

Class 14 – Marth, 2020

Sub-Neptune Formation: The View from Resonant Planets

Nick Choksi, Eugene Chiang

(Submitted on 6 Mar 2020)

The orbital period ratios of neighbouring sub-Neptunes are distributed asymmetrically near first-order resonances. There are deficits of systems---"troughs" in the period ratio histogram---just short of commensurability, and excesses---"peaks"---just wide of it. We reproduce quantitatively the strongest peak-trough asymmetries, near the 3:2 and 2:1 resonances, using dissipative interactions between planets and their natal disks. Disk eccentricity damping captures bodies into resonance and clears the trough, and when combined with disk-driven convergent migration, draws planets initially wide of commensurability into the peak. The migration implied by the magnitude of the peak is modest; reductions in orbital period are $\sim 10\%$, supporting the view that sub-Neptunes complete their formation more--or--less in situ. Once captured into resonance, sub-Neptunes of typical mass $\sim 5\text{--}15 M_{\oplus}$ stay captured (contrary to an earlier claim), as they are immune to the overstability that afflicts lower mass planets. Driving the limited, short-scale migration is a gas disk whose surface density is fairly constant inside 1 AU and depleted relative to a solar-composition disk by $3\text{--}5$ orders of magnitude. Such gas-poor but not gas-empty environments are quantitatively consistent with sub-Neptune core formation by giant impacts (and not, e.g., pebble accretion). While disk-planet interactions at the close of the planet formation era adequately explain the 3:2 and 2:1 asymmetries at periods $\gtrsim 5\text{--}15$ days, supporting modification by stellar tides appears necessary at shorter periods, particularly for the 2:1.

Comments: Submitted to MNRAS
Subjects: **Earth and Planetary Astrophysics (astro-ph.EP)**
Cite as: **arXiv:2003.03388** [astro-ph.EP]
(or **arXiv:2003.03388v1** [astro-ph.EP] for this version)

Most Stars (and Planets?) Are Born in Intense Radiation Fields

Eve J. Lee, Philip F. Hopkins

(Submitted on 6 Mar 2020)

Protostars and young stars are strongly spatially "clustered" or "correlated" within their natal giant molecular clouds (GMCs). We demonstrate that such clustering leads to the conclusion that the incident bolometric radiative flux upon a random young star/disc is enhanced (relative to volume-averaged fluxes) by a factor which increases with the total stellar mass of the complex. Because the Galactic cloud mass function is top-heavy, the typical star in our Galaxy experienced a much stronger radiative environment than those forming in well-observed nearby (but relatively small) clouds, exceeding fluxes in the Orion Nebular Cluster by factors of $\gtrsim 30$. Heating of the circumstellar disc around a median young star is dominated by this external radiation beyond ~ 50 AU. And if discs are not well-shielded by ambient dust, external UV irradiation can dominate over the host star down to sub-AU scales. Another consequence of stellar clustering is an extremely broad Galaxy-wide distribution of incident flux (spanning > 10 decades), with half the Galactic star formation in a substantial "tail" towards even more intense background radiation. We also show that the strength of external irradiation is amplified super-linearly in high-density environments such as the Galactic centre, starbursts, or high-redshift galaxies.

Comments: MNRAS Letters in press
Subjects: **Solar and Stellar Astrophysics (astro-ph.SR)**; Earth and Planetary Astrophysics (astro-ph.EP); Astrophysics of Galaxies (astro-ph.GA)
Cite as: **arXiv:2003.03390** [astro-ph.SR]
(or **arXiv:2003.03390v1** [astro-ph.SR] for this version)

Anibal Sierra, Susana Lizano

(Submitted on 6 Mar 2020)

It is known that the millimeter dust thermal emission of protoplanetary disks is affected by scattering, such that for optically thick disks the emission decreases with respect to the pure absorption case and the spectral indices can reach values below 2. The latter can also be obtained with temperature gradients. Using simple analytical models of radiative transfer in thin slabs, we quantify the effect of scattering, vertical temperature gradients, and dust settling on the emission and spectral indices of geometrically thin face-on accretion disks around young stars. We find that in vertically isothermal disks with large albedo ($\omega_{\nu} \gtrsim 0.6$), the emergent intensity can increase at optical depths between 10^{-2} and 10^{-1} . We show that dust settling has important effects on the spectral indices in the optically thick regime, since the disk emission mainly traces small dust grains in the upper layers of the disk. The $\lambda = 870 \mu\text{m}$ emission of these small grains can hide large grains at the disk mid plane when the dust surface density is larger than $\sim 3.21 \text{ g cm}^{-2}$. Finally, because of the change of the shape of the spectral energy distribution, optically thick disks at 1.3 mm and grains with sizes between $300 \mu\text{m} < a_{\text{max}} < 1 \text{ mm}$ have a 7 mm flux $\sim 60\%$ higher than the extrapolation from higher millimeter frequencies, assumed when scattering is neglected. This effect could provide an explanation to the excess emission at $\lambda = 7 \text{ mm}$ reported in several disks.

Subjects: **Earth and Planetary Astrophysics (astro-ph.EP)**; Solar and Stellar Astrophysics (astro-ph.SR)
Cite as: **arXiv:2003.02982** [astro-ph.EP]
(or **arXiv:2003.02982v1** [astro-ph.EP] for this version)

Global 3-D Radiation Magnetohydrodynamic Simulations for FU Ori's Accretion Disk and Observational Signatures of Magnetic Fields

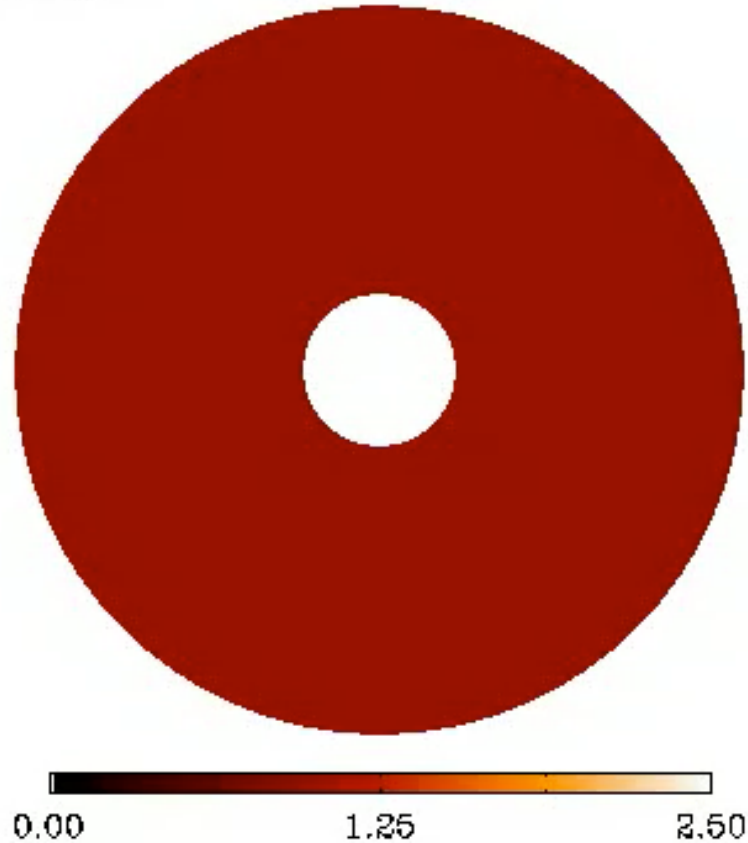
Zhaohuan Zhu, Yan-Fei Jiang, James M. Stone

(Submitted on 3 Dec 2019 (v1), last revised 6 Mar 2020 (this version, v2))

FU Ori is the prototype of FU Orionis systems which are outbursting protoplanetary disks. Magnetic fields in FU Ori's accretion disks have previously been detected using spectropolarimetry observations for Zeeman effects. We carry out global radiation ideal MHD simulations to study FU Ori's inner accretion disk. We find that (1) when the disk is threaded by vertical magnetic fields, most accretion occurs in the magnetically dominated atmosphere at $z \sim R$, similar to the "surface accretion" mechanism in previous locally-isothermal MHD simulations. (2) A moderate disk wind is launched in the vertical field simulations with a terminal speed of $\sim 300\text{--}500 \text{ km/s}$ and a mass loss rate of $1\text{--}10\%$ the disk accretion rate, which is consistent with observations. Disk wind fails to be launched in simulations with net toroidal magnetic fields. (3) The disk photosphere at the unit optical depth can be either in the wind launching region or the accreting surface region. Magnetic fields have drastically different directions and magnitudes between these two regions. Our fiducial model agrees with previous optical Zeeman observations regarding both the field directions and magnitudes. On the other hand, simulations indicate that future Zeeman observations at near-IR wavelengths or towards other FU Orionis systems may reveal very different magnetic field structures. (4) Due to energy loss by the disk wind, the disk photosphere temperature is lower than that predicted by the thin disk theory, and the previously inferred disk accretion rate may be lower than the real accretion rate by a factor of $\sim 2\text{--}3$.

