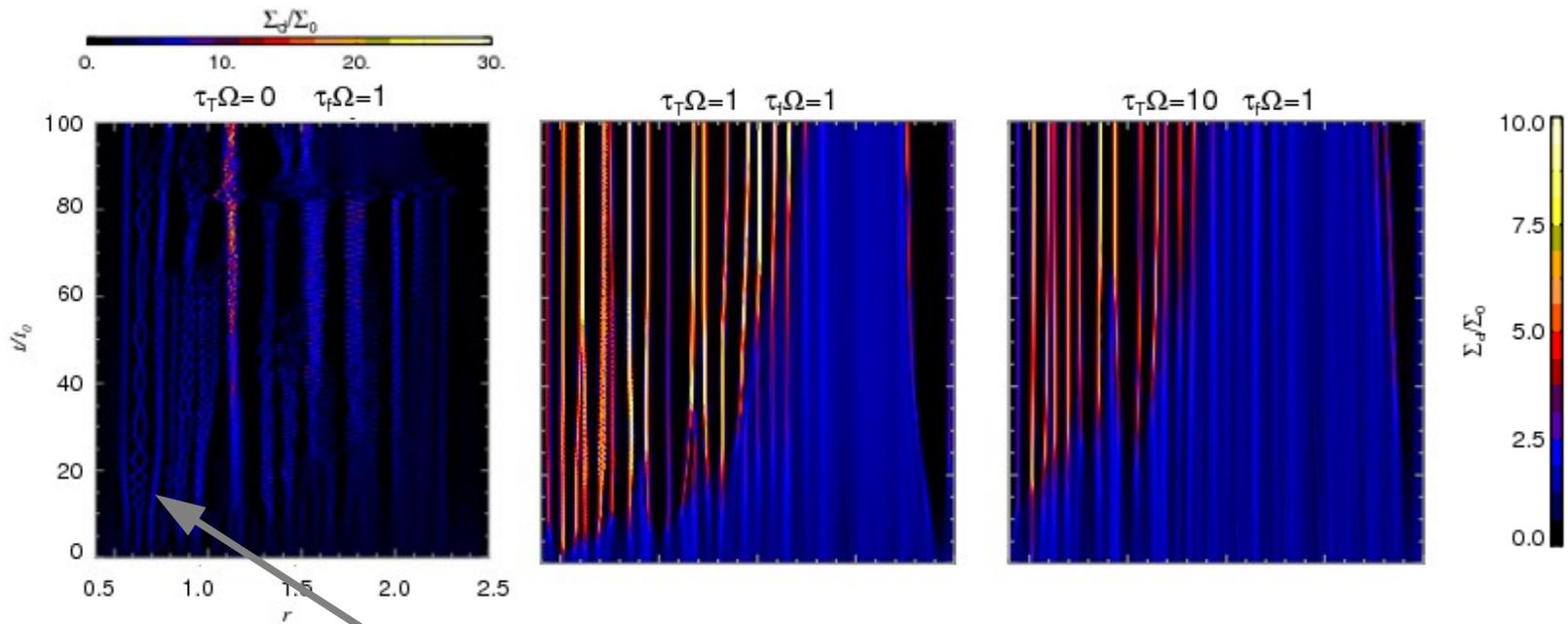
Oscillations

Thermal coupling time



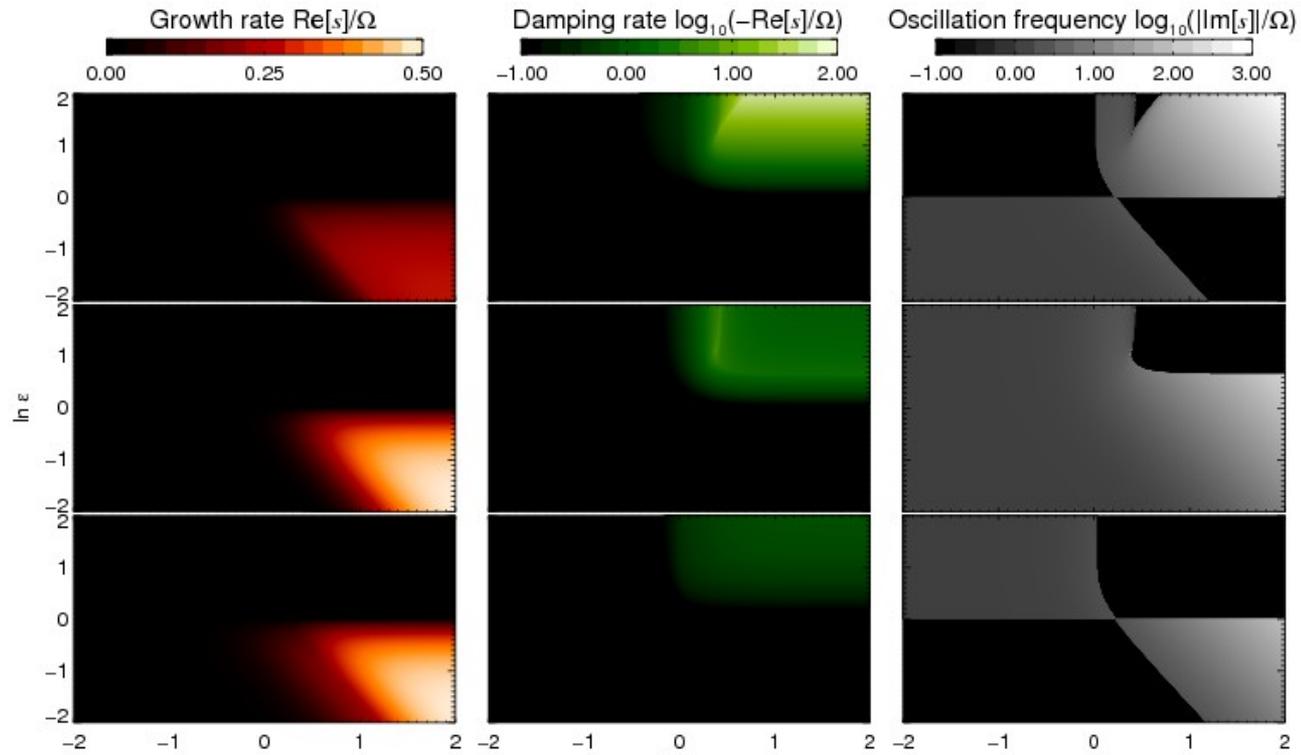
Oscillations

Thermal coupling time



Oscillations appear
with decreasing thermal time.

Solutions



Dispersion relation

$$A\omega^5 + B\omega^4 + C\omega^3 + D\omega^2 + E\omega + F = 0$$

$$A=1$$

$$B=2\epsilon + 2$$

$$C=\epsilon^2 + \epsilon(n^2+2) + 3$$

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$$\epsilon = \Sigma_d / \Sigma_g$$

