Physics of Accretion Disks



Wladimir Lyra

California State University Jet Propulsion Laboratory









UPPSALA UNIVERSITET











Keplerian disks can't develop hydrodynamical turbulence

Most astrophysical discs are close to Keplerian

- nearly circular
- angular velocity profile $\Omega \propto r^{-3/2}$

• angular momentum $(r^2 \Omega)$ increases outwards

Stable to axis-symmetric disturbances (Rayleigh criterion)



Rayleigh instability



Normally, angular velocity decreases outwards but angular momentum increases outwards.

Ring A losing angular momentum needs to jump to an orbit of lower angular momentum, which is inwards.

But IF *both* angular velocity and momentum decrease outwards, then when ring A loses angular, the orbit of lower angular momentum it must jump to is past B.

Likewise, B gaining angular momentum must jump to a higher angular momentum orbit, which is past A.

A and B must swap. The situation is UNSTABLE.

Gas parcels joined by a spring



The restoring force destabilizes the flow

Otherwise Rayleigh-stable, Keplerian motion is destabilized if a restoring force connects two gas parcels.



The restoring force resists the shear, trying to enforce rigid rotation.

The parcel inwards is tugged back, which enforces it to rotate at the angular velocity it had at the equilibrium position.

The angular momentum it loses is given to the outer gas parcel. Losing angular momentum, it jumps to an inner orbit.

The outer parcel gains angular momentum and jumps to an even more outward orbit.

The situation is unstable.

Magnetorotational Instability (MRI)



Magnetic fields in a conducting, rotating plasma behave EXACTLY like springs!

Turbulence and Accretion in 3D Global MHD Simulations of Stratified Protoplanetary Disk

Magneto-Rotational Instability

Turbulence in disks is enabled by the Magneto-Rotational Instability (Balbus & Hawley, 1991)



Magneto-Rotational Instability



Saturated State of MRI

Energy budget



FIG. 6.—Sketch of the energy budget. Energy is tapped from the Keplerian motion and goes into magnetic and kinetic energy, and is finally converted into heat. The numbers give the approximate energy fluxes in units of $\langle \frac{1}{2}B^2\Omega \rangle$.

$$\frac{d}{dt} \left(\frac{1}{2\mu_0} B^2 \right) = -\frac{3}{2} \Omega_0 \frac{1}{\mu_0} B_x B_y - \boldsymbol{u} \cdot (\boldsymbol{J} \times \boldsymbol{B}) - \eta \mu_0 J^2$$
$$\frac{d}{dt} \left(\frac{1}{2} \rho \boldsymbol{u}^2 \right) = \frac{3}{2} \Omega_0 \rho \boldsymbol{u}_x \boldsymbol{u}_y + \rho \boldsymbol{u} \cdot \boldsymbol{g} - \boldsymbol{u} \cdot \nabla \boldsymbol{p} + \boldsymbol{u} \cdot (\boldsymbol{J} \times \boldsymbol{B}) - 2 \nu \rho S^2$$
$$\frac{d}{dt} \rho \boldsymbol{e} = -p \nabla \cdot \boldsymbol{u} + 2 \nu \rho S^2 + \eta \mu_0 J^2 + \rho Q$$
$$\frac{d}{dt} E_{tot} = \frac{3}{2} \Omega_0 \left(\rho \boldsymbol{u}_x \boldsymbol{u}_y - \frac{1}{\mu_0} B_x B_y \right) + \rho Q$$

Dead zones





Next class

The MRI is dead Long live the *Thermal Instabilities*



Vertical shear instability



$$\Omega = \Omega_{\rm K} \left[1 + \frac{1}{2} \left(\frac{H}{R} \right)^2 \left(p + q + \frac{q}{2} \frac{Z^2}{H^2} \right) \right]$$

$$d\Omega/dz \neq 0$$
; $\kappa_z^2 < 0$

Violates the Rayleigh criterion: Unstable



Nelson et al. (2013)