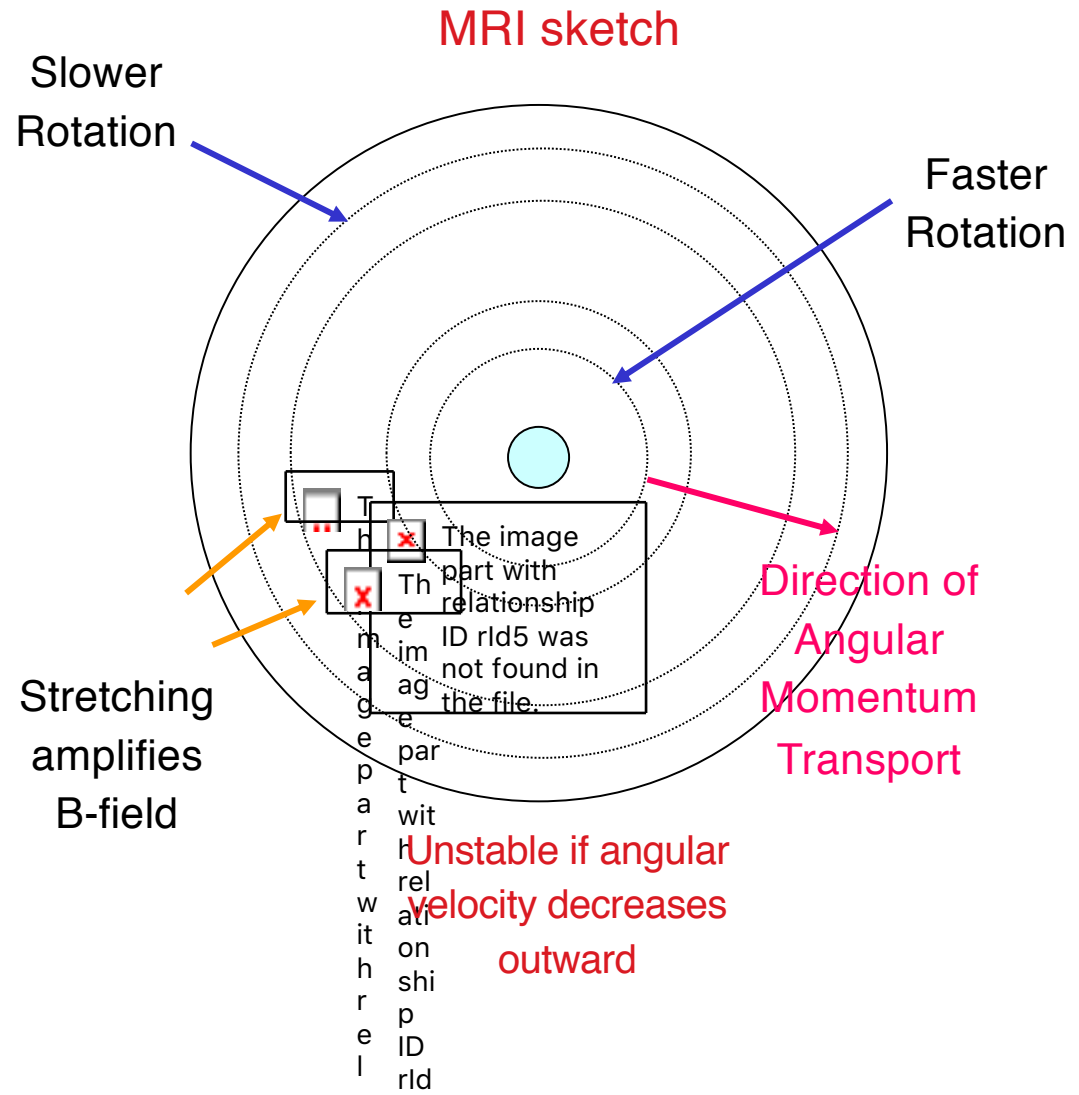
Balbus & Hawley (1991) – Magnetorotational Instability as source of turbulence

$t=0.2/88\text{yr}$



Magnetized disk

Lyra et al. (2008a)



Saturated State of MRI

Energy budget

VIEW

Abstract

Citations (672)

References (63)

Co-Reads

Similar Papers

Volume Content

Graphics

Metrics

Export Citation

Dynamo-generated Turbulence and Large-Scale Magnetic Fields in a Keplerian Shear Flow

Show affiliations

Brandenburg, Axel; Nordlund, Ake; Stein, Robert F.; Torkelsson, Ulf

The nonlinear evolution of magnetized Keplerian shear flows is simulated in a local, three-dimensional model, including the effects of compressibility and stratification. Supersonic flows are initially generated by the Balbus-Hawley magnetic shear instability. The resulting flows regenerate a turbulent magnetic field which, in turn, reinforces the turbulence. Thus, the system acts like a dynamo that generates its own turbulence. However, unlike usual dynamos, the magnetic energy exceeds the kinetic energy of the turbulence by a factor of 3-10. By assuming the field to be vertical on the outer (upper and lower) surfaces we do not constrain the horizontal magnetic flux. Indeed, a large-scale toroidal magnetic field is generated, mostly in the form of toroidal flux tubes with lengths comparable to the toroidal extent of the box. This large-scale field is mainly of even (i.e., quadrupolar) parity with respect to the midplane and changes direction on a timescale of 30 orbits, in a possibly cyclic manner. The effective Shakura-Sunyaev alpha viscosity parameter is between 0.001 and 0.005, and the contribution from the Maxwell stress is 3-7 times larger than the contribution from the Reynolds stress.