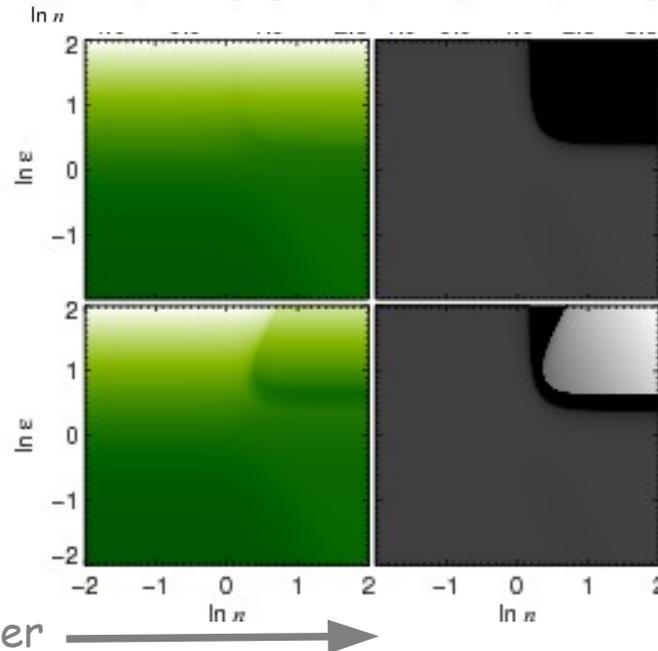
$$n = kH$$

$$\omega = s/\Omega$$

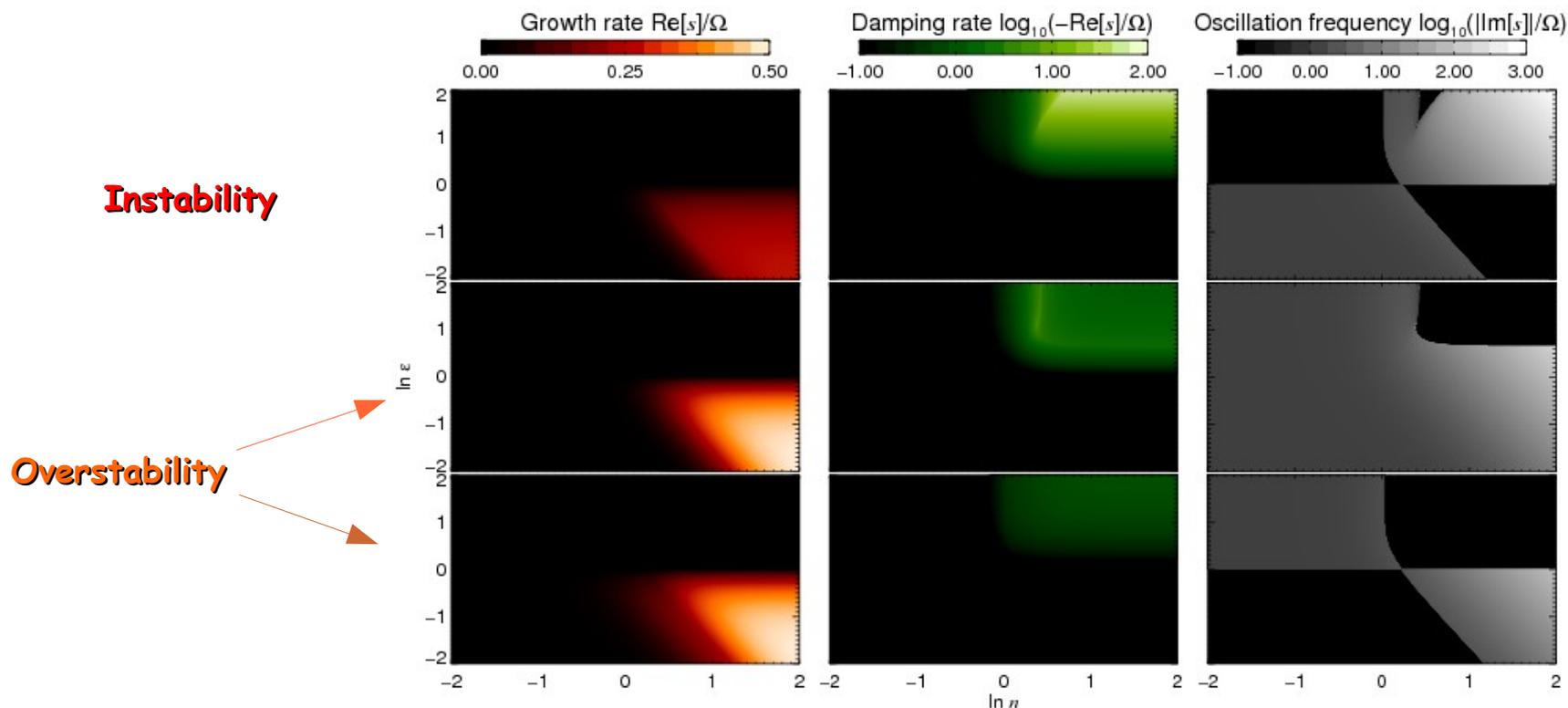
Dust-to-gas
ratio



Wavenumber



Solutions



Dispersion relation

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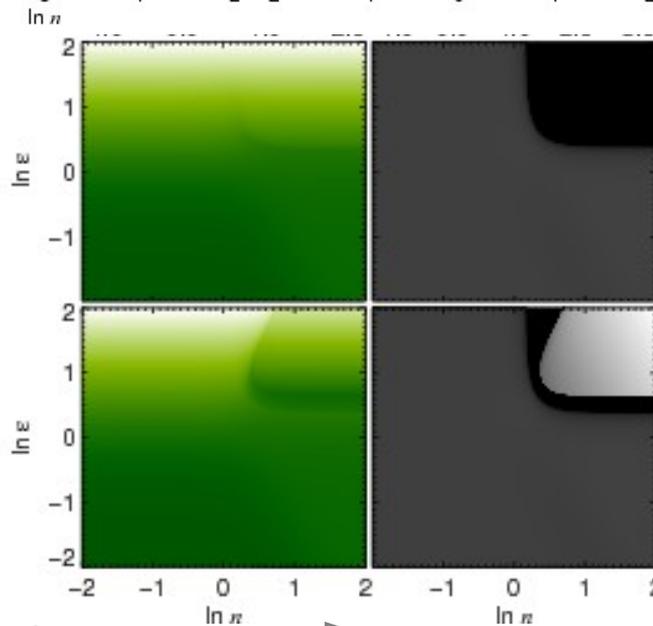
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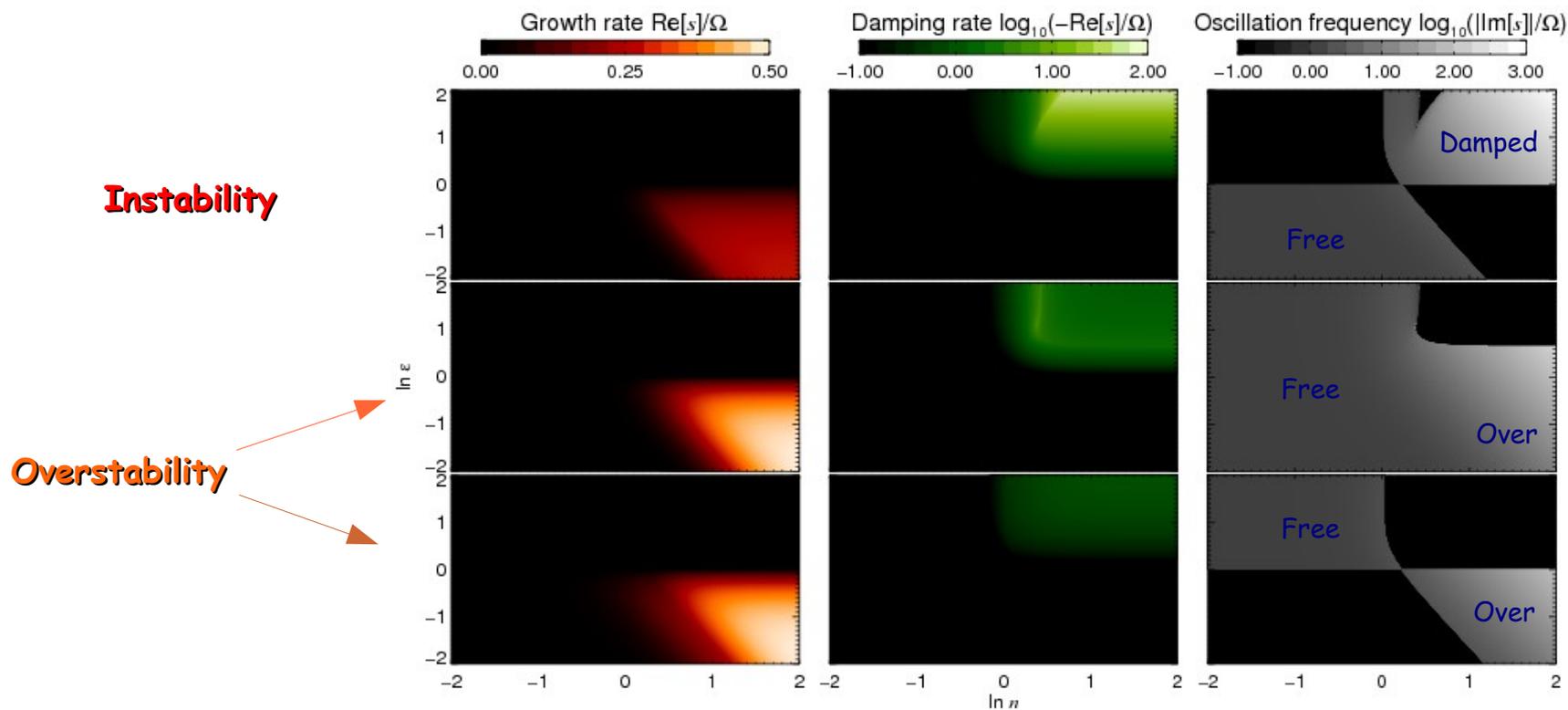
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Dust-to-gas
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Solutions