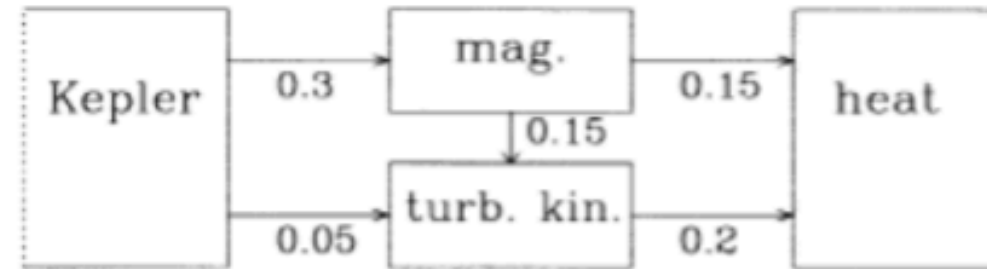
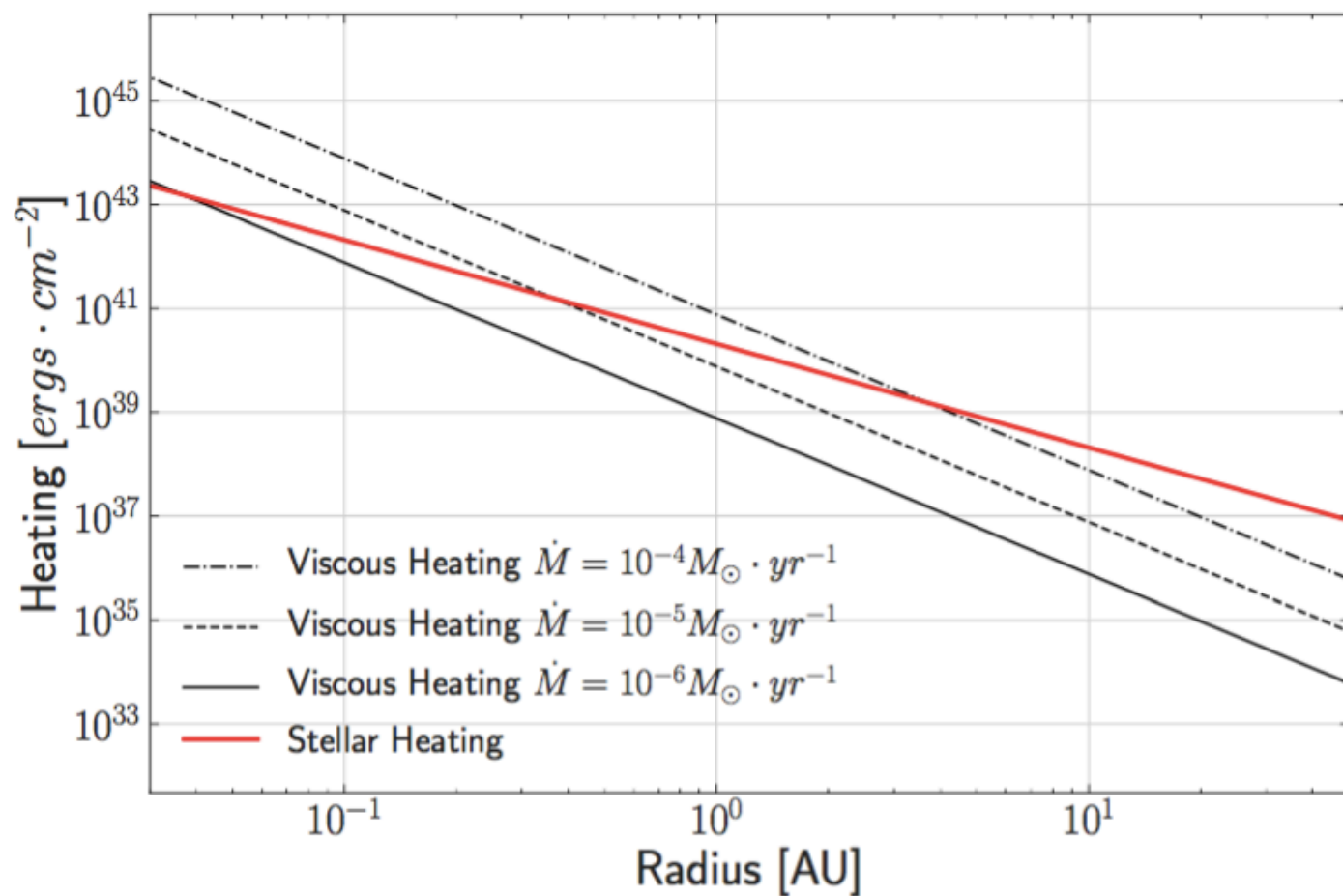


FIG. 6.—Sketch of the energy budget. Energy is tapped from the Keplerian motion and goes into magnetic and kinetic energy, and is finally converted into heat. The numbers give the approximate energy fluxes in units of $\langle \frac{1}{2} B^2 \Omega \rangle$.

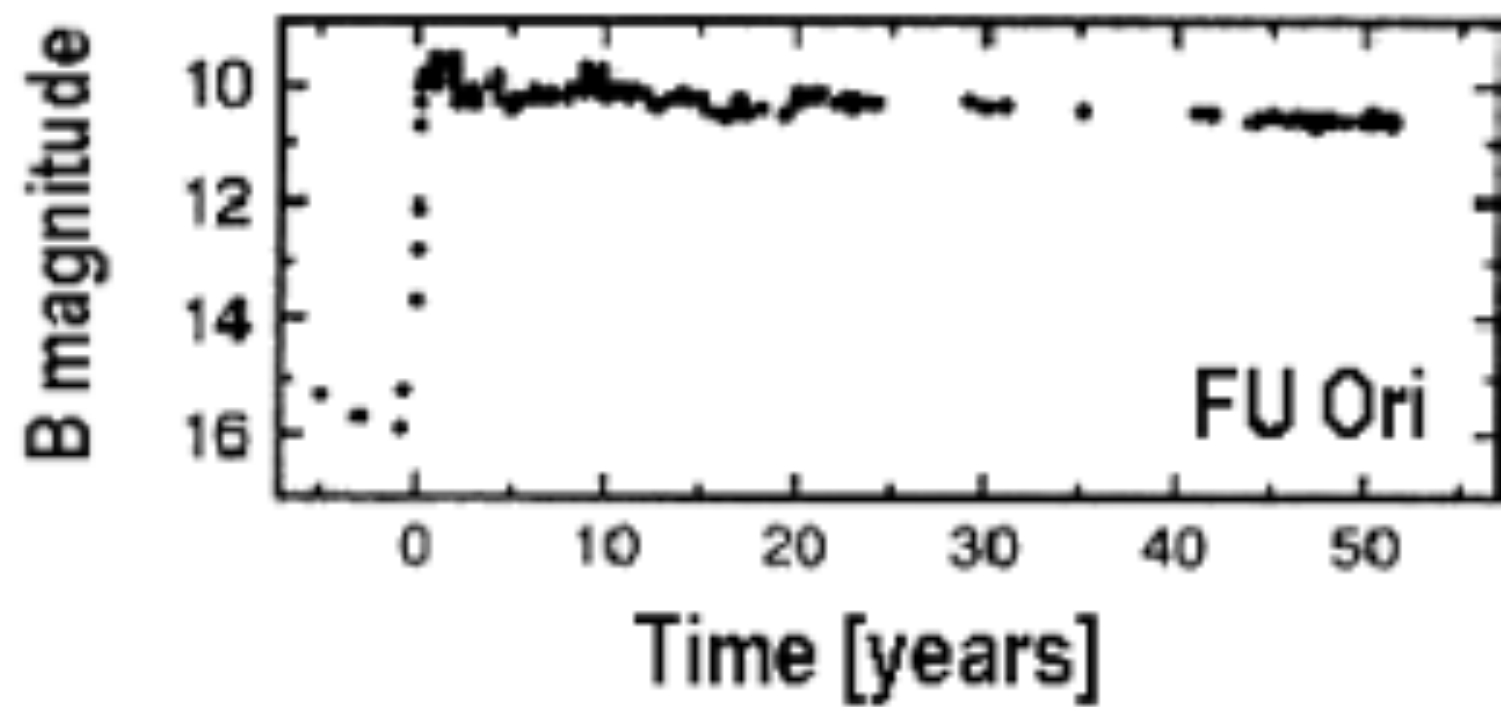
$$\begin{aligned} \frac{d}{dt} \left(\frac{1}{2 \mu_0} B^2 \right) &= -\frac{3}{2} \Omega_0 \frac{1}{\mu_0} B_x B_y - \mathbf{u} \cdot (\mathbf{J} \times \mathbf{B}) - \eta \mu_0 J^2 \\ \frac{d}{dt} \left(\frac{1}{2} \rho u^2 \right) &= \frac{3}{2} \Omega_0 \rho u_x u_y + \rho \mathbf{u} \cdot \mathbf{g} - \mathbf{u} \cdot \nabla p + \mathbf{u} \cdot (\mathbf{J} \times \mathbf{B}) - 2 \nu \rho S^2 \\ \frac{d}{dt} \rho e &= -p \nabla \cdot \mathbf{u} + 2 \nu \rho S^2 + \eta \mu_0 J^2 + \rho Q \\ \frac{d}{dt} E_{tot} &= \frac{3}{2} \Omega_0 \left(\rho u_x u_y - \frac{1}{\mu_0} B_x B_y \right) + \rho Q \end{aligned}$$

Viscous heating

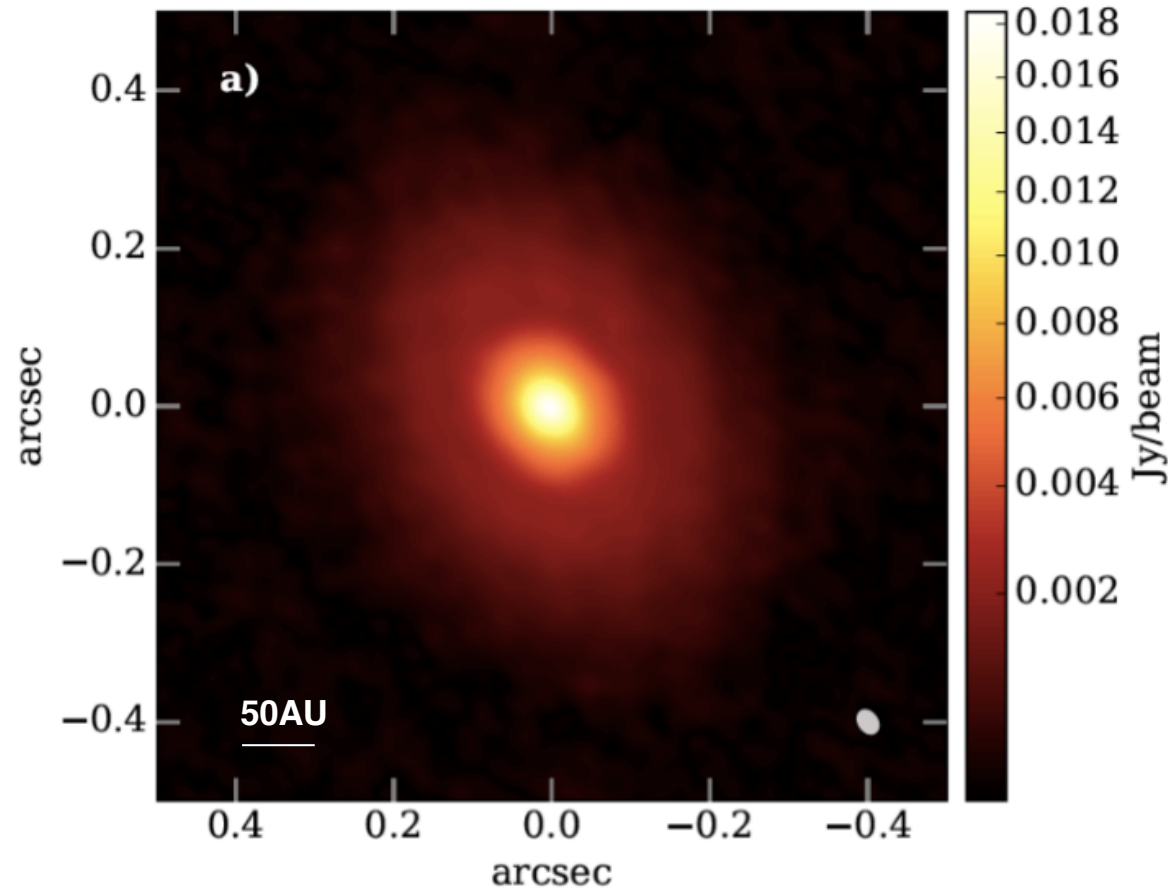
$$q^+ = \frac{3}{8\pi} \dot{M} \Omega_K^2$$



FU Orionis



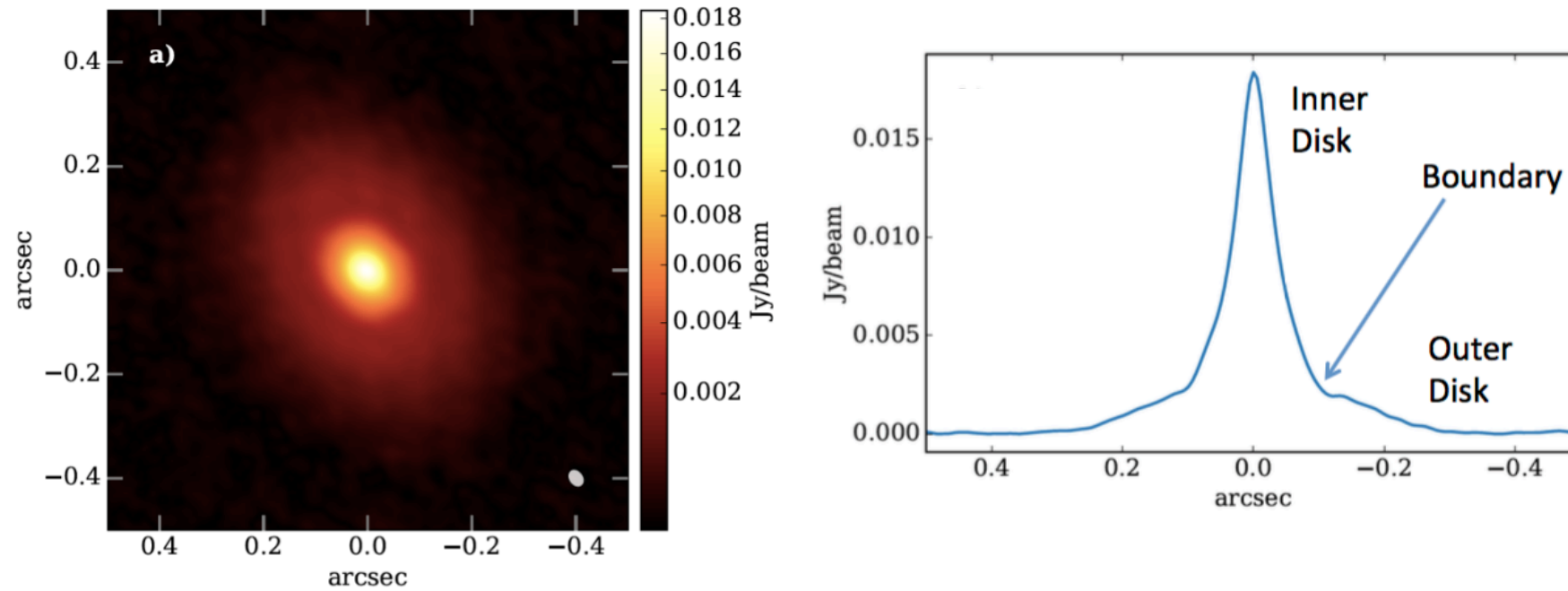
V883Ori



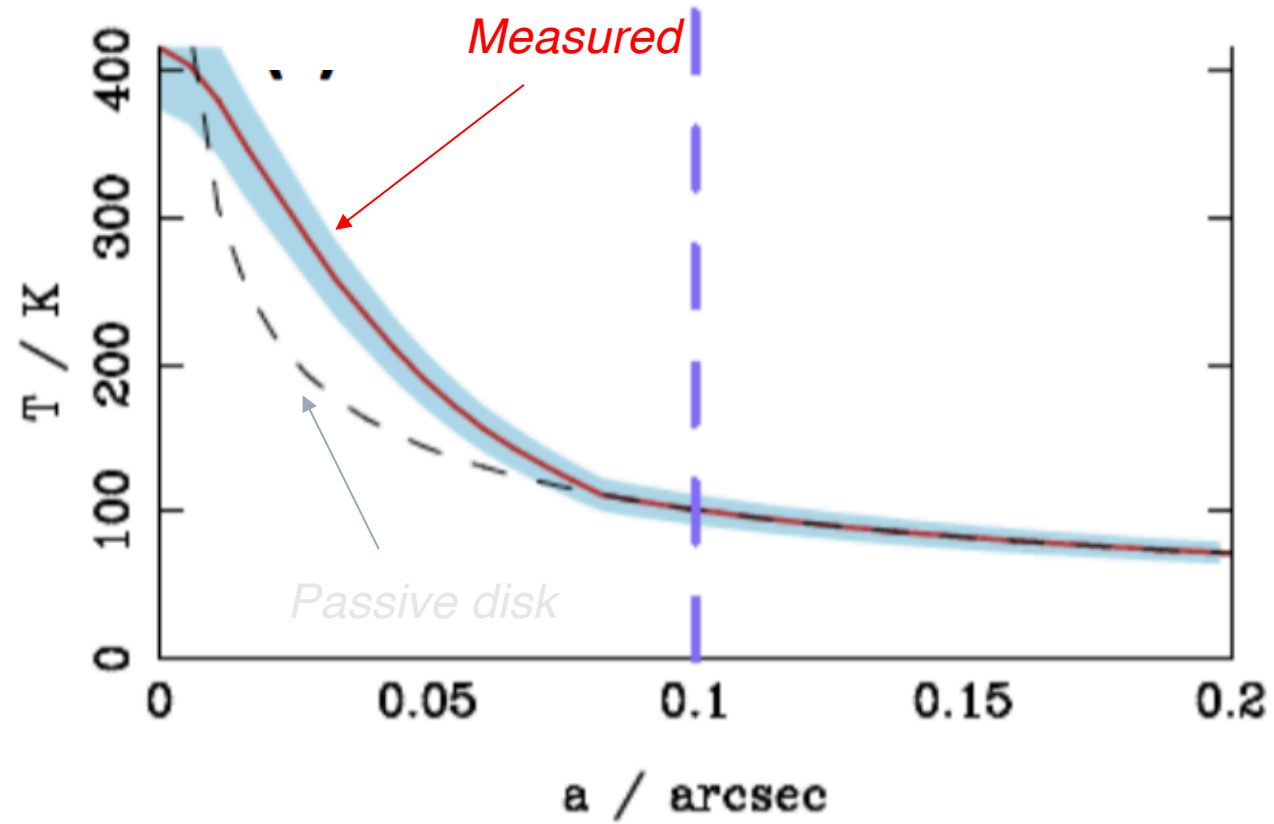
Class I protostar
(Age ~ 0.5 Myr)
($L^* \sim 6 L_{\text{sun}}$)