Dispersion relation

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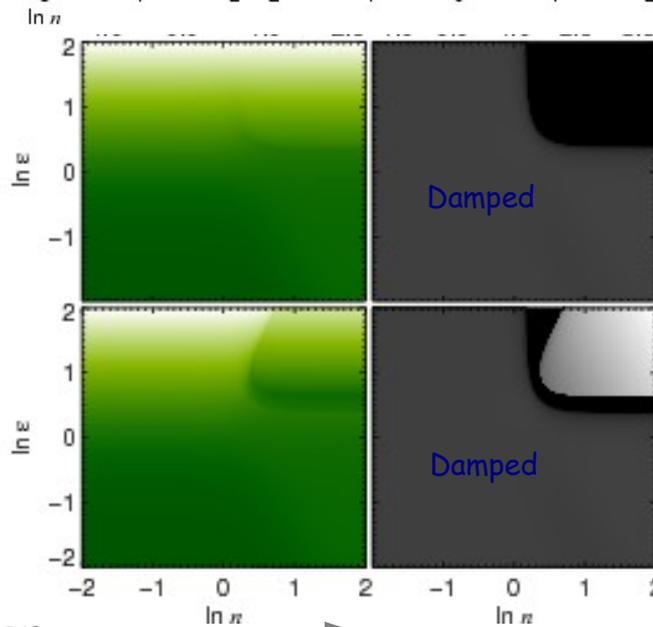
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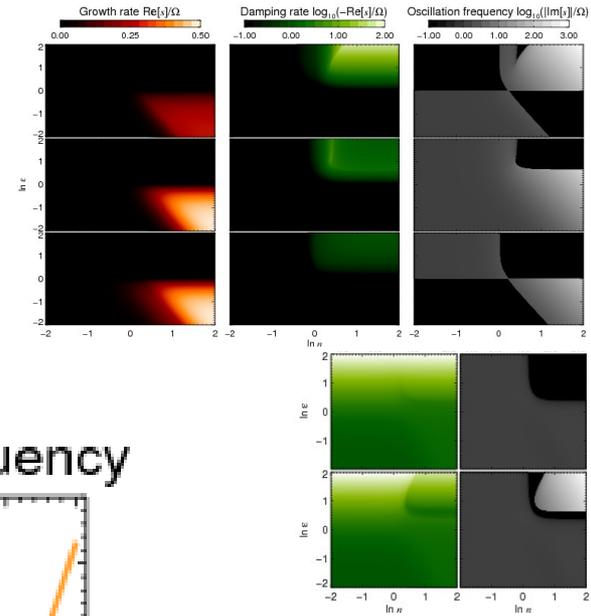
$$\omega = s/\Omega$$

Dust-to-gas
ratio

Wavenumber



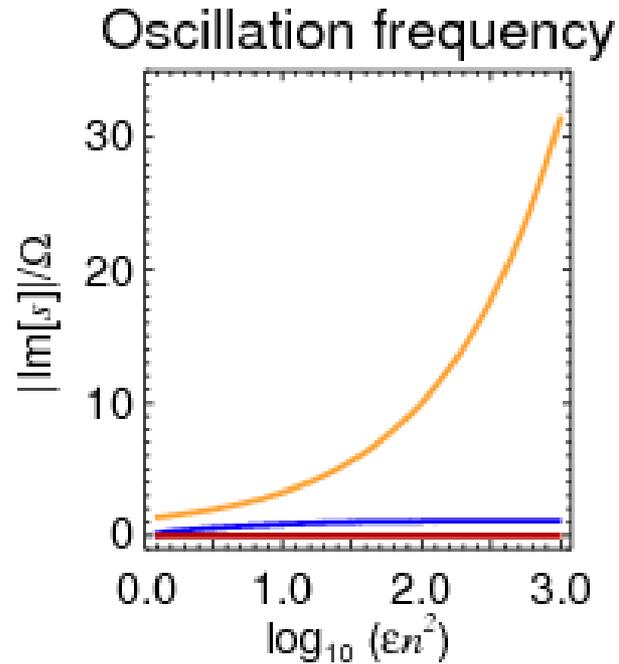
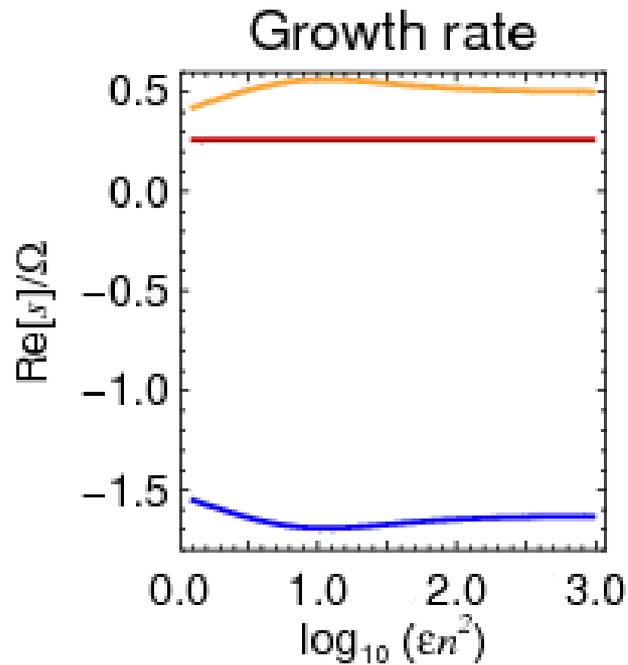
Solutions



Overstability

Instability

Oscillations

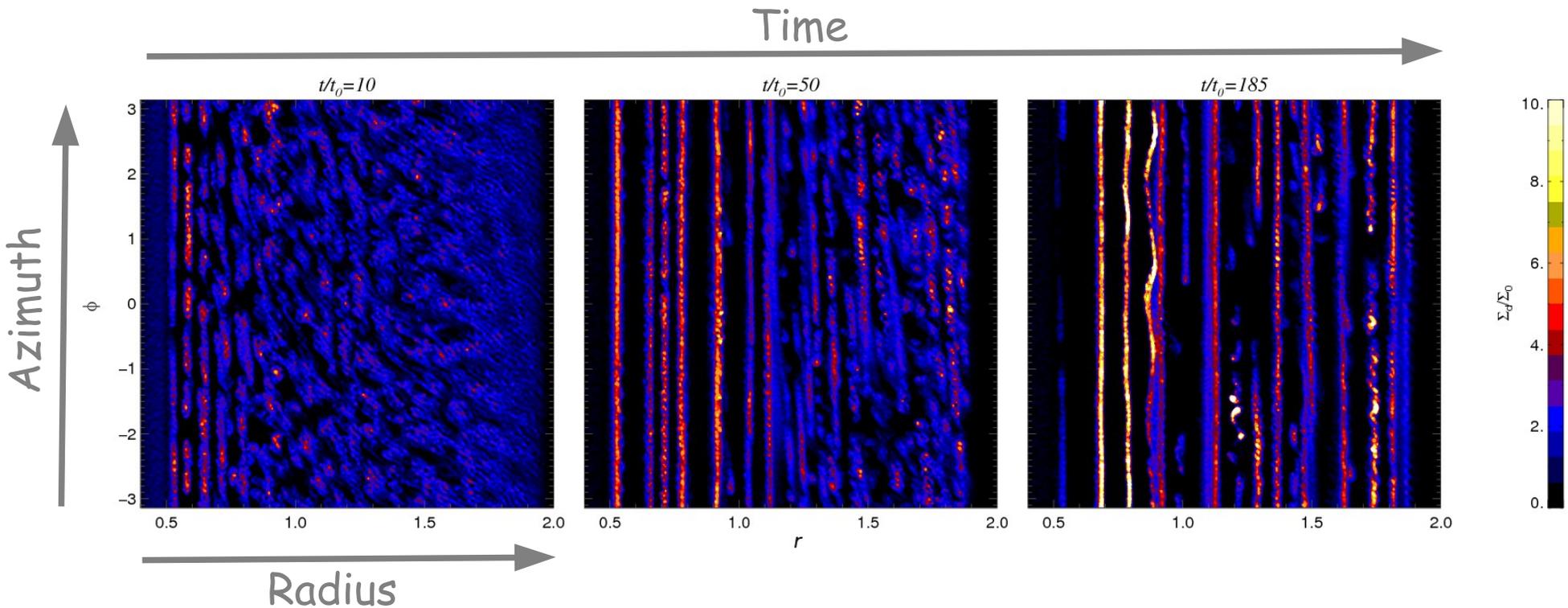


Max growth rate: $\Omega/2$.
Million-fold amplification in five orbits!

A very powerful instability.