$M_* = 1.3 \pm 0.1 M_{\text{sun}}$
 $M_{\text{disk}} \sim 0.3 M_{\text{sun}}$
 $d = 414 \pm 7$ pc
 $L_{\text{disk}} = 400 L_{\text{sun}}$

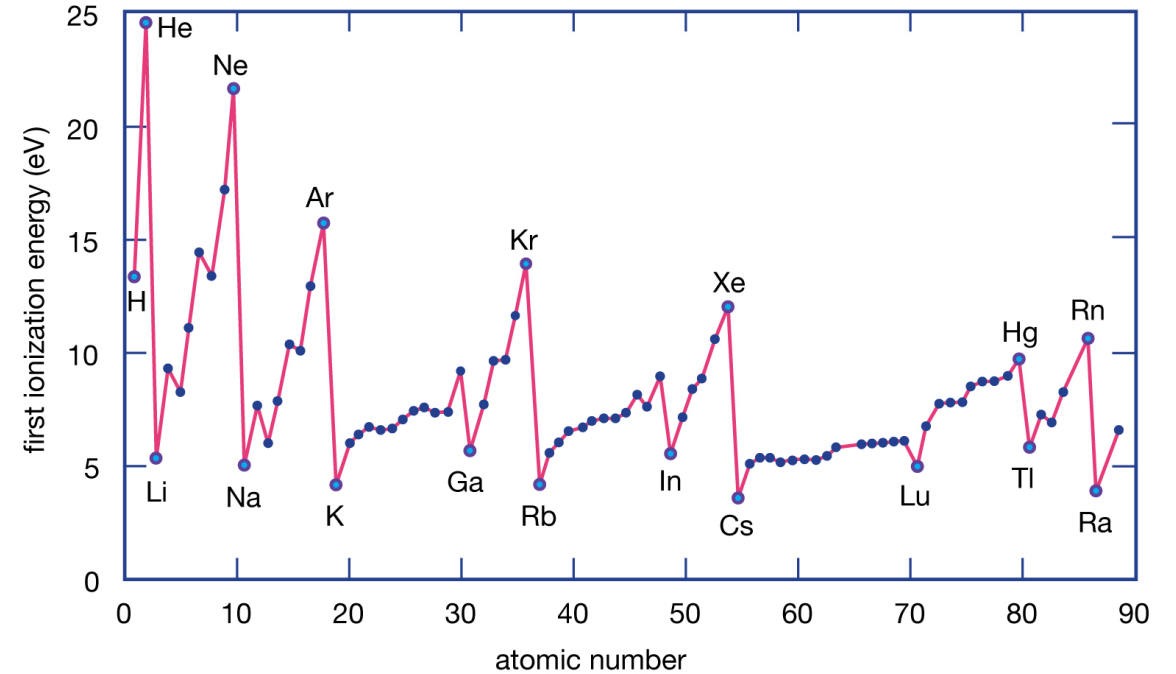
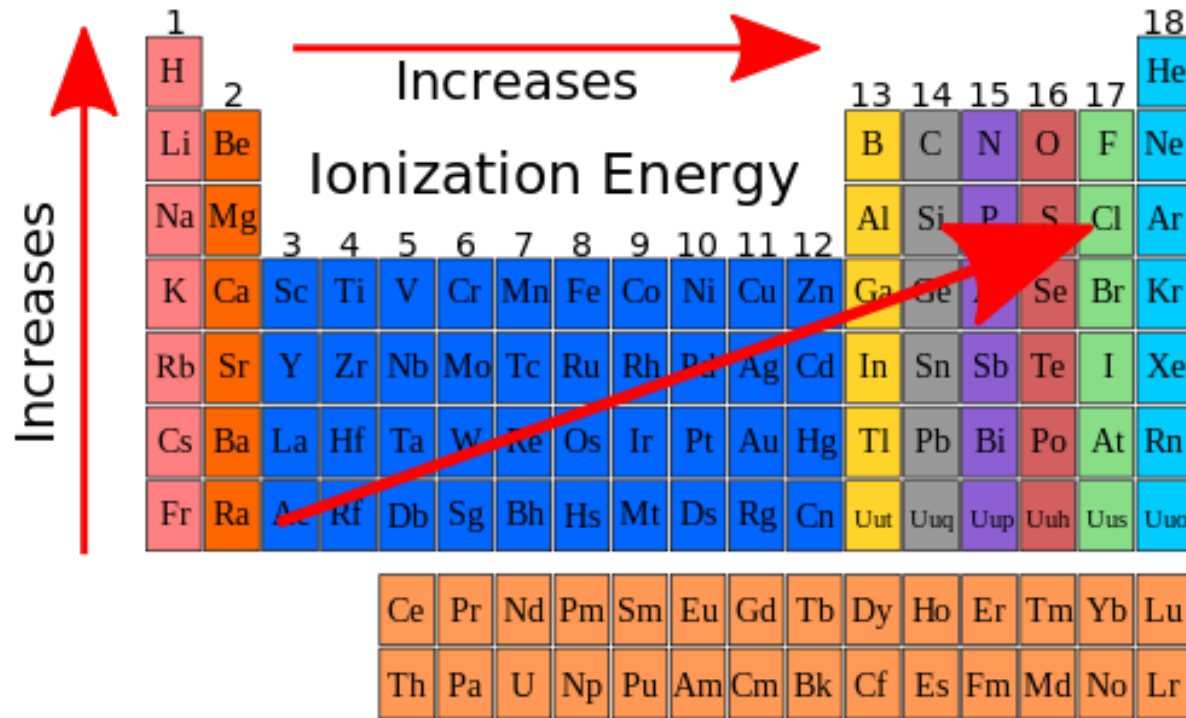
V883Ori



Extra source of heating needed



Ionization energy



© 2012 Encyclopædia Britannica, Inc.

Thermal ionization

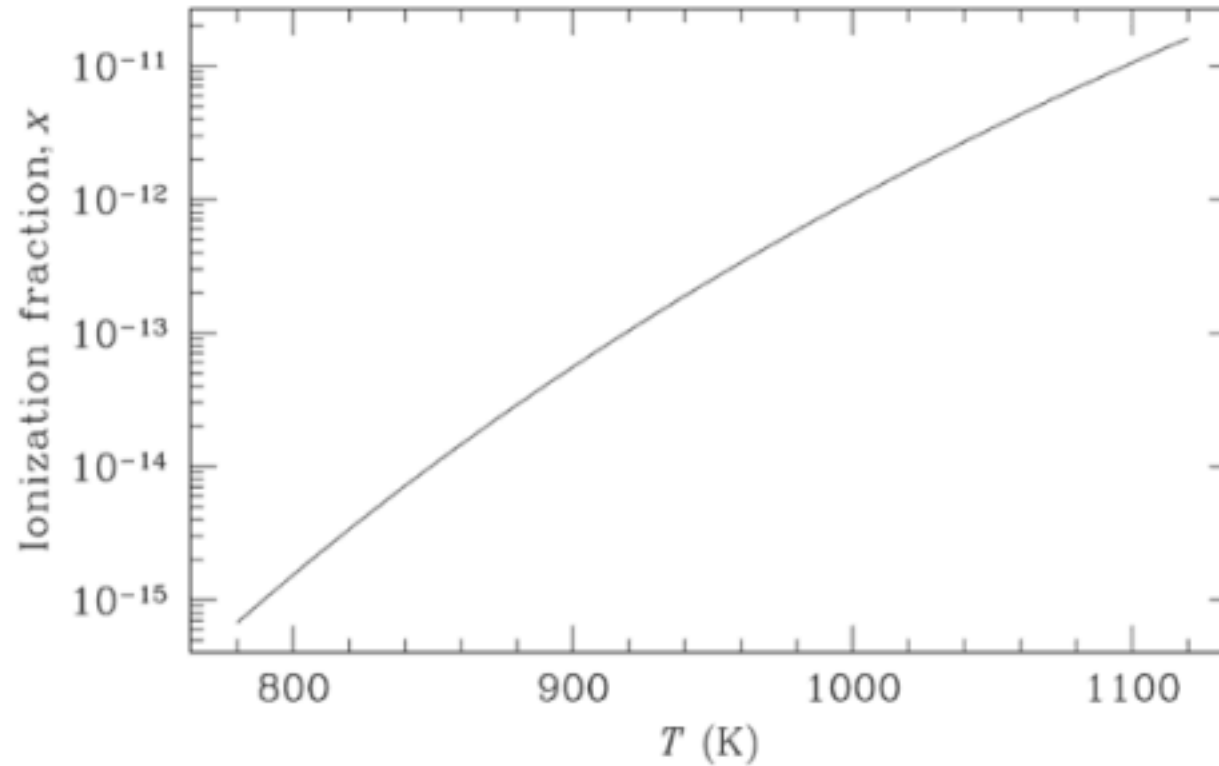
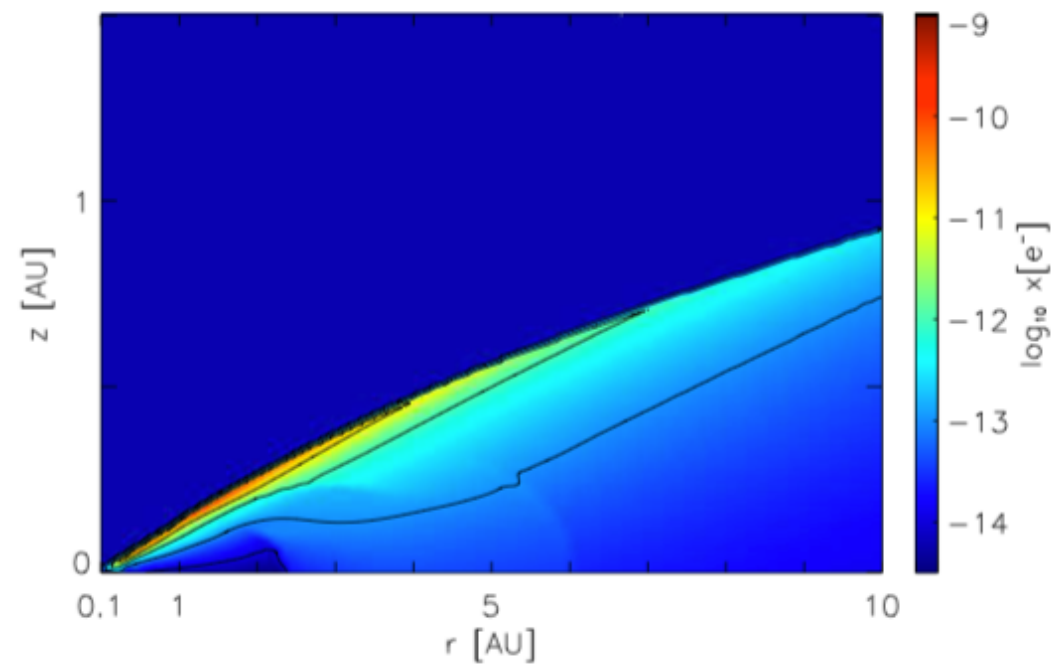
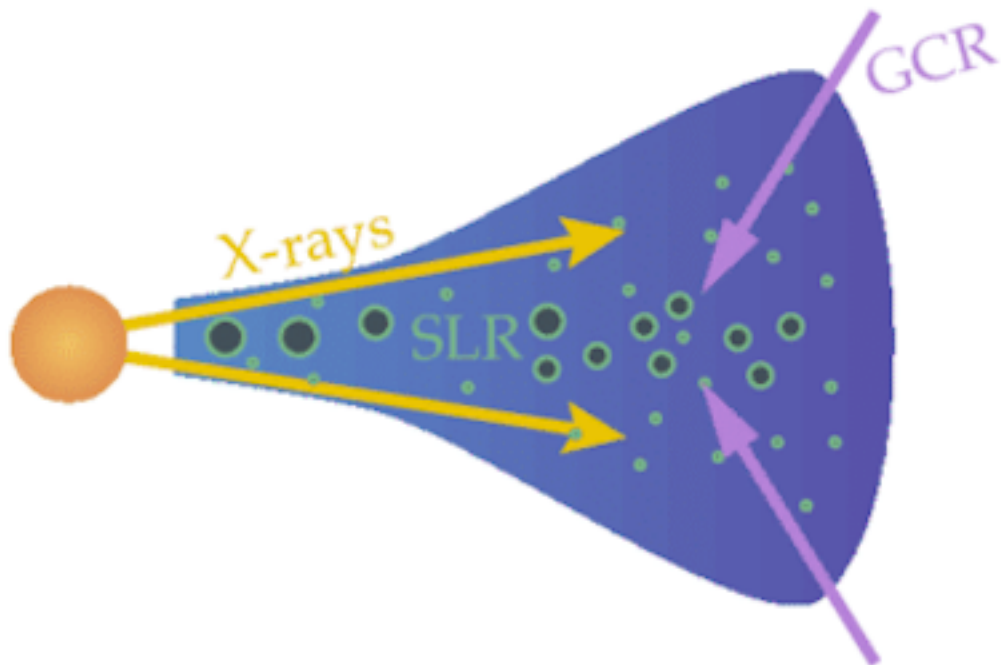


Figure 2.8 The Saha equation prediction for the ionization state of the disk due to the thermal ionization of potassium (ionization potential $\chi = 4.34 \text{ eV}$). The fractional abundance of potassium has been taken to be $f = 10^{-7}$, and the total number density of neutrals $n_n = 10^{15} \text{ cm}^{-3}$.

Non-thermal ionization



Dead zones and “Layered” accretion

VIEW

Abstract

Citations (870)

References (30)

Co-Reads

Similar Papers

Volume Content

Graphics

Metrics

Export Citation

THE ASTROPHYSICAL JOURNAL, 457:355–362, 1996 January 20
© 1996. The American Astronomical Society. All rights reserved. Printed in U.S.A.