The model in 2D: Eccentric rings

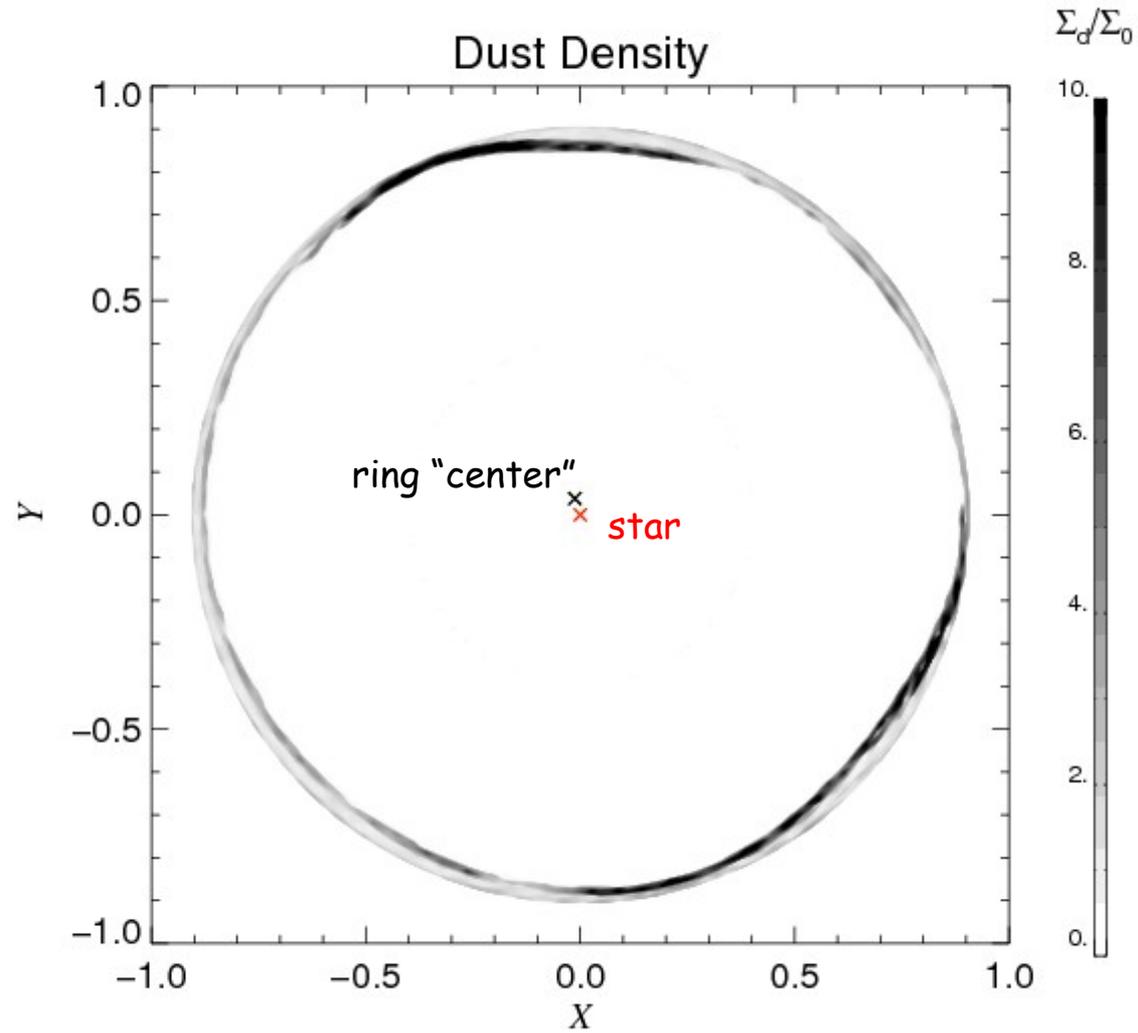
Growth of **axisymmetric** modes
+
Damping of **nonaxisymmetric** modes.
= **Rings !!!**



Epicyclic oscillations
make the ring appear **eccentric !!!**



Ring eccentricity



Eccentricity $e=0.04$

Summarizing

Gravitational collapse of an interstellar cloud

Outward transport of angular momentum through turbulence generated by the MRI. Dust coagulates into pebbles and boulders, sedimenting towards the midplane.

Rocks in the turbulent medium are trapped in transient pressure maxima and undergo collapse into planetesimals and dwarf planets.

Vortices may be excited in the dead zone. Inside them, the first dozens of Mars-mass embryos are formed. IMF ~ -2

Opacity transitions develop into regions of convergent migration. Low mass planets converge to these zones by inward/outward migration.

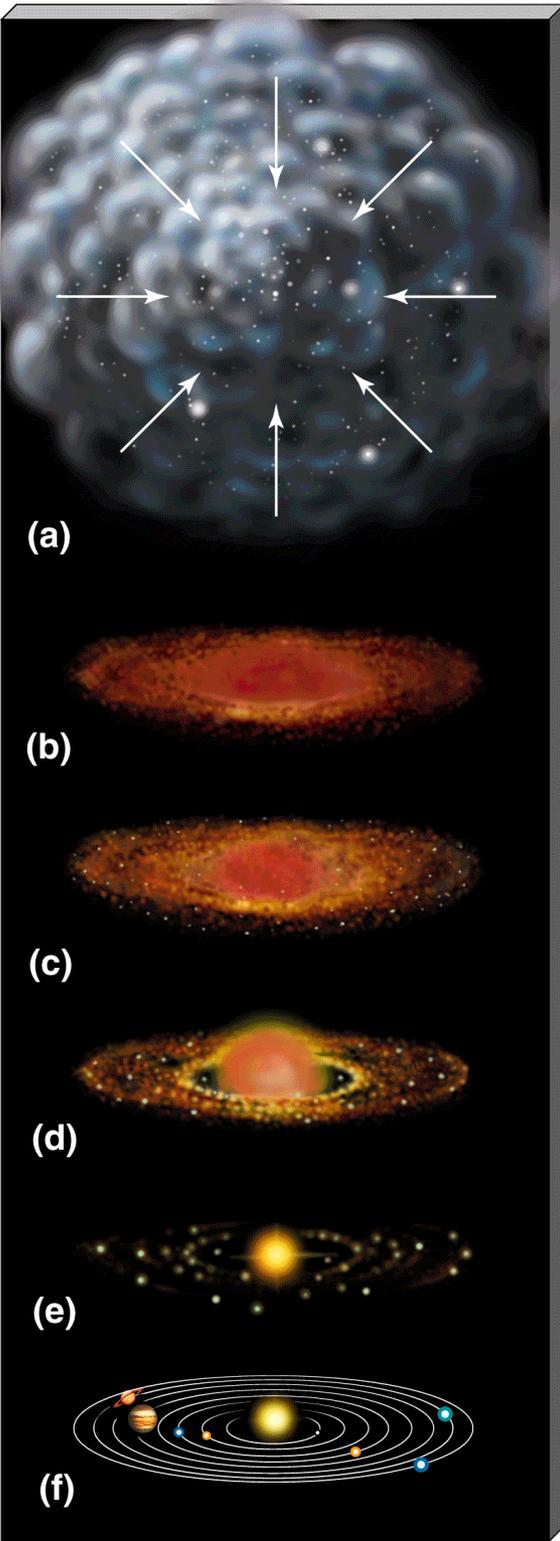
Convergent migration leads to resonances, these are disrupted by turbulent forcing. Collisions between embryos gives rise to oligarchs.

The disk thins due to photoevaporation. Planets released into stable orbits.

N-body interactions and stochastic forcing during disk evaporation produce the system's final architecture.

Debris disks with gas are subject to a thermo-centrifugal instability

The instability generates sharp eccentric rings. Caution before shouting "planet!". Not all that glitters is gold.



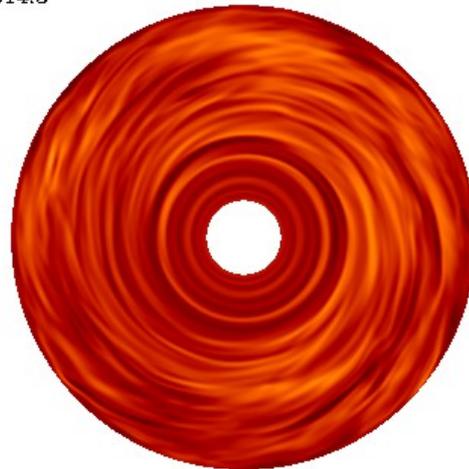
Summarizing

Gravitational collapse of an interstellar cloud

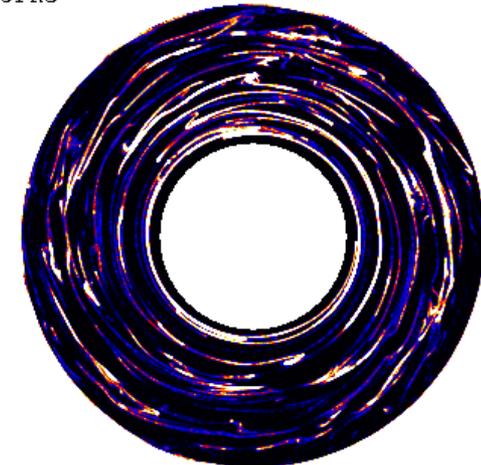
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Rocks in the turbulent medium are trapped in transient pressure maxima and undergo collapse into planetesimals and dwarf planets. **Gas** **Solids**