LAYERED ACCRETION IN T TAURI DISKS

CHARLES F. GAMMIE

Center for Astrophysics, MS-51, 60 Garden Street, Cambridge, MA 02138

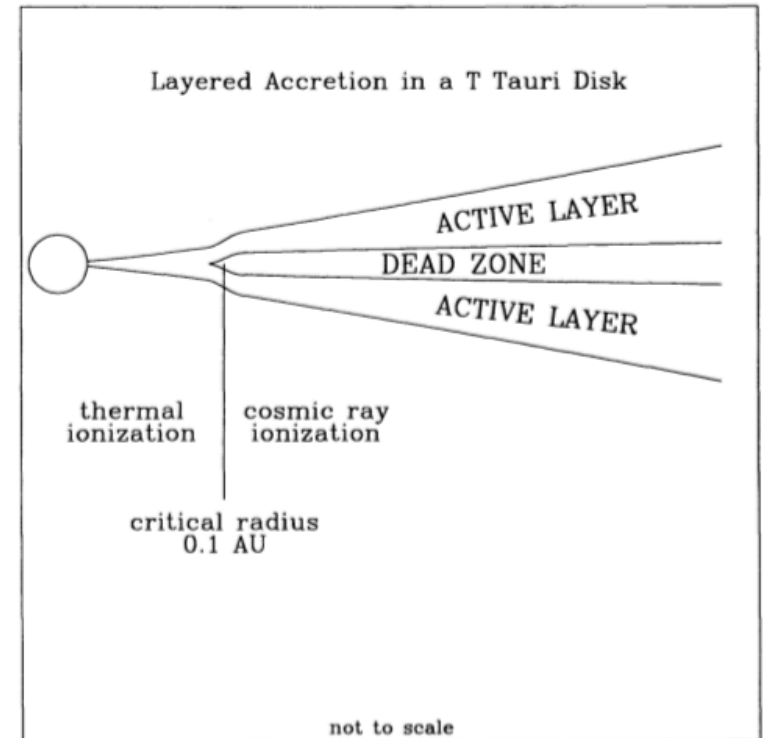
Received 1995 May 8; accepted 1995 July 28

ABSTRACT

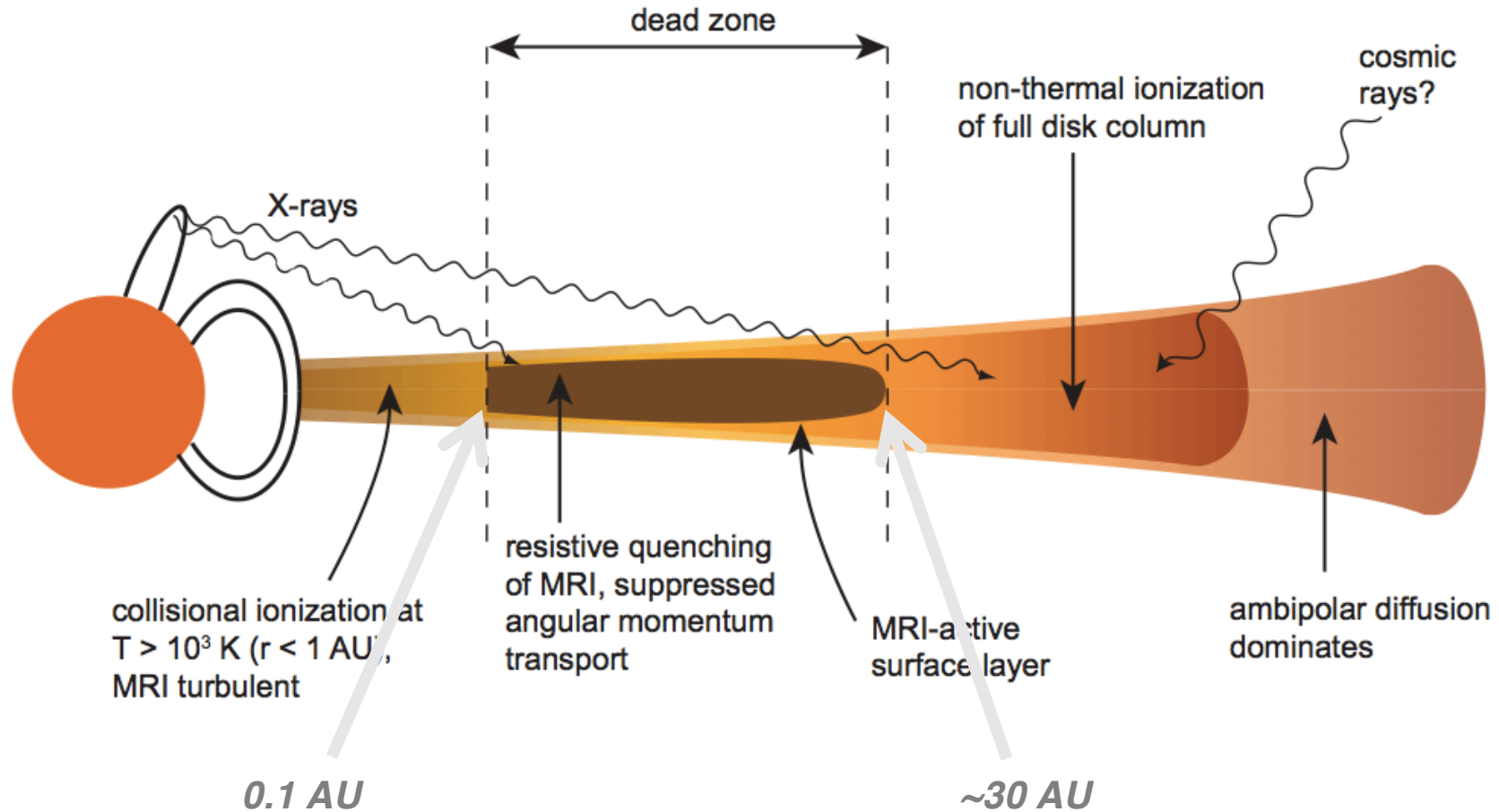
We put forward a model for accretion disks around T Tauri stars. The model assumes that angular momentum transport is driven by magnetic fields and can occur only in those parts of the disk that are sufficiently ionized that the gas can couple to the magnetic field. These regions lie at $R \lesssim 0.1$ AU, where collisional ionization is effective, and at $R \gtrsim 0.1$ AU in a layer of thickness $\approx 100 \text{ g cm}^{-2}$ at the surface of the disk where cosmic-ray ionization is effective.

The model predicts that the stellar accretion rate is about $10^{-8} M_{\odot} \text{ yr}^{-1}$, independent of the rate of infall onto the disk. Matter that is not accreted onto the star accumulates in the inner few AU of the disk at a rate of about $10^{-3} M_{\odot}$ in 10^4 yr. Given this buildup it is unlikely that accretion is steady. The effective temperature profile is $T_e \sim r^{-1/2}$ outside of 0.1 AU, which differs from the canonical $T_e \sim r^{-3/4}$. We calculate the expected spectral energy distribution for the disk and show that this temperature profile produces an infrared excess. Finally, we discuss some of the leading uncertainties in the theory.

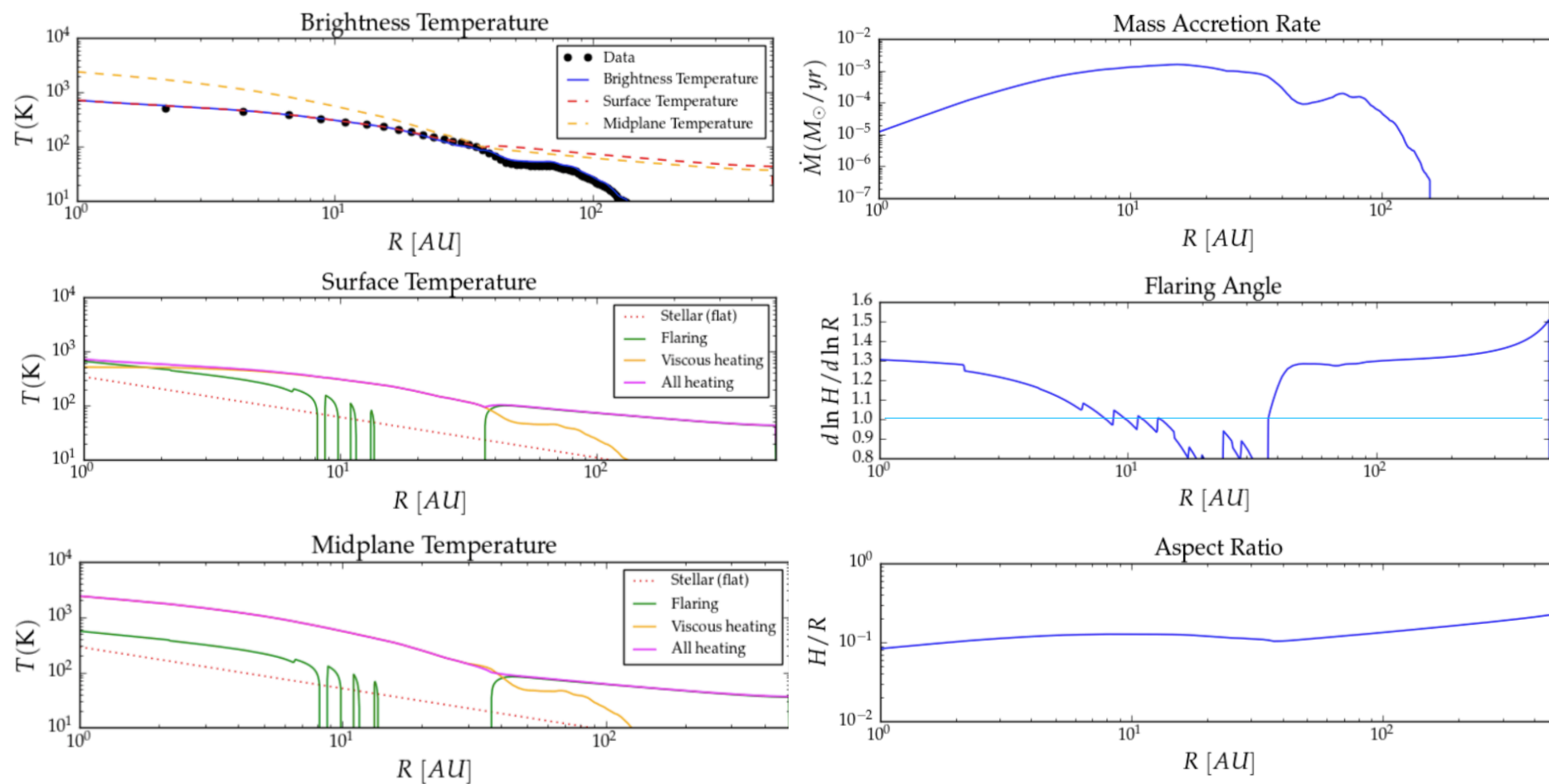
Subject headings: accretion, accretion disks — stars: magnetic fields — stars: pre-main-sequence



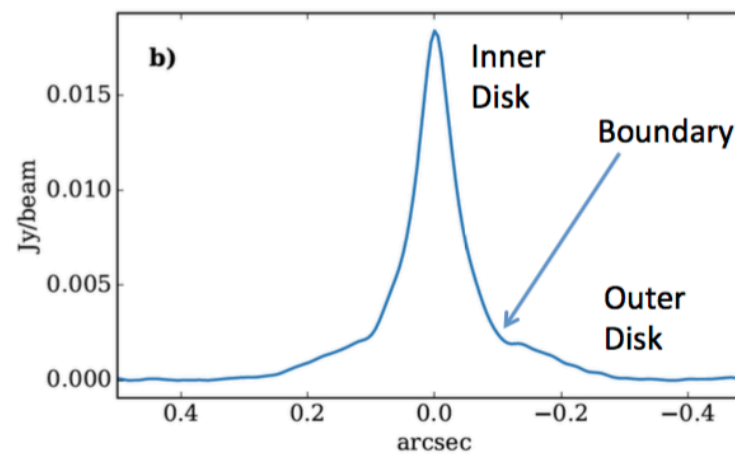
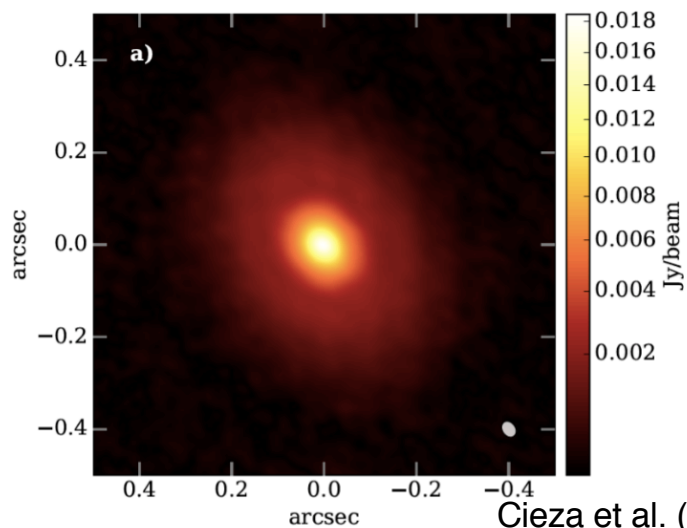
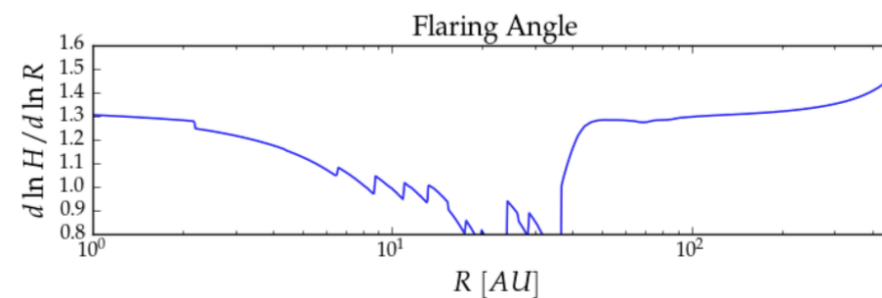
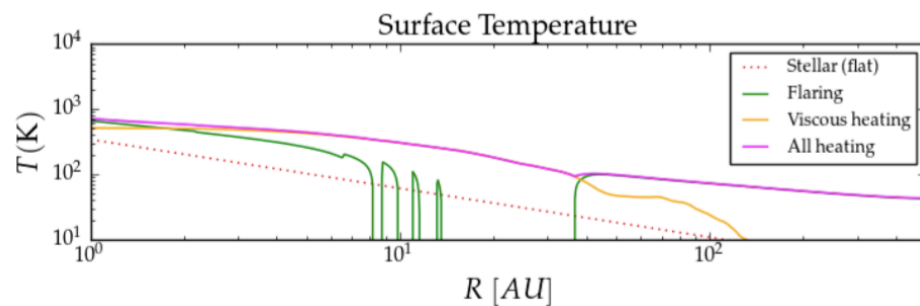
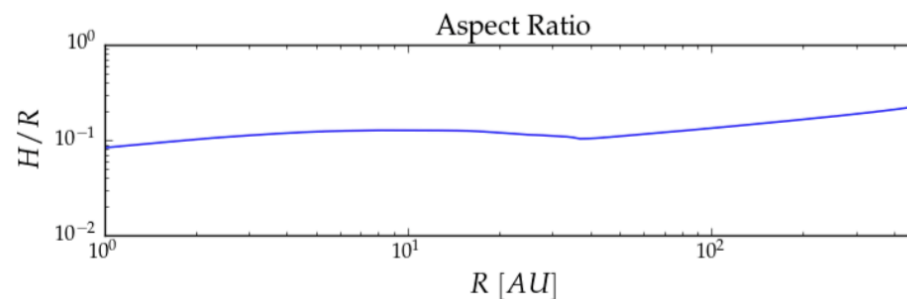
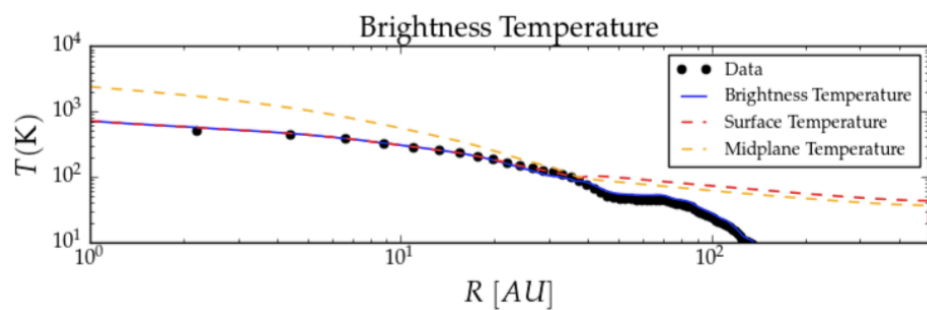
Dead zones



Best fit



Self-shadowing



Cieza et al. (2016), Alarcon et al. (in prep)