Vorticity mass $t=614.3$



one. $t=614.3$



Mars-

planets

Opacity

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Convection
forcing

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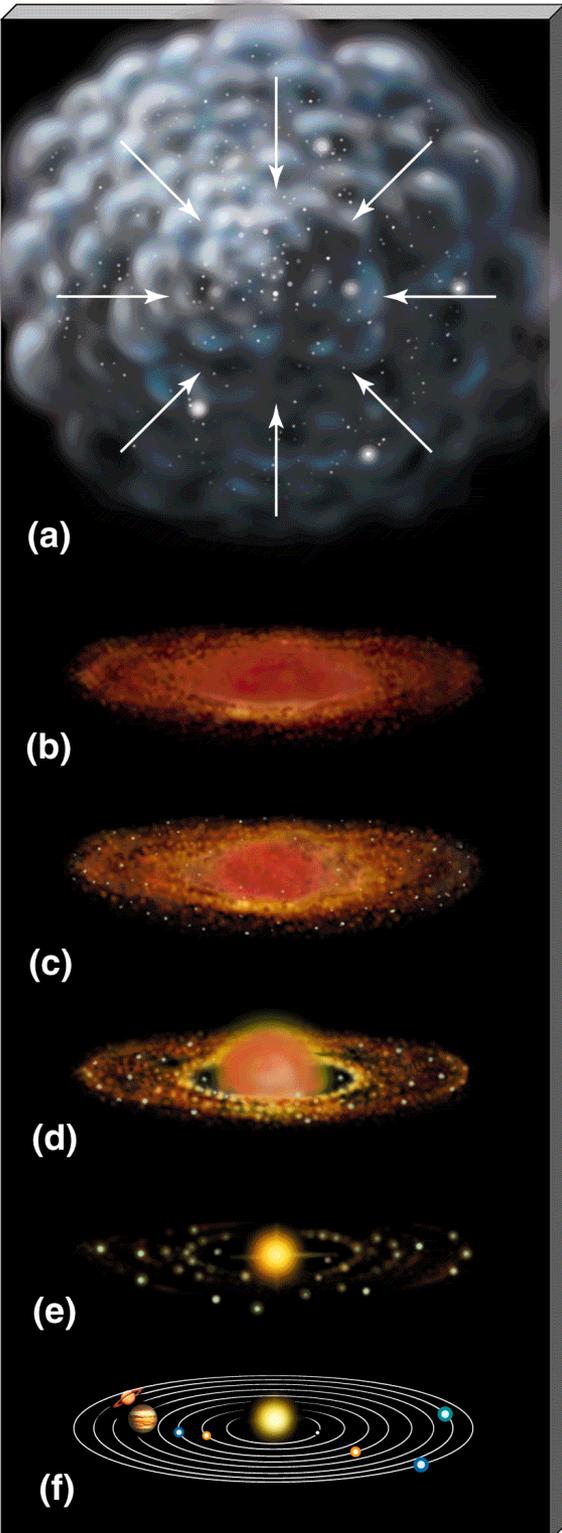
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n. Pl

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Opacity transition converges to the

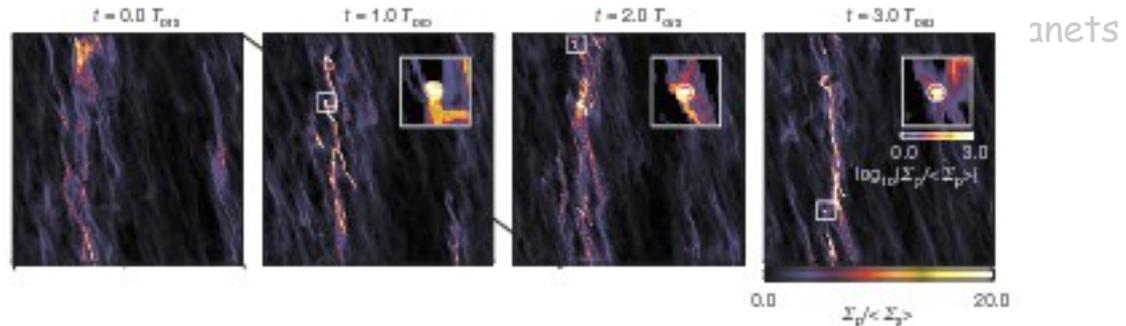
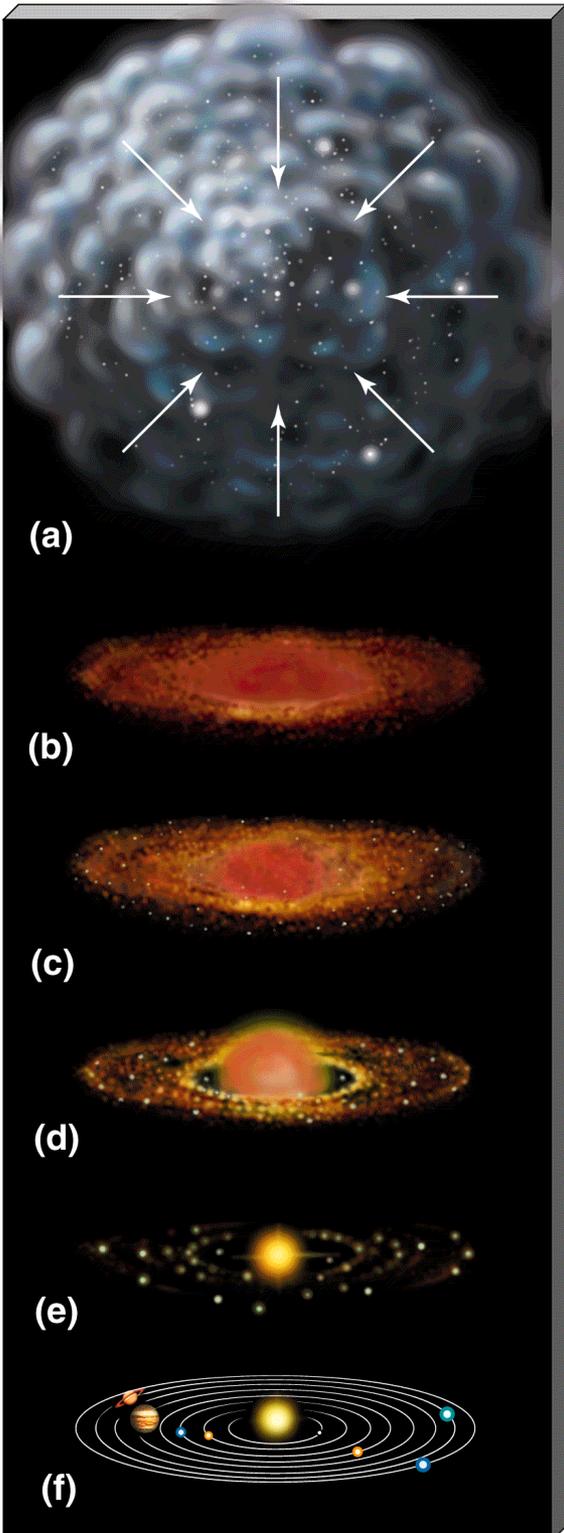
Convergent migration forcing. Collisions

The disk thins due

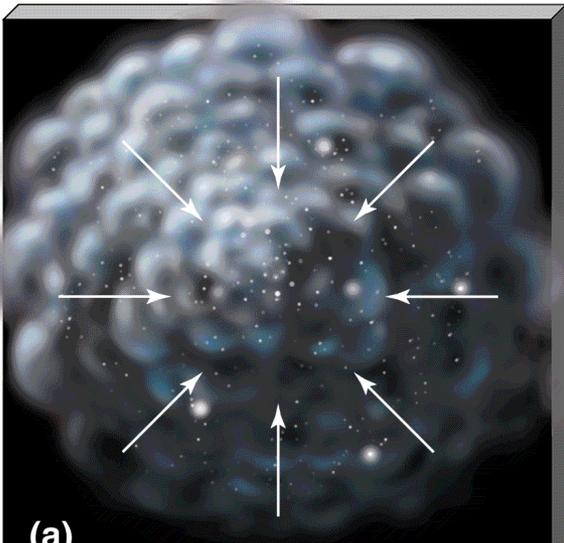
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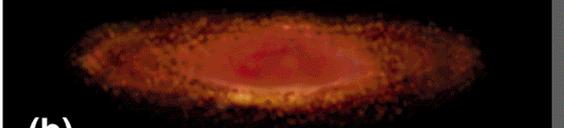
The instability generates sharp eccentric rings. Caution before shouting "planet!". Not all that glitters is gold.



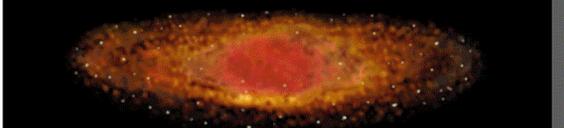
Summarizing



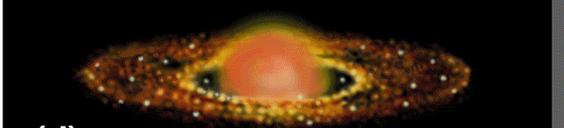
(a)



(b)



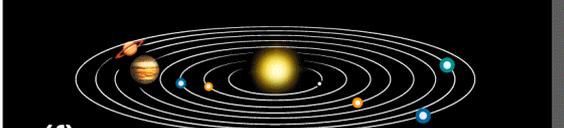
(c)



(d)



(e)

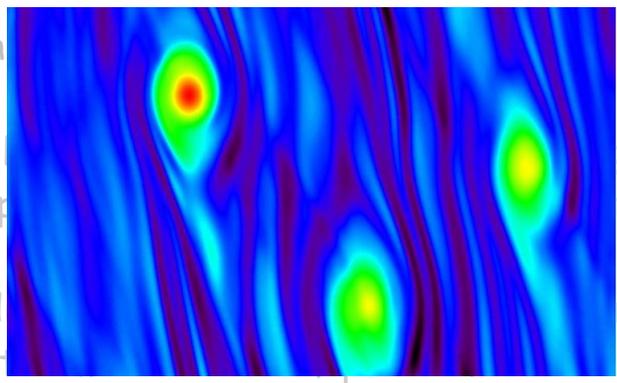


(f)

Gravitational collapse of an

Outward transport of angular momentum is generated by the MRI. Dust coagulates into pebbles and migrates towards the midplane.

Rocks in the turbulent medium are trapped at vortices (density maxima) and undergo collapse into planet embryos.



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are maxima and

Vortices may be excited in the dead zone. Inside them, the first dozens of Mars-mass embryos are formed. IMF ~ -2

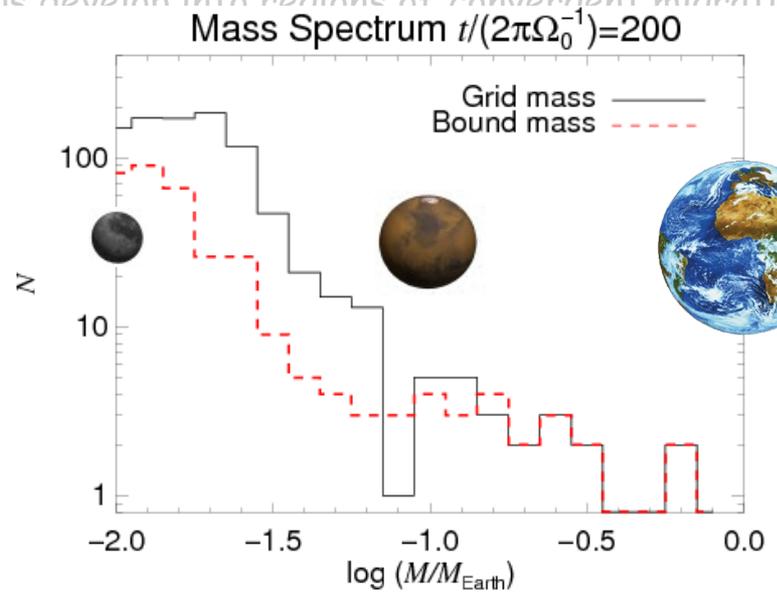
Opacity transitions develop into regions of convergent migration. Low mass planets converge to the

Convergent migration is forced by turbulent

The disk thins

N-body interactions in the system's final

Debris disks with



by turbulent

stable orbits.

Migration produce the

ability

The instability generates sharp eccentric rings. Caution before shouting "planet!". Not all that glitters is gold.

Summarizing

Gravitational collapse of

Outward transport of an MRI. Dust coagulates into

Rocks in the turbulent medium undergo collapse into planets

Vortices may be excited mass embryos are formed

Opacity transitions develop into regions of convergent migration. Low mass planets converge to these zones by inward/outward migration.

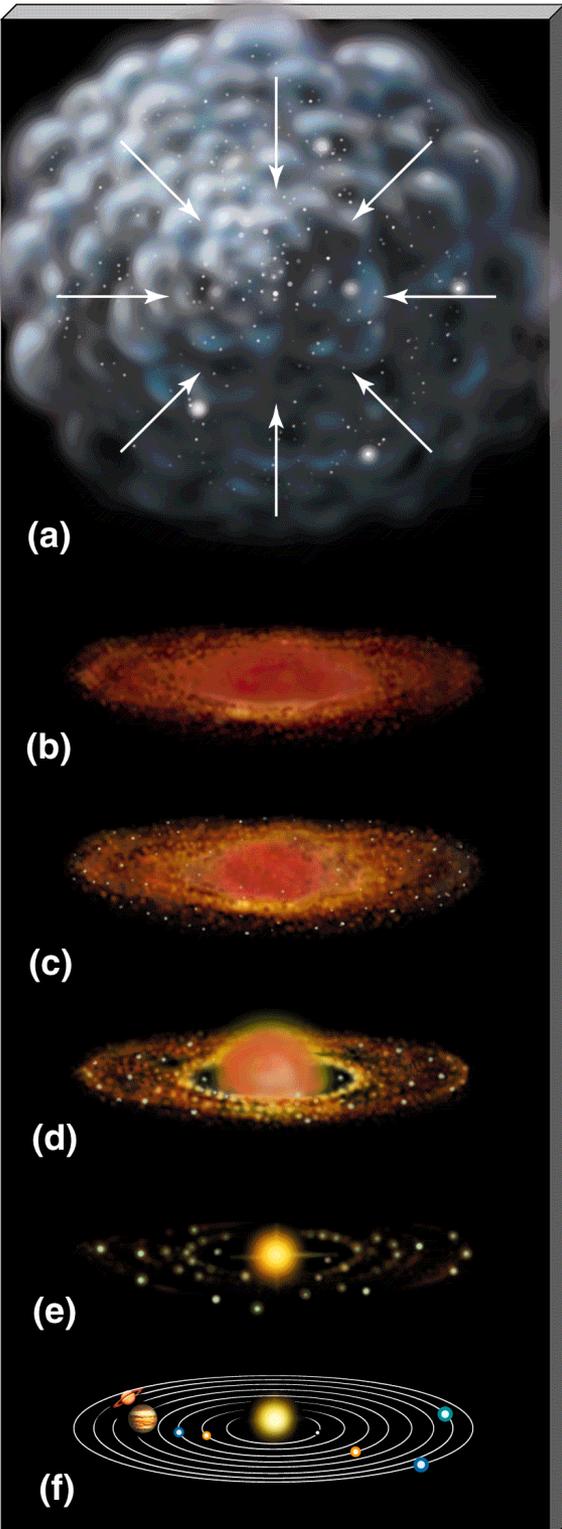
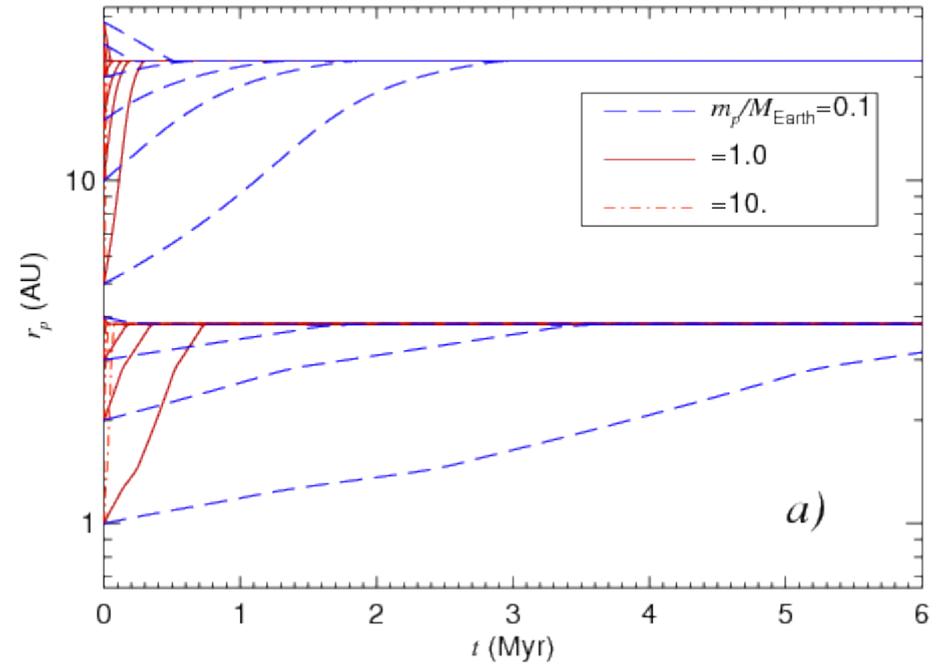
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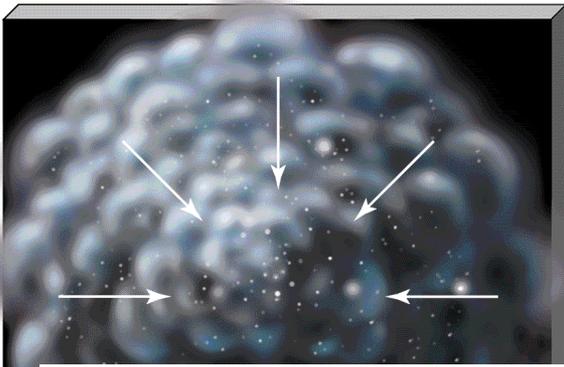
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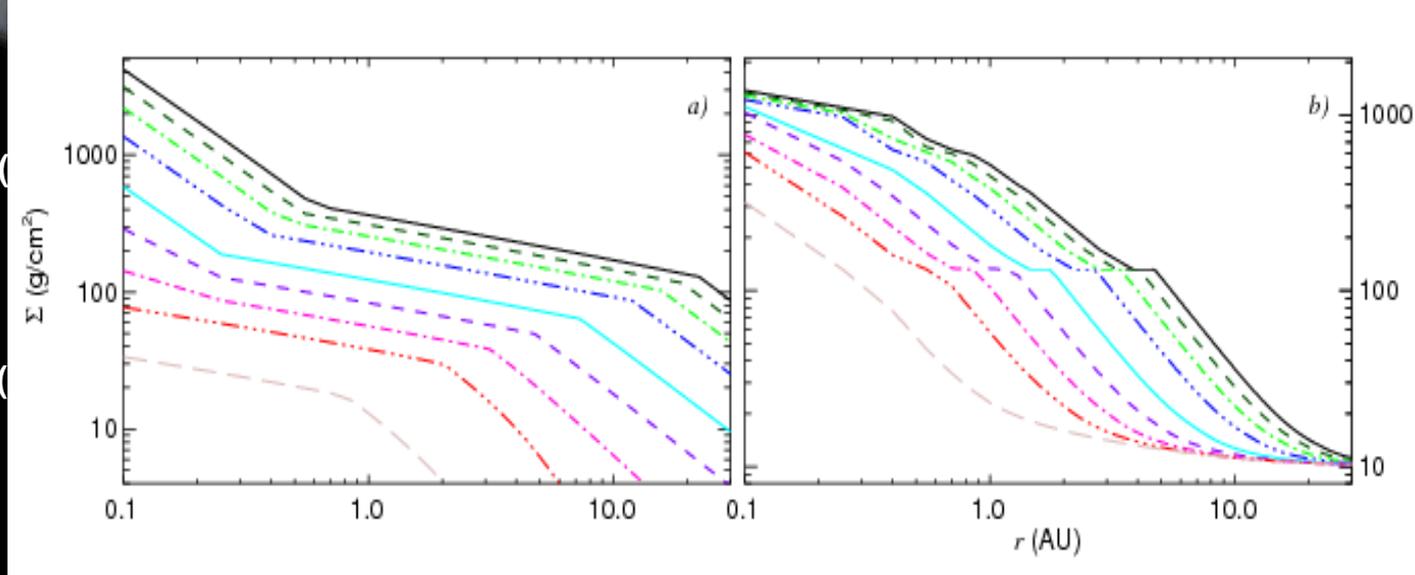


Summarizing



Gravitational collapse of an interstellar cloud

Outward transport of angular momentum through turbulence and MRI. Dust coagulates into pebbles and boulders, sediment



planet migration.

planet migration.

planet migration.

planet migration.

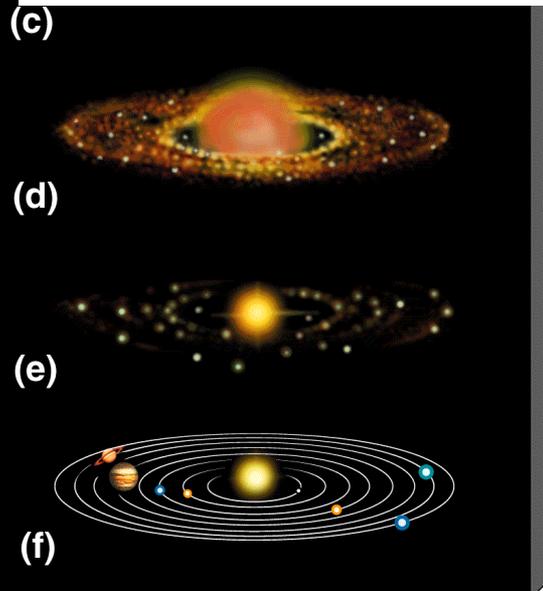
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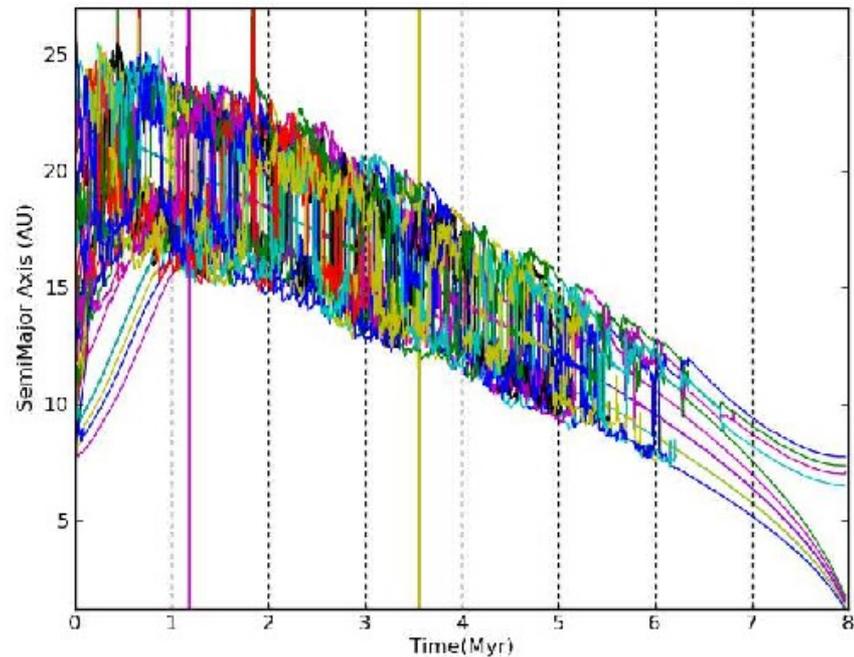
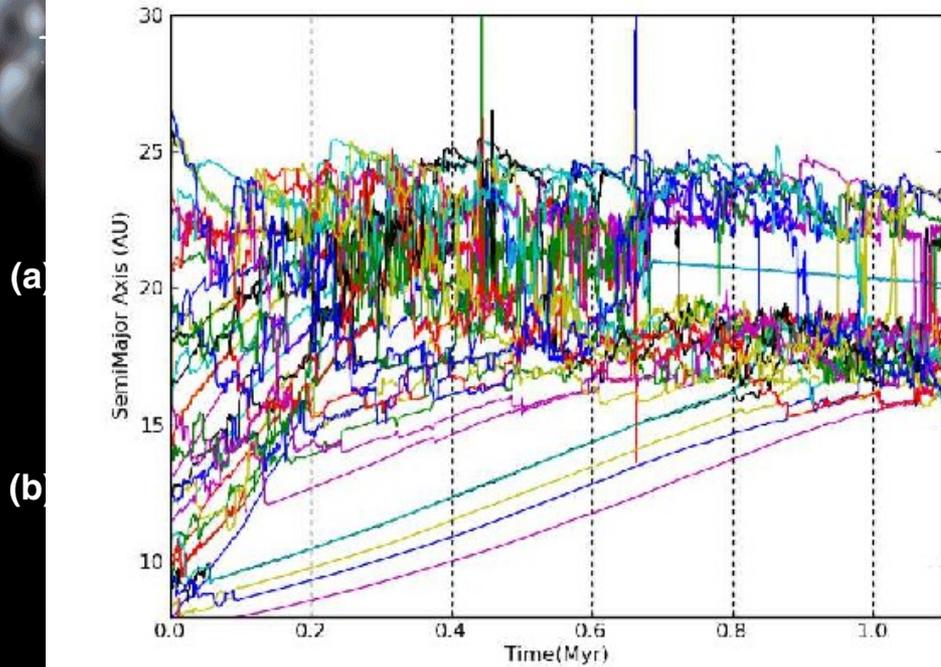
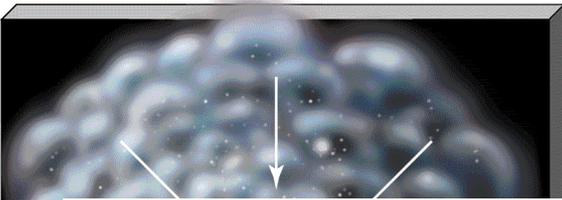
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Summarizing

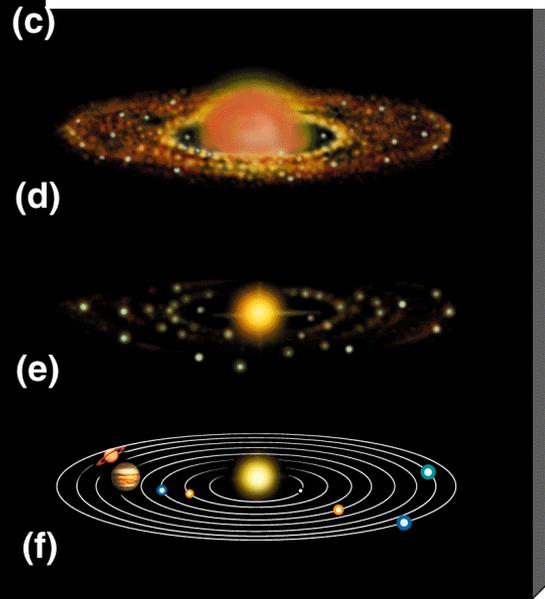
Gravitational collapse of an interstellar cloud



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plane.

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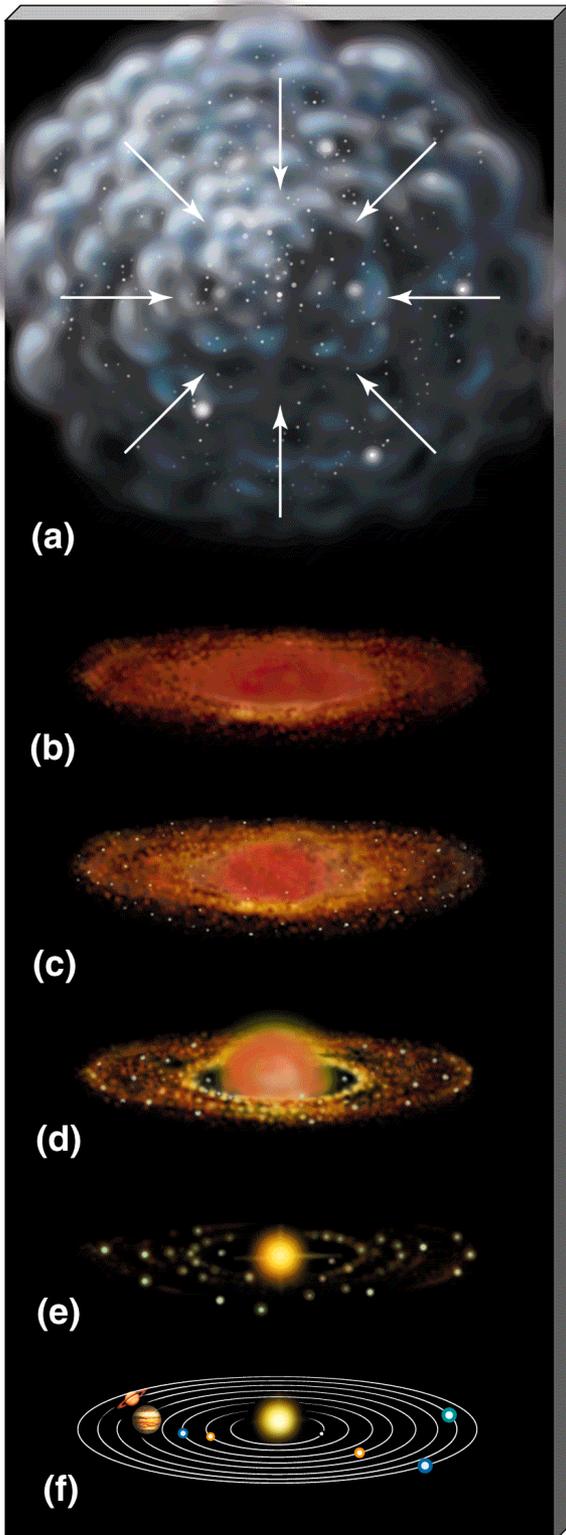
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Summariz

Gravitational collapse

Outward transport
MRI. Dust coagulation

Rocks in the turbulent
undergo collapse in

Vortices may be ex
mass embryos are

Opacity transitions develop into regions of convergence
Dust heats gas
Heated gas = high pressure region
High pressure concentrates dust

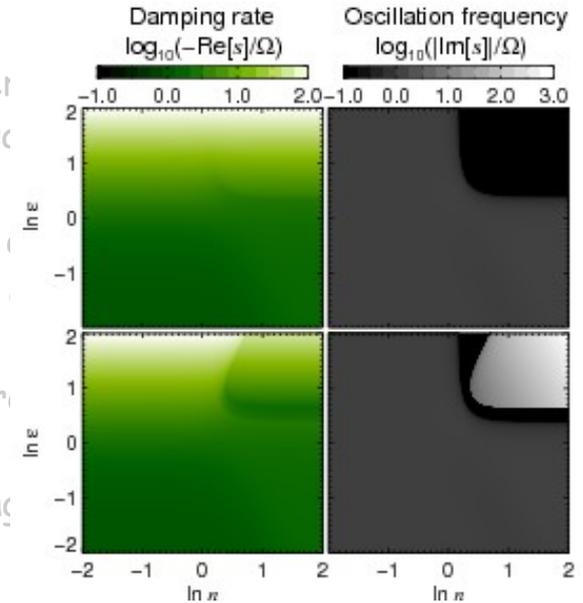
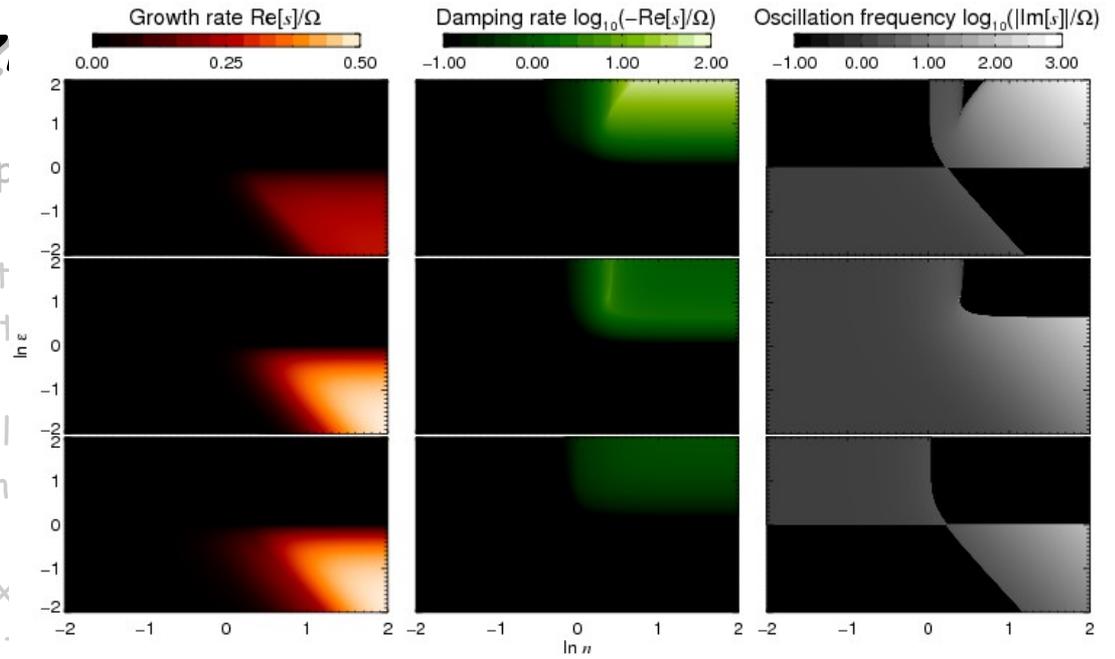
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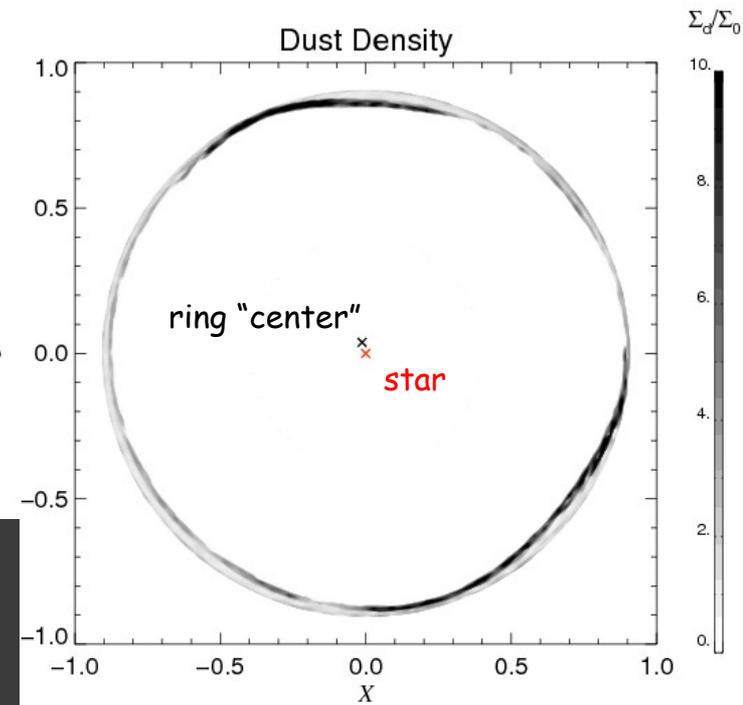
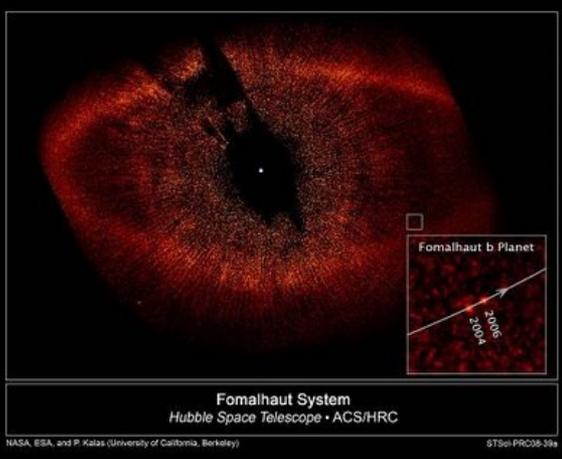


Summarizing

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Outward transport of angular momentum
MRI. Dust coagulates into pebbles and boulders

Rocks in the turbulent medium are trapped



convergent migration. Low mass planets
and migration.

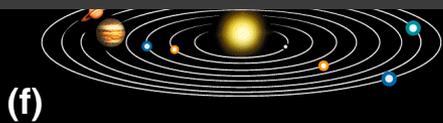
these are disrupted by turbulent
noise to oligarchs.

planets released into stable orbits.

during disk evaporation produce the

thermo-centrifugal instability

The instability generates sharp eccentric rings. **Caution before shouting "planet!"**.
Not all that glitters is gold.



$t/T_0=10$

$t/T_0=100$

$t/T_0=250$

Gas Density

Dust Density

